



Investigating the potential of MgMOF-74 membranes for CO₂ capture

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ABSTRACT

MgMOF-74 is a metal–organic framework (MOF) with exposed metal cation sites that has one-dimensional 1.1 nm sized hexagonal-shaped channels. On the basis of information available in the published literature, it appears that MgMOF-74 has significant advantages over other MOFs, with respect to its uptake capacity for CO₂. The primary objective of the present communication is to investigate the performance of MgMOF-74 membranes in separating CO₂/H₂, CO₂/N₂, CO₂/CH₄, and CH₄/H₂ mixtures, that are important in carbon capture. To achieve this objective all the parameters required for modeling MgMOF-74 membrane permeation were obtained using molecular simulations. Specifically, Configurational-Bias Monte Carlo (CBMC) simulations were used to determine pure component adsorption isotherms, and isosteric heats of adsorption. Molecular dynamics (MD) simulations were performed to determine the self-diffusivities, $D_{i,self}$, and the Maxwell–Stefan ($M-S$) diffusivities, \mathcal{D}_i , of guest molecules.

The MD simulations show that the zero-loading diffusivity $\mathcal{D}_i(0)$ is consistently lower, by up to a factor of 10, than the values of the Knudsen diffusivity, $D_{i,Kn}$. The ratio $\mathcal{D}_i(0)/D_{i,Kn}$ is found to correlate with the isosteric heat of adsorption, which in turn is a reflection of the binding energy for adsorption at the pore walls. The stronger the binding energy, the lower is the ratio $\mathcal{D}_i(0)/D_{i,Kn}$.

Using the Maxwell–Stefan formulation for binary mixture permeation, along with data inputs from CBMC and MD simulations, the permeation selectivities for CO₂/H₂, CO₂/N₂, CO₂/CH₄, and CH₄/H₂ mixtures were determined for a range of upstream pressures. The model calculations show that increased upstream pressures lead to significant enhancement in permeation selectivities; this enhancement is directly traceable to diffusional correlations within the 1D channels. Such correlations have the effect of slowing-down the more mobile partner species in the mixtures.

MgMOF-74 membrane permeation selectivities for CO₂/H₂, and CO₂/N₂ mixtures are higher than those reported in the published literature with zeolite membranes.

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1. Introduction

In recent years there has been considerable research on the development of novel porous materials such as metal–organic frameworks (MOFs) [1], and zeolitic imidazolate frameworks (ZIFs) [2,3]. MOFs, and ZIFs offer considerable potential for application in a wide variety of separation applications [4–10]. A substantial amount of the published research on metal–organic frameworks (MOFs) is dedicated to the issue of CO₂ capture by selective adsorption [6,9,11–14]. Of particular interest in CO₂ capture applications are MOFs with exposed metal cation sites such as MMOF-74 (=M₂(dobdc)) with M=Mg, Zn, Ni, or Co [15–20]. Yazaydin et al. [20] have screened a total of 14 different MOFs and ZIFs and concluded that MgMOF-74 has the highest CO₂ uptake capacity. For adsorption of CO₂ and CH₄, Bao et al. [21] have found significantly

higher capacities for MgMOF-74 when compared to 13X zeolite. Furthermore, the CO₂/CH₄ adsorption selectivity for MgMOF-74 was found to be comparable to that of 13X zeolite. MgMOF-74 yields CO₂/CH₄ adsorption selectivities that are significantly higher than those reported for CuBTC in the experiments of Hamon et al. [22]. For separation of CO₂/H₂ mixtures, Herm et al. [23] compared the adsorption selectivities of five different MOFs: MgMOF-74, MOF-177, CuBTri, BeBTB, and Co(BDP), with traditionally used adsorbents such as NaX. Their study concludes that while the adsorption selectivity of MgMOF-74 is higher than that of NaX, and has the additional advantage of having a significantly higher capacity. The molecular simulation studies of Gallo et al. [24] and Keskin [25], had concluded, respectively, that ZnMOF-74 and CoMOF-74 yield the highest adsorption selectivities values for CH₄/H₂ separation.

One technological solution is to allow the gaseous mixture to permeate through a porous crystalline layer, with a thickness of a few micrometers [13,25–33]. For the development of such membrane devices, it is essential to have information not only

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on adsorption, but also on the diffusivities of the guest molecules inside the pores.

The primary objective of the present communication is to investigate the performance of MgMOF-74 membranes in separating CO₂/H₂, CO₂/N₂, CO₂/CH₄, and CH₄/H₂ mixtures, that are important in CO₂ capture. To achieve this objective all the parameters required for characterizing the permeation characteristics of MgMOF-74 membranes were obtained using molecular simulations. Specifically, Configurational-Bias Monte Carlo (CBMC) simulations were used to determine pure component adsorption isotherms, and isosteric heats of adsorption. Molecular dynamics (MD) simulations were performed to determine the diffusivities of guest molecules. Besides H₂, N₂, CO₂, and CH₄, we also investigated the diffusion characteristics of higher alkanes: ethane (C2), propane (C3), and n-butane (nC4) in order to establish a correlation between diffusivities and binding energies. Furthermore, for comparison purposes and to obtain a broader perspective on membrane permeation, the adsorption and diffusion characteristics of ZnMOF-74 [15,16,19,20], that is iso-structural with MgMOF-74, were also investigated.

Supplementary material accompanying this publication gives information on the specific force fields used, unit cell dimensions, accessible pore volume, characteristic pore dimensions, molecular simulation methodologies, pore landscapes and snapshots, along with the entire data base of simulation data on isotherms, heats of adsorption, and diffusivities.

2. Maxwell–Stefan model and diffusivities

For modeling of intra-crystalline mixture diffusion in membrane permeation devices it is commonly accepted that the fundamentally correct approach is to relate the fluxes N_i , defined in terms of the cross-sectional area of the membrane, to the chemical potential gradients by use of the Maxwell–Stefan (M – S) equations [34–36]

$$-\phi \frac{c_i}{RT} \frac{d\mu_i}{dz} = \sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_j N_i - x_i N_j}{\mathcal{D}_{ij}} + \frac{N_i}{\mathcal{D}_i}; \quad i = 1, 2, \dots, n \quad (1)$$

where ϕ represents the fractional pore volume of the porous material, and the concentrations c_i are defined in terms of moles per m³ of accessible pore volume. The x_i in Eq. (1) is the component mole fractions of the adsorbed phase

$$x_i = \frac{c_i}{c_t}; \quad i = 1, 2, \dots, n \quad (2)$$

There are two kinds of diffusivities that are needed. Firstly we need information on the \mathcal{D}_i , that characterize species i –wall interactions. In many cases the \mathcal{D}_i corresponds to the value of the pure component i ; consequently this can be estimated from unary diffusion data. In some exceptional cases the \mathcal{D}_i in the mixture can be lower than that of pure components [37–39]. In the published literature there is only one experimental study, that reports the M – S diffusivities of CO₂ and CH₄ in MgMOF-74 [21].

Secondly, we need data on the \mathcal{D}_{ij} ; these are exchange coefficients representing interaction between components i with component j . At the molecular level, the \mathcal{D}_{ij} reflect how the facility for transport of species i correlates with that of species j . Conformity with the Onsager reciprocal relations prescribes

$$\mathcal{D}_{ij} = \mathcal{D}_{ji} \quad (3)$$

The relative contributions of the first and second right members of Eq. (1) are dictated by the *degree of correlations*, defined as $\mathcal{D}_i/\mathcal{D}_{ij}$; the higher the degree of correlations, the more important is the contribution of the first right member. We will see later

that correlation effects are of paramount importance in MgMOF-74 membrane permeation.

In earlier published formulations on multicomponent diffusion in zeolites and MOFs [40,41], the M – S equations are set up differently using loadings, q_i , expressed as moles per kg of framework; the inter-relationship between c_i and q_i is

$$c_i = \frac{q_i}{V_p} \quad (4)$$

where V_p is the accessible pore volume. For MgMOF-74 the value of V_p , obtained from molecular simulations using the helium probe insertion technique suggested by Talu and Myers [42,43], is 0.78 cm³/g. Furthermore, the earlier formulations use fractional occupancies, θ_i , defined by

$$\theta_i \equiv \frac{c_i}{c_{i,sat}} \quad (5)$$

in place of the adsorbed phase mole fractions, x_i . For a detailed comparison of the earlier approaches with the current one, using Eq. (1), see Supplementary material accompanying the paper of Krishna and Van Baten [34]. There are important advantages in using the current approach using moles per accessible pore volume, c_i . Firstly, Eq. (1), is applicable equally to micro- and meso-porous materials, as explained in our earlier papers [34,35,44–47]. Secondly, the use of the current formulation allows a much clearer interpretation of the M – S diffusivities \mathcal{D}_{ij} ; these are relatable to the corresponding M – S diffusivity in the *binary* fluid mixture [34,35].

The MD data on the M – S diffusivities, \mathcal{D}_i , for H₂, N₂, CO₂, CH₄, C2, C3, and nC4 diffusion in the 1D channels of MgMOF-74 are shown in Fig. 1a as a function of the pore concentrations, c_i . It is interesting to note that the hierarchy of \mathcal{D}_i values: H₂ > N₂ > CH₄ > C2 > C3 ≈ CO₂ ≈ nC4 does not appear to correlate with either the molar mass, M_i , or the kinetic diameter; this needs further investigation and analysis.

