

*Supplementary Material (1) to accompany **Investigating the Influence of Diffusional Coupling on Mixture Permeation across Porous Membranes***

# Experimental Membrane Permeation Data; Comparisons with Maxwell-Stefan Simulations with Simulation Data Inputs

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## 1. Abstract

Explicit analytic expressions for the permeation fluxes of binary mixtures across porous membranes are derived.

## 2. Introduction

In this first of two Supplementary Materials to accompany our article *Investigating the Influence of Diffusional Coupling on Mixture Permeation across Porous Membranes*, we derive analytic expressions for the permeation fluxes across porous membranes using the Maxwell-Stefan diffusion formulation.

The Figures accompanying this Appendix A contain the experimental membrane permeation data, along with inputs (pure component isotherm fits, Maxwell-Stefan data on diffusivities  $D_1$ ,  $D_2$ , and  $D_{12}$  used in the simulations to match the binary mixture permeation experiments.

## 3. Maxwell-Stefan equations for mixture diffusion

The Maxwell-Stefan (M-S) equations for  $n$ -component diffusion inside porous crystalline materials can be written as

$$-\phi \frac{c_i}{RT} \nabla \mu_i = \sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_j N_i - x_i N_j}{D_{ij}} + \frac{N_i}{D_i}; \quad i = 1, 2, \dots, n \quad (1)$$

In equation (1) the  $D_i$  is the M-S diffusivities of species  $i$ , portraying the interaction between component  $i$  in the mixture with the surface, or wall of the structure. The fluxes  $N_i$  defined in equation (1) are expressed in terms of the number of moles of species  $i$  transported per  $\text{m}^2$  of *crystalline material* per second. The  $D_{ij}$  are M-S exchange coefficients representing interaction between component  $i$  with component  $j$ . The  $c_i$  are the loadings, defined in terms of moles per  $\text{m}^3$  of *accessible* pore volume, within the pore.

The pore concentrations  $c_i$  are related to the molar loadings  $q_i$  by the expression

$$c_i = \frac{\rho q_i}{\phi} = \frac{q_i}{V_p} \quad (2)$$

where  $\rho$  represents the framework density,  $\phi$  represents the fractional pore volume, and  $V_p$  is the  $\text{m}^3$  accessible pore volume per kg of framework. It is to be noted that the  $\phi$  appears in the M-S equations (1) because the  $c_i$  are defined in terms of pore volume and not the total volume of the crystals. The  $c_i$  are useful measures when comparing different materials. We will have occasion to use both  $c_i$  and the molar loadings  $q_i$ .

An equivalent formulation of the M-S equations uses the molar loadings  $q_i$  as concentration measures

$$-\rho \frac{q_i}{RT} \nabla \mu_i = \sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_j N_i - x_i N_j}{D_{ij}} + \frac{N_i}{D_i}; \quad i = 1, 2, \dots, n \quad (3)$$

in which the M-S diffusivities,  $D_i$  and  $D_{ij}$ , have the same significance, and magnitudes as in equations (1).

The  $x_i$  in equations (1) and (3) represent the component mole fractions in the adsorbed phase within the pores

$$x_i = q_i / q_t = c_i / c_t; \quad 1, 2, \dots, n \quad (4)$$

where  $q_t$  and  $c_t$  are the *total* mixture loadings and pore concentrations, respectively:

$$q_t = \sum_{i=1}^n q_i; \quad c_t = \sum_{i=1}^n c_i;$$

The Onsager reciprocal relations require

$$D_{ij} = D_{ji} \quad (5)$$

By defining an  $n$ -dimensional square matrix  $[B]$  with elements

$$B_{ii} = \frac{1}{D_i} + \sum_{\substack{j=1 \\ j \neq i}}^n \frac{x_j}{D_{ij}}; \quad B_{ij} = -\frac{x_i}{D_{ij}}; \quad i, j = 1, 2, \dots, n \quad (6)$$

We can recast eq. (1) into the following form



$$-\phi \frac{c_i}{RT} \nabla \mu_i = -\rho \frac{q_i}{RT} \nabla \mu_i = \sum_{j=1}^n B_{ij} N_j; \quad i=1,2,..n \quad (7)$$

The chemical potential gradients  $\nabla \mu_i$  can be related to the gradients in the pore concentrations,  $c_i$ , and the gradients of the molar loadings,  $q_i$ , by defining thermodynamic correction factors  $\Gamma_{ij}$

$$\frac{q_i}{RT} \nabla \mu_i = \sum_{j=1}^n \Gamma_{ij} \nabla q_j; \quad \frac{c_i}{RT} \nabla \mu_i = \sum_{j=1}^n \Gamma_{ij} \nabla c_j; \quad \Gamma_{ij} = \frac{q_i}{p_i} \frac{\partial p_i}{\partial q_j} = \frac{c_i}{p_i} \frac{\partial p_i}{\partial c_j}; \quad i, j = 1, \dots, n \quad (8)$$

In the general case, the elements of the matrix of thermodynamic correction factors  $\Gamma_{ij}$  must be determined by numerical differentiation of the equations describing adsorption equilibrium. For description of adsorption equilibrium, the Ideal Adsorbed Solution Theory (IAST) of Myers and Prausnitz [1] is of sufficient accuracy and is used consistently in this work.

Equation (7) can be re-written in  $n$ -dimensional matrix notation as

$$-\phi [\Gamma] \nabla(c) = -\rho [\Gamma] \nabla(q) = [B](N) \quad (9)$$

In proceeding further with our development we denote the inverse of  $[B]$  as  $[\Delta]$ :

$$[B]^{-1} \equiv [\Delta] \quad (10)$$

For the special case of a binary mixture,  $n=2$ , the four elements of the matrix  $[B]$  are given explicitly as follows

$$B_{11} = \frac{1}{D_1} + \frac{x_2}{D_{12}} \quad (11)$$

$$B_{22} = \frac{1}{D_2} + \frac{x_1}{D_{12}} \quad (12)$$

$$B_{12} = -\frac{x_1}{D_{12}} \quad (13)$$

$$B_{21} = -\frac{x_2}{D_{12}} \quad (14)$$

The elements of the matrix  $[\Delta]$  can be derived from the inverse of  $[B]$

$$[\Delta] \equiv [B]^{-1} = \begin{bmatrix} \frac{1}{D_1} + \frac{x_2}{D_{12}} & -\frac{x_1}{D_{12}} \\ -\frac{x_2}{D_{12}} & \frac{1}{D_2} + \frac{x_1}{D_{12}} \end{bmatrix}^{-1} \quad (15)$$

The matrix inversion can be carried out explicitly to give the following expressions

$$\Delta_{11} = D_1 \frac{\left(1 + \frac{x_1 D_2}{D_{12}}\right)}{1 + \frac{x_1 D_2}{D_{12}} + \frac{x_2 D_1}{D_{12}}} \quad (16)$$

$$\Delta_{22} = D_2 \frac{\left(1 + \frac{x_2 D_1}{D_{12}}\right)}{1 + \frac{x_1 D_2