In the limit of vanishingly small loadings, $c_i \rightarrow 0$, the zero-loading diffusivity value $\mathcal{D}_i(0)$ is dictated primarily by molecule-wall collisions. When the reflections are purely *diffuse* in nature, i.e. the angle of reflection bears no relation to the angle of incidence at which the molecule strikes the pore wall, the $\mathcal{D}_i(0)$ value corresponds to that obtained by the classic Knudsen formula

$$D_{i,Kn} = \frac{d_p}{3} \sqrt{\frac{8RT}{\pi M_i}} \quad (6)$$

Eq. (6) holds in the limiting case when the molecule does not adsorb at pore walls [34,35,48–50].

MgMOF-74 has one-dimensional (1D) hexagonal-shaped channels; see Fig. 2. The effective channel diameter d_p was determined following the method of Delaunay triangulation, described in the work by Foster et al. [51]. The maximum diameter of a sphere that can pass through the channels is determined to be 1.1 nm. Taking $d_p = 1.1$ nm, the application of Eq. (6) yields $D_{i,Kn} = 65 \times 10^{-8} \text{ m}^2 \text{ s}^{-1}$, which value is slightly higher than the value of $\mathcal{D}_i(0) \approx 55 \times 10^{-8} \text{ m}^2 \text{ s}^{-1}$ obtained from MD. For ZnMOF-74, that has channel dimensions almost identical to that of MgMOF-74, we note that the $\mathcal{D}_i(0)$ for H₂ is practically the same as the Knudsen value; see Fig. 2b. H₂ has a poor adsorption strength, and it is therefore not surprising that the diffuse reflectance assumption is a reasonable one, and the Knudsen prescription holds.

For species other than H₂, adsorption at the pore wall is significant, the values of $\mathcal{D}_i(0)$ show considerable deviation from the values of $D_{i,Kn}$. For example, the diffusivity of CO₂ in MgMOF-74 is seen to be five times lower than that of CH₄, in excellent agreement with the value obtained in the experiments of Bao et al. [21]. However, on the basis of Eq. (6) we should expect this factor to be only lower by a factor $\sqrt{44/16} = 1.66$. In MgMOF-74 the values of $\mathcal{D}_i(0)/D_{i,Kn}$ are 0.43, 0.25, 0.17, 0.108, 0.123, and 0.124 for N₂, CH₄,

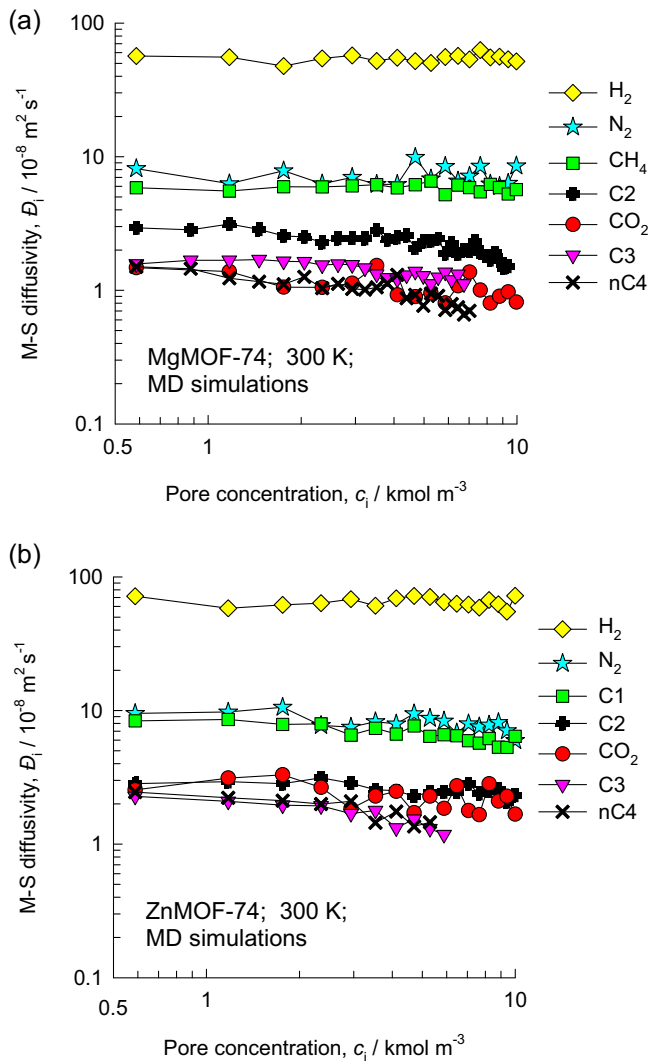


Fig. 1. MD simulations of M - S diffusivities, D_i , of H_2 , N_2 , CO_2 , CH_4 , C_2 , C_3 , and nC_4 in the 1D channels of (a) MgMOF-74, and (b) ZnMOF-74 at 300 K, expressed as a function of the pore concentrations, c_i .

C_2 , CO_2 , C_3 , and nC_4 , respectively. In ZnMOF-74 the corresponding values $D_i(0)/D_{i,Kn}$ values for these species are 0.49, 0.28, 0.24, 0.18, 0.18, and 0.174.

Let us examine whether the extent of departures from the Knudsen prescription correlates with the binding energy of the molecules. A measure of the binding energy is the isosteric heats of adsorption, $-\Delta H_{st}$. CBMC simulation data on $-\Delta H_{st}$ are presented in Fig. 3a and b for MgMOF-74 and ZnMOF-74 respectively. We note that CO_2 has the highest heat of adsorption, reflecting the strongest binding energy amongst all the investigated species. For MgMOF-74, the ratio $D_i(0)/D_{i,Kn}$ is seen to correlate well with the isosteric heat of adsorption; see Fig. 4a. This is a rational result; the higher the binding energy, the higher is the sticking tendency of that species with the pore wall, leading to greater departure from the Knudsen prescription of diffuse reflectance. It is also interesting to note that the $D_i(0)/D_{i,Kn}$ values tend to reach an asymptote value of about 0.1. The higher the binding energy the higher is the probability that a molecule hops to a neighboring adsorption site rather than return to the bulk gas phase after collision with the pore wall [35,49,52]. This bias is best appreciated by viewing video animations of MD simulations showing the hopping of H_2 , N_2 , CH_4 , C_2 , C_3 , and nC_4 within the 1D channel MgMOF-74; these Video animations have been provided as [Supplementary material](#) accompanying this pub-

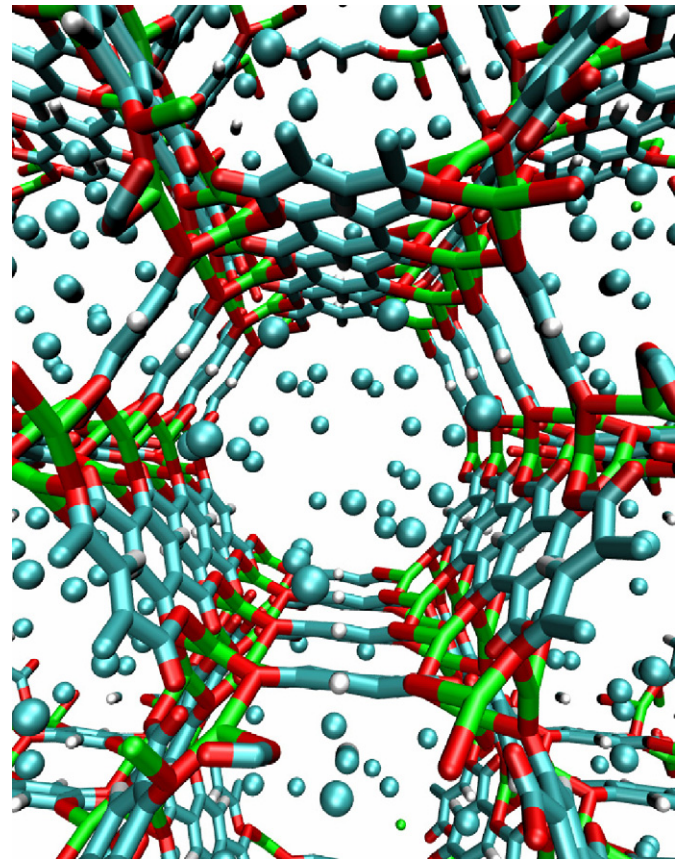


Fig. 2. Snapshots of adsorbed CH_4 molecules within the 1D channels of MgMOF-74.

lication. The trajectories followed by H_2 molecules, as seen in the animations, clearly show that a H_2 molecule that strikes the pore wall has a tendency to return to the bulk gas phase largely in keeping with the diffuse reflection scenario prescribed by the Knudsen theory. For all other species, the adsorption at the pore walls causes many molecules to jump to a neighboring site rather than return to the bulk gas phase.

Fig. 4b presents the corresponding correlation of the ratios $D_i(0)/D_{i,Kn}$ with the isosteric heats of adsorption for ZnMOF-74. In this case, the value of the ratio $D_i(0)/D_{i,Kn}$ for CO_2 is 0.18. This higher value as compared to MgMOF-74 is because of the lower binding energy of CO_2 ; this is reflected in the heat of adsorption value of 25 kJ mol^{-1} , compared to 32 kJ mol^{-1} for MgMOF-74.

The information in Fig. 4 provides a practical methodology for estimation of the $D_i(0)$ using the $D_{i,Kn}$ value as a starting point.

3. Correlation effects, and diffusion selectivities

Applying Eq. (1) to a mixture of tagged and untagged species i , we can derive the following expression for the self-diffusivities, $D_{i,self}$, for unary diffusion [35]

$$\frac{1}{D_{i,self}} = \frac{1}{D_i} + \frac{1}{D_{ii}} \quad (7)$$

The self-exchange coefficients, D_{ii} , are starting points for the estimation of the D_{ij} . The ratio D_i/D_{ii} is a measure of the *degree of correlations* for unary diffusion. Fig. 5a and b shows the MD simulated data for $D_{i,self}$, along with the ratio D_i/D_{ii} for various guest molecules in MgMOF-74. For any species, the D_i/D_{ii} increases with increasing pore concentration, c_i . As the concentration within the pores increases, the jumps of individual molecules become increasingly correlated because the number of vacant adsorption sites will

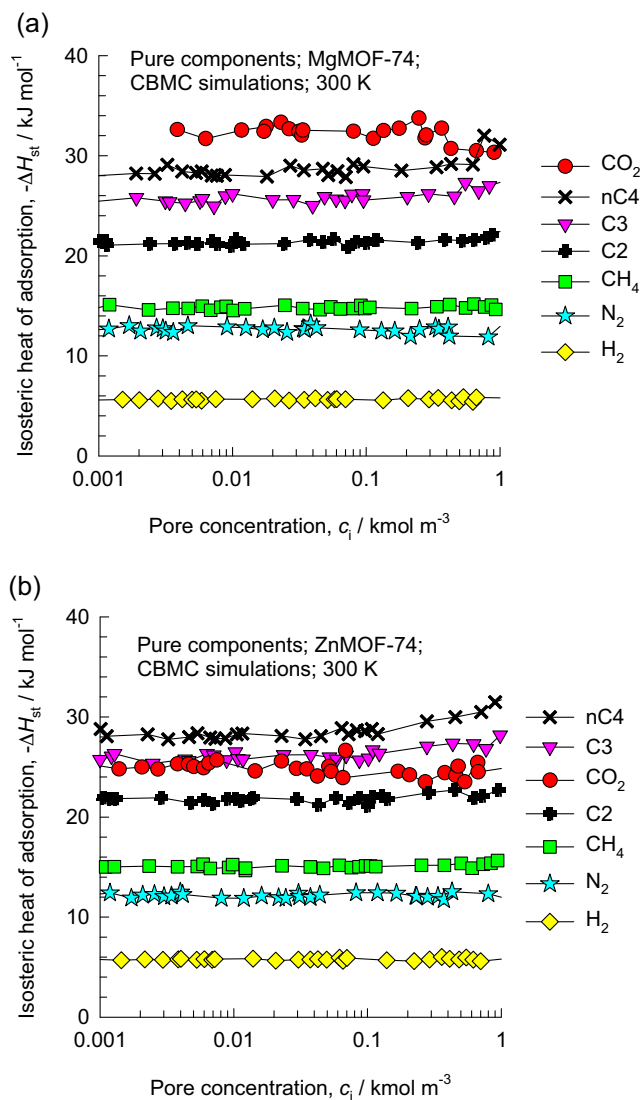


Fig. 3. CBMC simulation data on the isosteric heats of adsorption, $-\Delta H_{st}$, in (a) MgMOF-74, and (b) ZnMOF-74 of H_2 , N_2 , CO_2 , CH_4 , C_2 , C_3 , and nC_4 at 300 K.

decrease, and the probability of unsuccessful jumps will increase. Fig. 5c compares the degree of correlations for methane diffusion in a variety of zeolites and MOFs. The strongest degree of correlations are within the 0.74 nm sized 1D channels of MIL-53(Cr). The degree of correlations in MgMOF-74, are similar to that of ZnMOF-74, and both are comparable with those found for the intersecting channels of MFI zeolite. Published MD simulations also show that the inter-cage hopping of molecules across the narrow windows of DDR, CHA, LTA, and ITQ-29 zeolites is practically uncorrelated, and a good approximation is to take $\bar{D}_i/\bar{D}_{ii} \approx 0$ for such structures [38,40,41].

It is worth pointing out here that in our earlier work on the development of the M - S equations for micropore diffusion we had suggested the approximation $\bar{D}_i/\bar{D}_{ii} \approx 1$ [53,54]. Subsequently, MD simulation results became available to show the \bar{D}_i/\bar{D}_{ii} varies with pore concentrations for all guest-host combinations [40,55]. It is remarkable, therefore, to note that the recent paper by Lito et al. [56] invokes the assumption $\bar{D}_i/\bar{D}_{ii} = 1$ to describe diffusion across both MFI membranes. As seen from the data in Fig. 5c, and in the extensive MD simulation results for MFI presented earlier [40,55], \bar{D}_i/\bar{D}_{ii} for guest molecules such as CO_2 , and linear alkanes in MFI increases significantly with pore concentration. The perme-

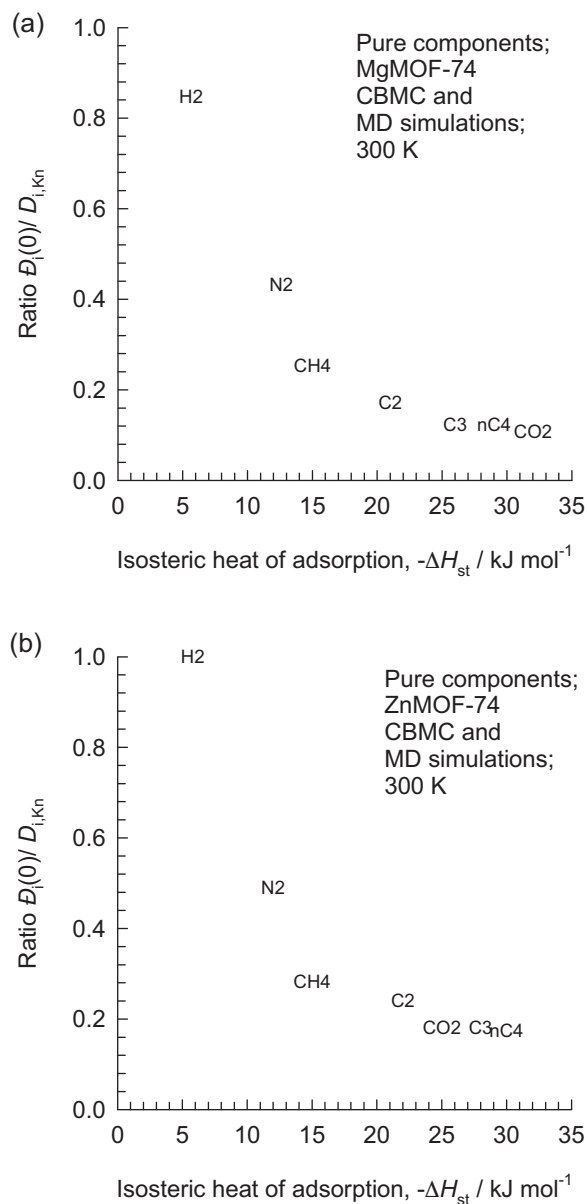


Fig. 4. The ratio of \bar{D}_i to the Knudsen diffusivity, $\bar{D}_i(0)/\bar{D}_{i,Kn}$, for (a) MgMOF-74, and (b) ZnMOF-74, plotted as function of the isosteric heat of adsorption $-\Delta H_{st}$ of the corresponding species.

ation calculations of Lito et al. [56] must therefore be viewed with caution. Indeed, an important message that emerges later in this paper is that the *increase* in the degree of correlations with pore concentrations is a key to the success of MgMOF-74 membrane separations.

For mixture diffusion, correlations in molecular jumps have the effect of slowing down the more mobile species because the tardier species does not vacate the adsorption site quick enough. In order to demonstrate the slowing down effect let us consider diffusion in binary mixture $CO_2(1)/H_2(2)$ with equal concentrations, i.e. $c_1 = c_2$. The slowing-down effect is quantified by the diffusion selectivity, S_{diff} . Starting with the M - S Eq. (1) the following expression for S_{diff} can be derived [13]

$$S_{diff} = \frac{\bar{D}_1}{\bar{D}_2} \frac{1 + (\bar{D}_2/\bar{D}_{12})}{1 + (\bar{D}_1/\bar{D}_{12})} \quad (8)$$

The derivation of Eq. (8) is provided in [Supplementary material](#), along with an explanation on how MD mixture simulations are

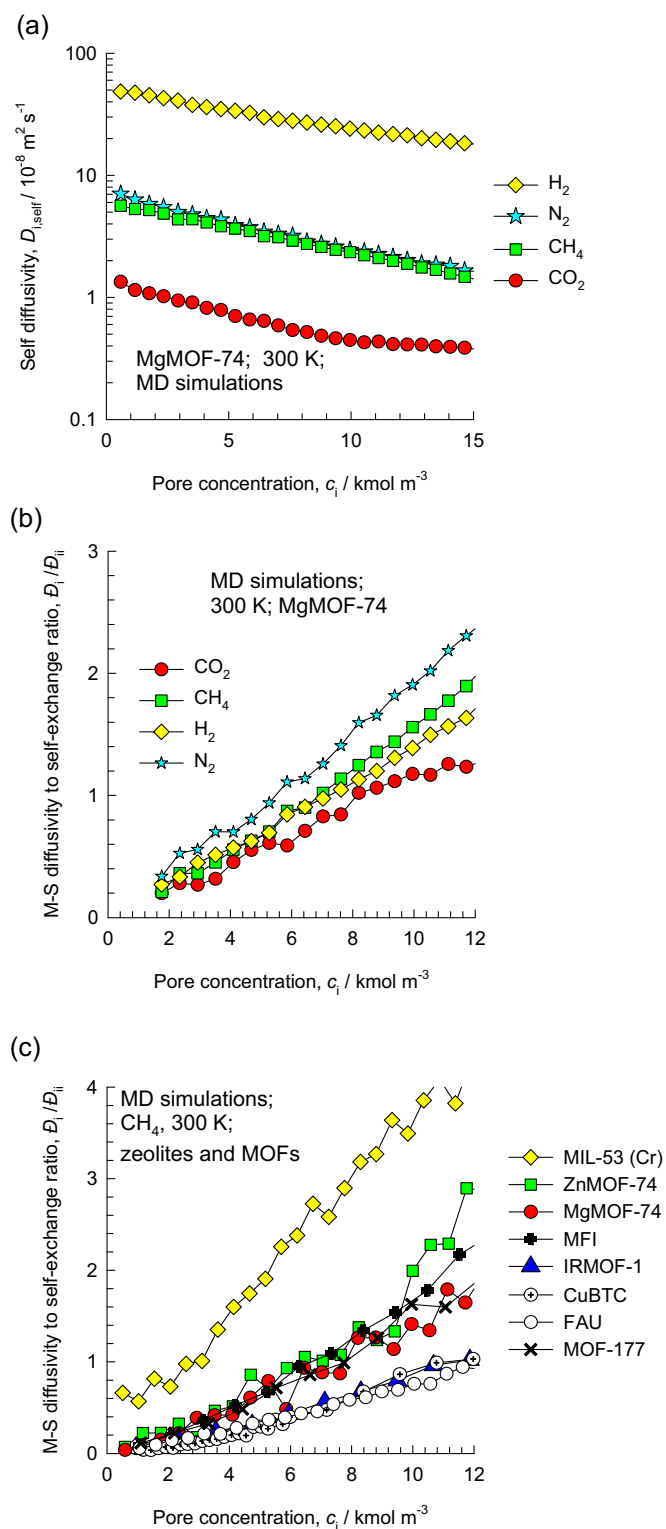


Fig. 5. MD simulations of (a) self-diffusivities $D_{i,self}$, and (b) the degree of correlations, D_i/D_{ii} , for diffusion of H_2 , N_2 , CO_2 , CH_4 , C_2 , C_3 , and nC_4 in the 1D channels of MgMOF-74 at 300 K, expressed as a function of the pore concentrations, c_i . (c) Comparison of the degree of correlations, D_i/D_{ii} , for diffusion of CH_4 in a variety of zeolites (MFI, FAU) and MOFs (IRMOF-1, CuBTC, MIL-53(Cr), MOF-177, MgMOF-74) at 300 K, expressed as a function of the pore concentrations, c_i .

used for the calculations. As an approximation, and for screening purposes, the right member of Eq. (8) is often assumed to be equal to the ratio of the self-diffusivities $D_{i,self}$ in the mixture [13,49,50]

$$S_{diff} = \frac{D_{1,self}}{D_{2,self}} \quad (9)$$

Eq. (9) is a reasonable approximation when the pore concentrations significantly lower than those corresponding to saturation conditions. For MgMOF-74, the CO_2 adsorption strength is high and this approximation is not a good one to use.

For a variety of MOFs, the right member of Eq. (8), determined from CO_2/H_2 mixture simulations, are plotted in Fig. 6a as a function of the total pore concentration, c_t . Also plotted in Fig. 6a is the value of the Knudsen selectivity

$$\frac{D_{1,Kn}}{D_{2,Kn}} = \sqrt{\frac{M_2}{M_1}} \quad (10)$$

For all the MOFs, the S_{diff} is lower the Knudsen selectivity at low pore concentrations; this is because CO_2 has a significantly higher binding energy than H_2 . Indeed, the hierarchy of values of S_{diff} in Fig. 6a is inverse to the hierarchy of the isosteric heats of adsorption of CO_2 in the corresponding MOFs; cf. Fig. 6b. The stronger the binding energy of CO_2 , the lower is the diffusion selectivity. With increasing pore concentration, the correlations become stronger. This has the effect of slowing down the more mobile H_2 ; consequently S_{diff} increases with increasing c_t . It is noteworthy that with increasing degree of correlations, the Knudsen selectivity value can be exceeded.

Fig. 7a and b presents data on the diffusion selectivity for diffusion of CO_2/N_2 , and CO_2/CH_4 mixtures in MgMOF-74 and ZnMOF-74. For both mixtures, S_{diff} for MgMOF-74 is significantly lower than for ZnMOF-74, due to the higher binding energy of CO_2 . For CH_4/H_2 mixture diffusion the S_{diff} values in MgMOF-74 and ZnMOF-74 are the same; see Fig. 7c. This is because the binding energies of CH_4 are practically identical in these two MOFs.

The data presented in Figs. 6 and 7 lead to the inescapable conclusion that the adsorption selectivity does not go hand in hand with diffusion selectivity. This conclusion is generally valid for “open” microporous materials such as zeolites, MOFs, and ZIFs that have pore sizes larger than about 0.74 nm [13,14,35]. For structures such as LTA, CHA, DDR, and ZIF-8 that have cages separated by narrow windows in the 0.35–0.41 nm size range, it is possible to have the diffusion and adsorption selectivities complement each other [40,57].

4. Membrane permeation

We now consider permeation of (a) $\text{CO}_2(1)/\text{H}_2(2)$, (b) $\text{CO}_2(1)/\text{N}_2(2)$, (c) $\text{CO}_2(1)/\text{CH}_4(2)$, and (d) $\text{CH}_4(1)/\text{H}_2(2)$ mixtures across a MgMOF-74 membrane with microporous crystalline layer of thickness $l = 50 \mu\text{m}$. The microporous layer consists of pure MOF-74 crystals, without any supporting matrix polymer. There is however a practical issue relating to the orientation, and alignment of all the 1D channels. For the purposes of the calculations presented here, it is assumed that all the crystals are aligned in such a way that the 1D channels are parallel to the direction of gas flow. For all mixtures the species 1 has the higher adsorption strength, while the species 2 has higher diffusivity.

The first step is to re-write the chemical potential gradients in Eq. (1) in a more practically usable form. The chemical potential gradients can be related to the more conventional gradients of con-

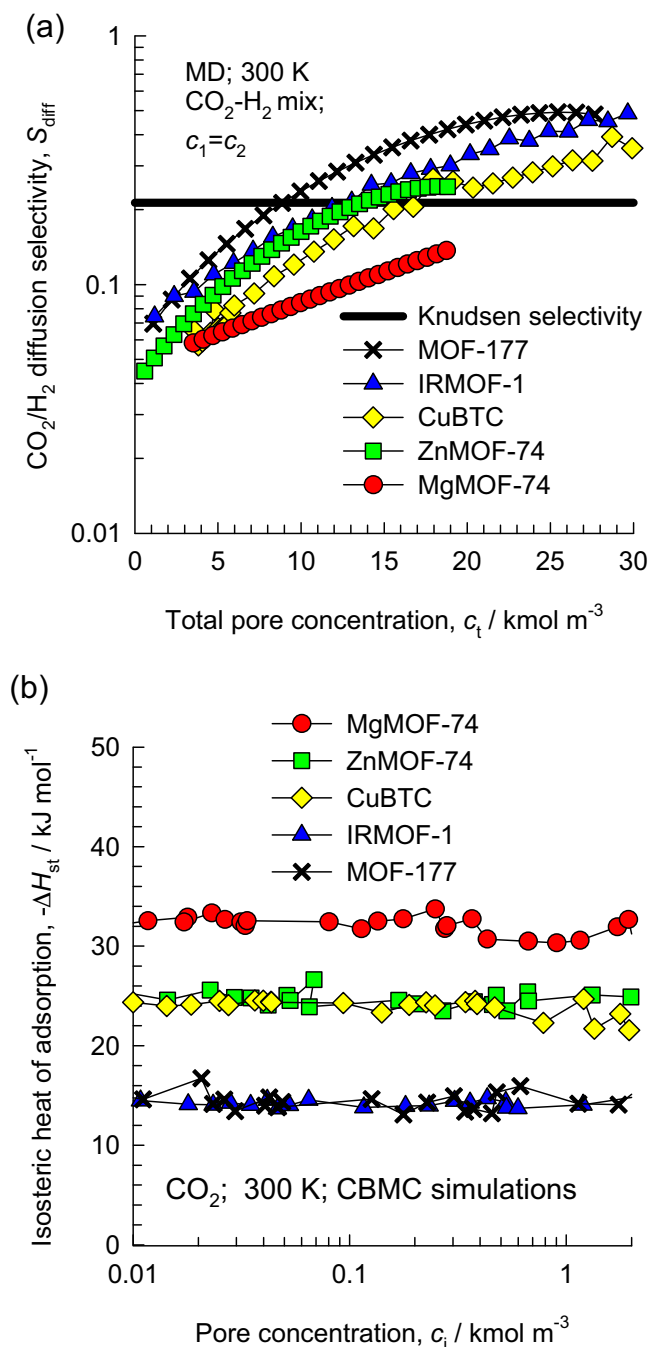


Fig. 6. (a) The diffusion selectivity, S_{diff} , for equimolar ($c_1 = c_2$) CO_2/H_2 mixtures in MgMOF-74, ZnMOF-74, MOF-177, IRMOF-1, and CuBTC at 300 K, plotted as a function of the total pore concentration, $c_t = c_1 + c_2$. (b) CBMC simulation data on the isosteric heats of adsorption, $-\Delta H_{st}$ of CO_2 in various MOFs plotted against the pore concentrations, c_i . The MD data for MOF-177, IRMOF-1, and CuBTC are from earlier work [13].

centrations within the pores by defining thermodynamic correction factors Γ_{ij}

$$\frac{c_i}{RT} \frac{d\mu_i}{dz} = \sum_{j=1}^2 \Gamma_{ij} \frac{dc_j}{dz}; \quad \Gamma_{ij} = \frac{c_i}{p_i} \frac{\partial p_i}{\partial c_j}; \quad i, j = 1, 2 \quad (11)$$

For determination of the adsorbed phase pore concentrations c_i and the elements Γ_{ij} , in the binary mixtures, the Ideal Adsorbed Solution Theory (IAST) of Myers and Prausnitz [58] was used. For application of the IAST, pure component isotherm fits are required.

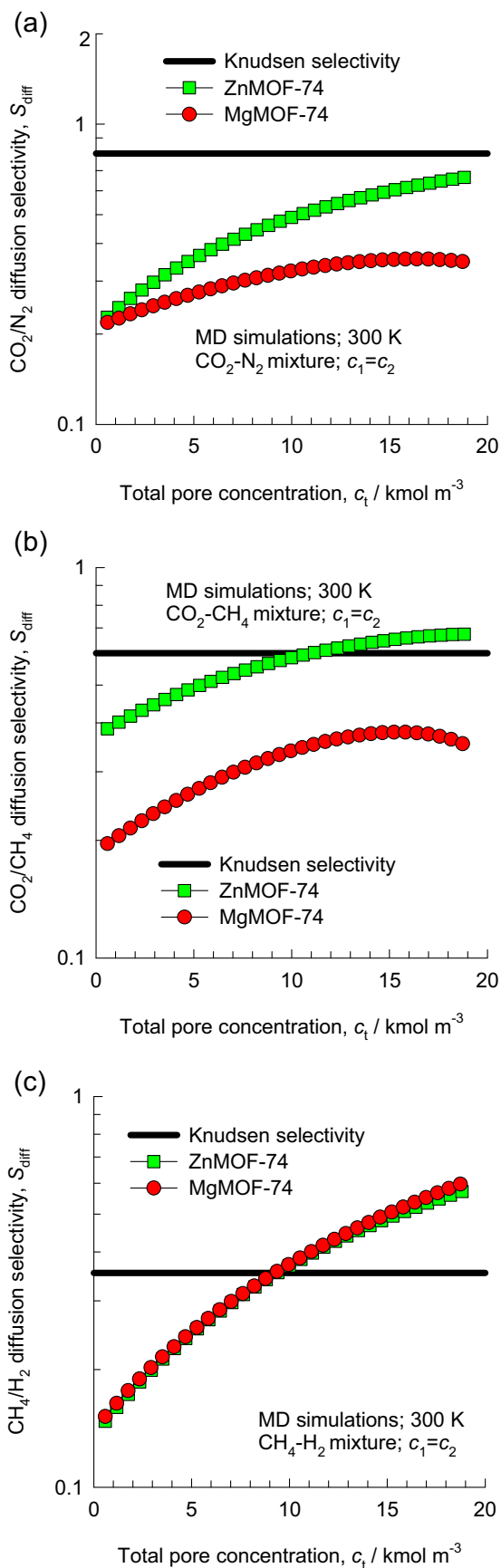


Fig. 7. The diffusion selectivity, S_{diff} , for equimolar ($c_1 = c_2$) (a) CO_2/N_2 , (b) CO_2/CH_4 , and (c) CH_4/H_2 mixtures in MgMOF-74 and ZnMOF-74 at 300 K, plotted as a function of the total pore concentration, $c_t = c_1 + c_2$.

For this purpose the dual-site Langmuir fits of the CBMC simulated pure component isotherms were employed; see Fig. 8a. In order to verify the accuracy of IAST calculations, we carried out CBMC simulations for CO₂/N₂, and CO₂/H₂ mixtures. Fig. 8b and c presents comparisons of the IAST predictions with the CBMC simulation data for component loadings. We see excellent agreement of the IAST with CBMC data; this confirms the applicability of IAST for practical purposes. Similar good agreement was earlier established for adsorption of a wide variety of mixtures in different zeolites [59,60].

In view of the fact that the accuracy of the CBMC simulations of the adsorption isotherms, especially of CO₂, are essential to the prediction of the permeation characteristics, we compare CBMC simulations of CO₂ isotherms at 300 K, and 313 K with the experimental results of Dietzel et al. [18] and Herm et al. [23]; see Fig. 9a and b. For both temperatures, the experimental data show an inflection at a loading of about 8 mol/kg that corresponds to one molecule of 1 CO₂ per atom of Mg in the framework. We note that for pressures exceeding 0.1 MPa, there is excellent agreement between the CBMC simulations and experiment. However, for pressures below 0.1 MPa, the experimental loadings are significantly higher than the simulated values. The reason for this deviation at pressures below 0.1 MPa can be traced to the fact that our force field implementation, following the work of Yazaydin et al. [20] does not explicitly account for orbital interactions and polarization. Such effects are particularly strong in the low pressure region; the influence of polarization is of lesser importance at higher pressures. Put another way, we would expect the CBMC simulation results to be reasonably good for high pressure membrane operation with upstream pressures that are considerably in excess of 0.1 MPa.

When the guest species do not have orbital interactions with the Mg atoms, the predictions of the CBMC simulations can be expected to be good. Fig. 9c compares CBMC simulations of the adsorption isotherms for CH₄ at 300 K with the experimental data of Dietzel et al. [18]; the agreement between the two sets is good for the entire pressure range.

An important consequence of the under-prediction of the CO₂ loadings for pressures below 0.1 MPa, is that the estimations of the adsorption selectivity from CBMC simulations can be expected to be conservative. To confirm this expectation, we compare the calculations of the adsorption selectivity for separation of 20/80 CO₂/H₂ mixtures using MgMOF-74 in the experimental study of Herm et al. [23] with corresponding CBMC mixture simulations; see Fig. 9d. We note that the CBMC mixture simulations predict somewhat lower values of S_{ads} over the range of pressures 0.1–10 MPa. On the basis of the results presented in Fig. 9d we conclude that the permeation selectivity results to be presented below are likely to be conservative.

At steady-state, the fluxes N_i obey

$$\frac{\partial N_i}{\partial z} = 0 \quad (12)$$

The steady state permeation fluxes N_i were determined by solving the set of Eqs. (1), (11) and (12) subject to the boundary conditions

$$\begin{aligned} z = 0; & \quad p_i = p_{i0}; \quad c_i = c_{i0} \\ z = \ell; & \quad p_i = p_{i\ell}; \quad c_i = c_{i\ell} \end{aligned} \quad (13)$$

using the numerical procedures described in our earlier works [54,61,62].

The M - S diffusivities \mathcal{D}_i were taken from the data presented in Fig. 1a. The binary exchange coefficient \mathcal{D}_{12} , in the mixtures were estimated on the basis of the pure component self-exchange coefficients, \mathcal{D}_{ii} , using the interpolation formula [34,36]

$$\mathcal{D}_{12} = (\mathcal{D}_{11})^{x_1} (\mathcal{D}_{22})^{x_2} \quad (14)$$

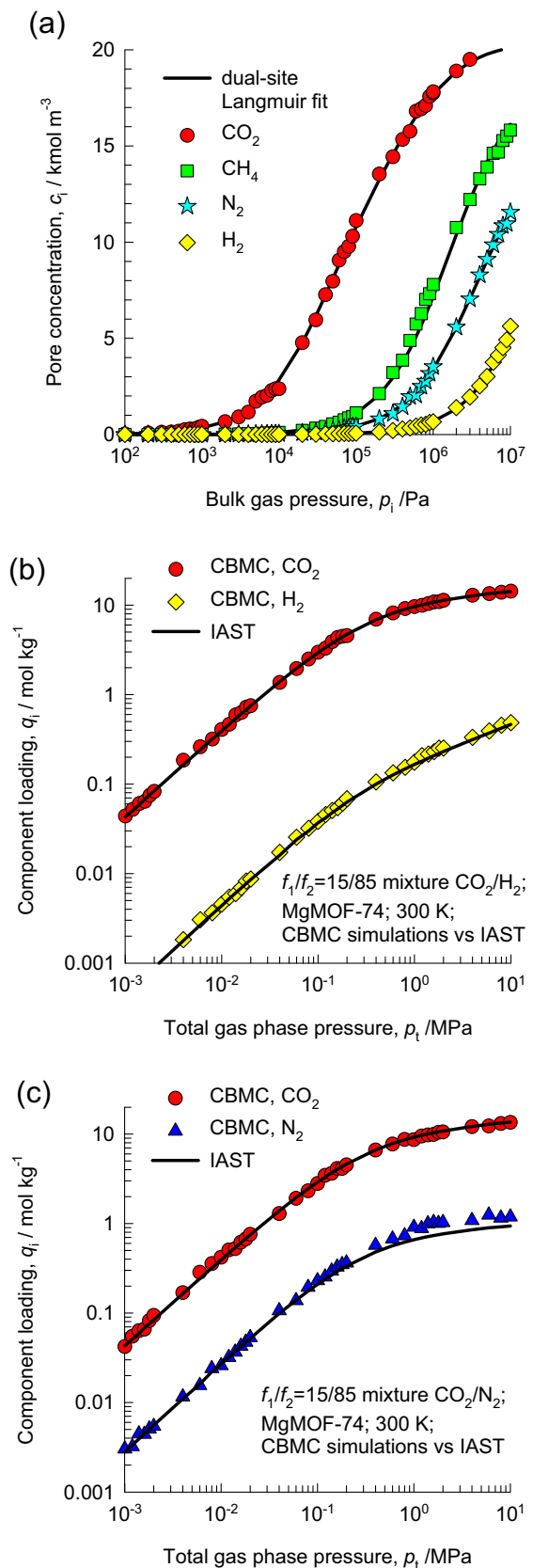


Fig. 8. (a) CBMC simulation data on the pure component adsorption isotherms of H₂, N₂, CH₄, and CO₂ in MgMOF-74 at 300 K. The continuous solid lines are the dual-Langmuir fits of the isotherms; the fit parameters are provided in Supplementary material accompanying this publication. (b, c) CBMC simulations for component loadings in CO₂/N₂, and CO₂/H₂ mixtures at 300 K with partial pressure $p_1/p_2 = 15/85$. The continuous solid lines are the IAST calculations using the dual-Langmuir fits of the CBMC simulated pure component isotherms.

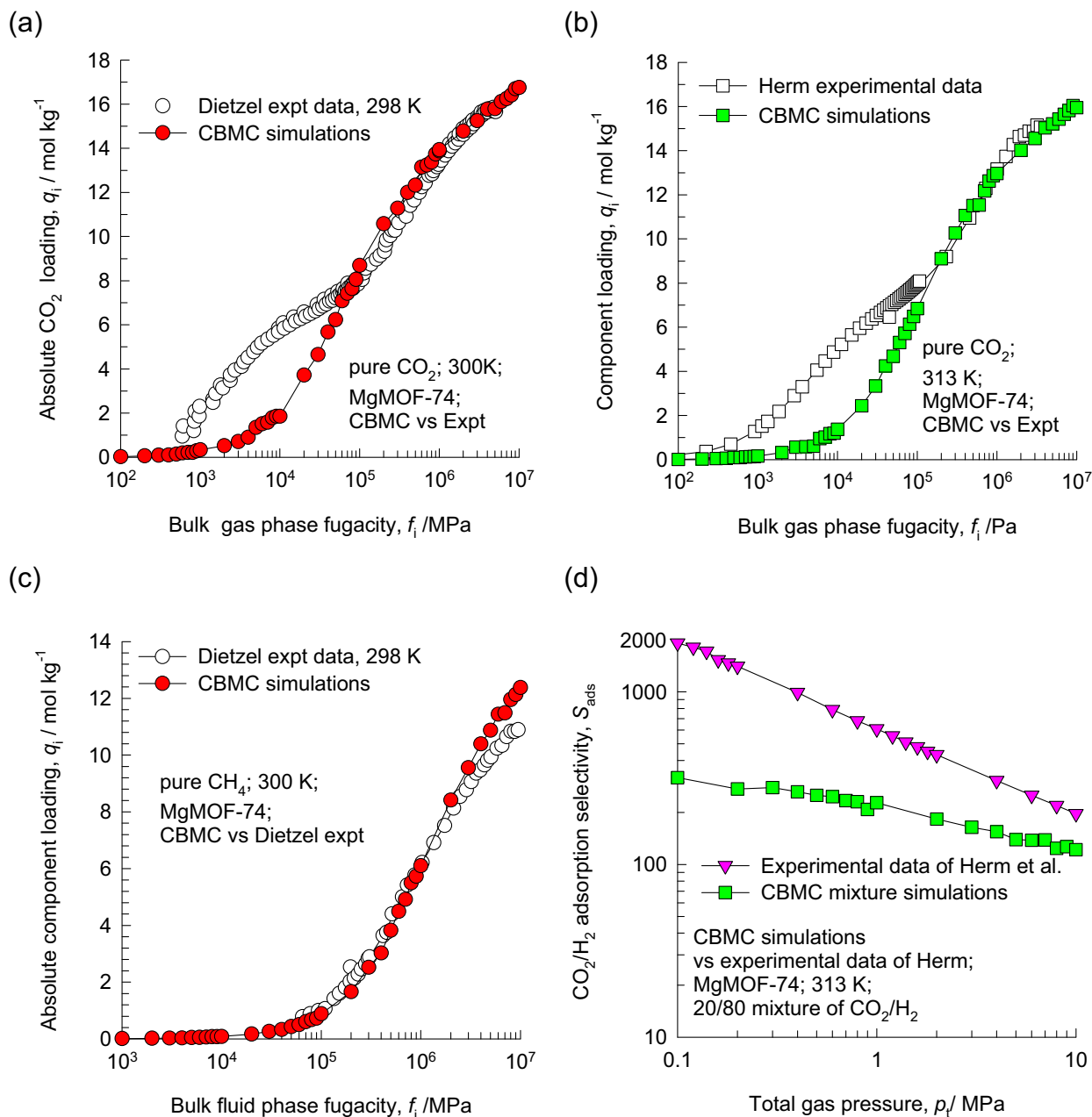


Fig. 9. (a, b) Comparison of the CBMC simulations of the pure component adsorption isotherms for CO₂ in MgMOF-74 at (a) 300 K, and (b) 313 K with experimental data of Dietzel et al. [18] and Herm et al. [23]. The published experimental data of Dietzel et al. is given in terms of excess loadings; these have been converted to absolute loadings by estimating the fluid densities within the pores using the Peng–Robinson equation of state. The experimental data on pore volume, 0.63 cm³/g was used in this conversion. (c) Comparison of the CBMC simulations of the pure component adsorption isotherms for CH₄ in MgMOF-74 at 300 K with experimental data of Dietzel et al. [18]. (d) Adsorption selectivity, S_{ads} , for 20/80 CO₂/H₂ mixture separation with MgMOF-74 at 313 K. Comparison of the experimental data of Herm et al. [23] with CBMC mixture simulations.

The data on \mathcal{D}_{ii} were obtained by fitting the $\mathcal{D}_i/\mathcal{D}_{ii}$ data plotted in Fig. 5b. All the data inputs are specified in the Supplementary material accompanying this publication.

Fig. 10 shows the calculations of the permeation selectivity, S_{perm} defined by

$$S_{perm} = \frac{N_1/N_2}{p_{10}/p_{20}} \quad (15)$$

where p_{10} are the partial pressures in the upstream compartment. Also shown in Fig. 10 are IAST calculations of the adsorption selectivity, S_{ads}

$$S_{ads} = \frac{c_{10}/c_{20}}{p_{10}/p_{20}} \quad (16)$$

where c_{i0} are the adsorption phase concentrations at the upstream face.

For all four mixtures, S_{ads} either decreases with increasing upstream pressure (CO₂/H₂, and CO₂/CH₄), or remains practically constant (CO₂/N₂, and CH₄/H₂). It is remarkable, therefore, to note that S_{perm} is significantly increased with p_{10} for all mixtures. This increase is directly attributable to the increase in the degree of correlations with increased loadings within the membrane. Put another way, with increased upstream pressure, the slowing-down effect is enhanced, resulting in a higher permeation selectivity. To underline this, Fig. 10 also presents the diffusion selectivity values, back-calculated using $S_{diff} = S_{perm}/S_{ads}$. The slowing-down effect is particularly dramatic for CO₂/H₂ mixtures. Here the S_{diff} increases from 0.04 to a value approaching 0.7 at $p_{10} > 1$ MPa. For the other

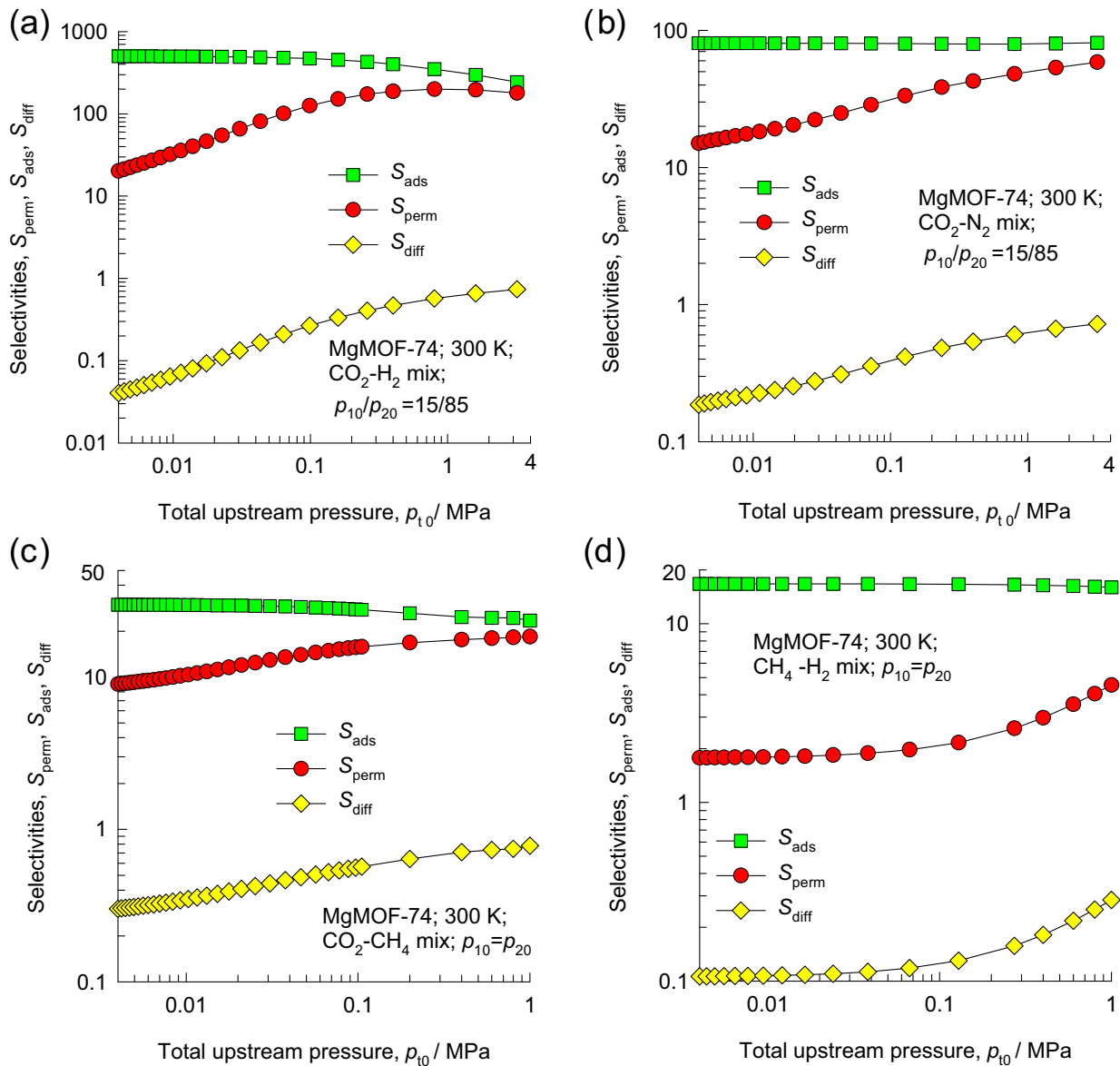


Fig. 10. Calculations for membrane permeation selectivity, S_{perm} , adsorption selectivity, S_{ads} , and diffusion selectivity, S_{diff} , for (a) CO_2/H_2 , (b) CO_2/N_2 , (c) CO_2/CH_4 , and (d) CH_4/H_2 mixtures across a MgMOF-74 membrane at 300 K, plotted as a function of the total upstream pressure, $p_{t0} = p_{10} + p_{20}$. For (a) and (b) the partial pressures in the gas phase in the upstream compartment satisfy $p_{10}/p_{20} = 15/85$. In (c) and (d) $p_{10} = p_{20}$. The downstream partial pressures were maintained at 100 Pa for each species in all cases. The fractional porosity of the crystalline layer, $\phi = 0.708$. The thickness of MgMOF-74 crystalline layer, $l = 50 \mu\text{m}$. The data on diffusivities, correlations, and dual-site Langmuir parameters for fitting of the pure component isotherms are specified in the Supplementary material accompanying this publication. The diffusion selectivities, S_{diff} , are obtained by dividing the S_{perm} values by S_{ads} .

three mixtures the increase in S_{diff} with increasing p_{t0} , though significant, is less dramatic.

The S_{diff} values for all mixtures presented in Fig. 10 are seen to be higher than the corresponding values presented in Figs. 6 and 7. The reasoning is that the data in Figs. 6 and 7 are for equimolar ($c_1 = c_2$) mixtures, whereas the concentrations within the membrane are such that $c_1 \gg c_2$, i.e. containing a preponderantly higher proportion of the tardier species 1. To illustrate this, Fig. 11 presents the profiles of pore concentrations of CO_2 and H_2 within the $50 \mu\text{m}$ thick MgMOF-74 layer for permeation of CO_2/H_2 mixture with a total upstream pressure of 1 MPa. We note that H_2 amounts to less than 2% of the total mixture concentration. This small proportion of the more mobile H_2 gets severely retarded in its motion within the channels.

The S_{perm} values of 200 estimated for CO_2/H_2 mixture at $p_{t0} > 1$ MPa, are more than an order of magnitude higher than those obtained in the experiments of Li et al. [63,64] for SAPO-34 mem-

branes. For CO_2/N_2 mixtures, the estimated values of $S_{perm} > 40$ at $p_{t0} > 1$ MPa, are comparable to those reported by White et al. [65] for NaY membranes, and about a factor two higher than those obtained for SAPO-34 [63,64] and DDR [66] membranes.

There is a severe drawback to the use of SAPO-34 and DDR membranes that arises because of the severe constraint at the window regions; this restriction leads to significantly lower membrane permeances, defined by

$$\Pi_i \equiv \frac{N_i}{p_{i0} - p_{il}} \quad (17)$$

Fig. 12 compares the permeances of CO_2 for CO_2/H_2 mixture permeation across MgMOF-74 and SAPO-34 membranes at 300 K. The Π_i values of SAPO-34 are more than two orders of magnitude lower because the severe constraint for inter-cage jumps across the $0.38 \text{ nm} \times 0.42 \text{ nm}$ sized windows of SAPO-34. When the window

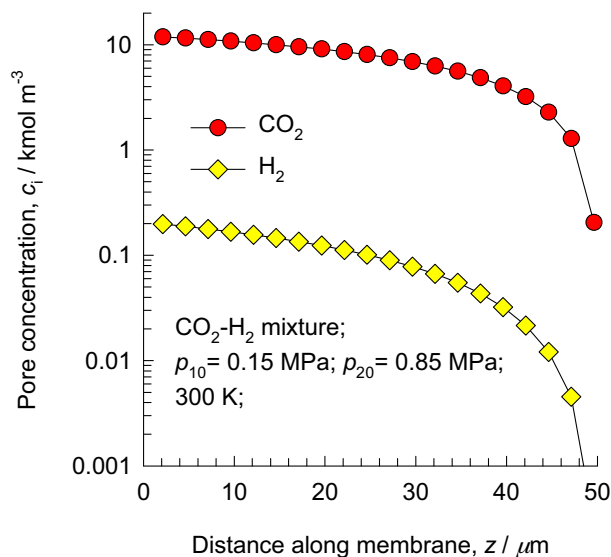


Fig. 11. Steady-state profiles of pore concentrations of CO_2 and H_2 within the 50 μm thick MgMOF-74 layer for permeation of CO_2/H_2 mixture. The upstream partial pressures are $p_{10} = 0.15$ MPa, $p_{20} = 0.85$ MPa.

size is still smaller, such as for ZIF-7 with 0.34 nm apertures, the mixture permeation is H_2 -selective [67].

For CO_2/CH_4 mixtures the S_{perm} value of 20 at $p_{t0} = 1$ MPa, is significantly higher than those reported in the experiments of Bux et al. [32] for ZIF-8; however, this value is significantly lower than the reported values of $S_{perm} \approx 50$ for SAPO-34 [63,64] and $S_{perm} = 300 - 500$ for DDR [66]. The narrow windows separating the cages of SAPO-34 and DDR membranes cause CH_4 to be severely constrained at the window regions, leading to S_{diff} values exceeding 100 [38,41]. In other words, the higher S_{perm} values for CO_2/CH_4 with SAPO-34 and DDR membranes are driven by diffusion selectivity.

The severe constraint experienced by CH_4 at the window regions of DDR, SAPO-34, ZIF-7, and ZIF-8 causes CH_4/H_2 mixture permeation to be H_2 -selective [13,33,63,64,67,68].

The higher permeance of MgMOF-74, coupled with its significantly high permeation selectivity are both attractive features

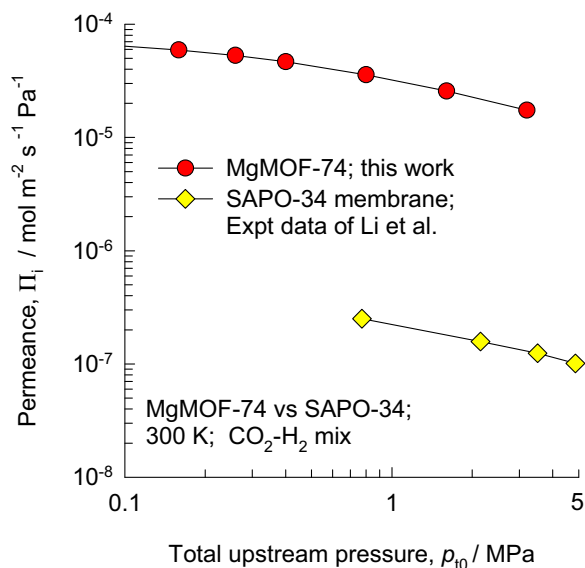


Fig. 12. Comparison of the permeances of CO_2 for CO_2/H_2 mixture permeation across MgMOF-74 and SAPO-34 membranes at 300 K.

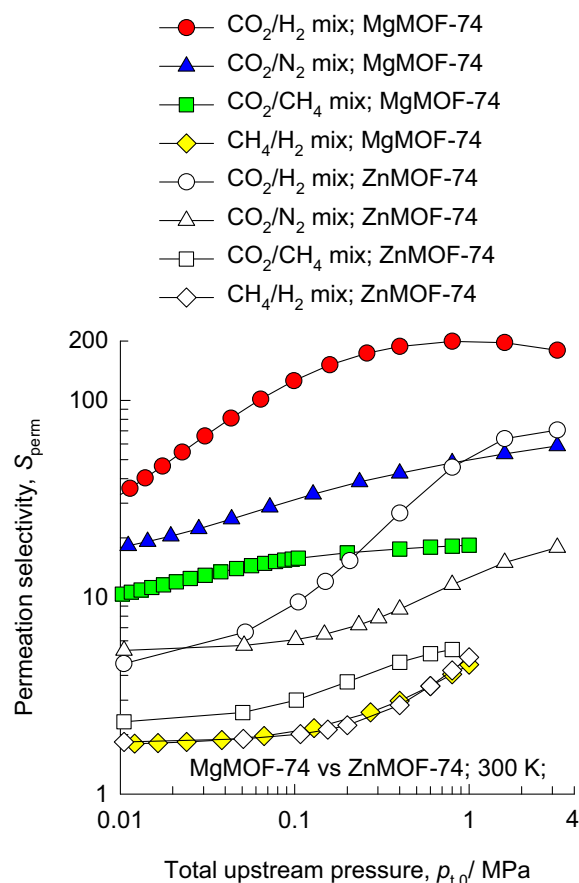


Fig. 13. Comparison of the permeation selectivity, S_{perm} , for CO_2/H_2 , CO_2/N_2 , and CH_4/H_2 mixtures across MgMOF-74 and ZnMOF-74 membranes. The upstream conditions are as in the legend to Fig. 10.

that warrant experimental verification for further technological exploitation, especially for CO_2/H_2 , and CO_2/N_2 mixtures.

Finally, it is instructive to compare the selectivities for mixture permeation across MgMOF-74 and ZnMOF-74 membranes; see data in Fig. 13. For CH_4/H_2 mixtures, the permeation selectivities across MgMOF-74 and ZnMOF-74 membranes are practically the same; this is because the adsorption strengths of the CH_4 are practically the same in these structures, as is evidenced in similar values of the isosteric heats of adsorption; cf. Fig. 3a and b. Indeed, the obtained S_{perm} values in the range of 2–4 for CH_4/H_2 mixtures were also obtained by Keskin [25] with CoMOF-74 and NiMOF-74 membranes. For CO_2/H_2 , CO_2/N_2 , and CO_2/CH_4 mixtures, the S_{perm} values for MgMOF-74 are significantly higher than those for ZnMOF-74; the reason can be traced to the significantly higher binding energy of CO_2 in MgMOF-74. The higher diffusion selectivities of ZnMOF-74, witnessed in Figs. 6 and 7, cannot compensate for the significantly lower S_{ads} . MgMOF-74 is the membrane of choice for CO_2 -selective permeation.

5. Conclusions

Molecular simulations have been used to investigate the characteristics of adsorption, and diffusion characteristics of a variety of guest molecules within the 1.1 nm channels of MgMOF-74. The following major conclusions can be drawn from this study.

With increasing binding energy of the guest molecules, the zero-loading diffusivity $D_i(0)$ falls increasingly below that predicted by the Knudsen formula, Eq. (6). The validity of Eq. (6) is restricted to

cases where the binding energy of the molecule is negligibly small, as is the case with H₂.

The ratio $D_i(0)/D_{i,Kn}$ for a given molecule correlates with the isosteric heat of adsorption; see Fig. 4. This correlation provides a good starting point for estimation of diffusivities.

The degree of correlations in molecular motions, D_i/D_{ii} , increases as the pore concentration increases; see Fig. 5.

For binary mixtures, the diffusion selectivity S_{diff} is lower than the Knudsen formula when species 1 has the higher binding energy; see Fig. 5.

For binary mixtures in which species 2 is more mobile, S_{diff} increases with increasing pore concentration due to increased correlation effects; see Figs. 6 and 7.

Adsorption and diffusion do not go hand in hand; the membrane permeation selectivity for binary mixtures S_{perm} is significantly lower than S_{ads} ; see Fig. 10.

For permeation of CO₂/H₂, CO₂/N₂, CO₂/CH₄, and CH₄/H₂ mixtures across a MgMOF-74 membrane, slowing-down effects become increasingly important at high upstream pressures, and consequently the S_{perm} is significantly enhanced, and tends to approach S_{ads} at $p_{t0} > 1$ MPa.

CO₂/N₂ permeation selectivities with MgMOF-74 membranes at $p_{t0} > 1$ MPa are about a factor two higher than those reported for SAPO-34 and DDR membranes.

Correlation effects are particularly important for CO₂/H₂ permeation, S_{perm} values of 200 at $p_{t0} > 1$ MPa are more than an order of magnitude higher than those reported for SAPO-34 membrane.

An important advantage of MgMOF-74 membranes is that due to the 1.1 nm channel sizes, the permeances are more than two orders of magnitude higher than for SAPO-34 and DDR membranes.

Another message that emerges from this study is that Maxwell–Stefan modeling of membrane permeation fluxes must take proper account of diffusion correlations. For more accurate modeling of membrane permeation, there may be a need to also account for inter-crystalline and grain boundary resistances in the polycrystalline membrane layers.

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Appendix A. Supplementary data

Supplementary data associated with this article can be found, in the online version, at doi:10.1016/j.memsci.2011.05.001.

Nomenclature

c_i	pore concentration of species i , mol m ⁻³
$c_{i,sat}$	saturation capacity of species i , mol m ⁻³
c_t	total concentration in mixture, mol m ⁻³
d_p	pore diameter, m
$D_{i,self}$	self-diffusivity of species i , m ² s ⁻¹
\bar{D}_i	Maxwell–Stefan diffusivity, m ² s ⁻¹

$\bar{D}_i(0)$	zero-loading M – S diffusivity, m ² s ⁻¹
$D_{i,Kn}$	Knudsen diffusivity of species i , m ² s ⁻¹
\bar{D}_{ii}	self-exchange coefficient, m ² s ⁻¹
\bar{D}_{12}	M – S exchange coefficient for binary mixture, m ² s ⁻¹
$-\Delta H_{St}$	isosteric heat of adsorption, J mol ⁻¹
l	thickness of membrane, m
M_i	molar mass of species i , kg mol ⁻¹
n	number of components in mixture
N_i	molar flux of species i defined in terms of the membrane area, mol m ⁻² s ⁻¹
p_{i0}	partial pressure of species i in upstream compartment, Pa
p_{il}	partial pressure of species i in downstream compartment, Pa
p_t	total system pressure in upstream compartment, Pa
q_i	molar loading species i
R	gas constant, 8.314 J mol ⁻¹ K ⁻¹
S_{ads}	adsorption selectivity
S_{diff}	diffusion selectivity
S_{perm}	permeation selectivity
T	temperature, K
V_p	accessible pore volume, m ³ kg ⁻¹
x_i	mole fraction of species i based on loading within pore
z	distance along the membrane, m

Greek letters

Γ_{ij}	thermodynamic factors, dimensionless
ϕ	fractional pore volume, dimensionless
μ_i	molar chemical potential, J mol ⁻¹
Π_i	permeance of species i , mol m ⁻² s ⁻¹ Pa ⁻¹
θ_i	fractional occupancy of component i , dimensionless

Subscripts

0	referring to upstream face of membrane
l	referring to downstream face of membrane
i	referring to component i
p	referring to pore
t	referring to total mixture
Kn	referring to Knudsen

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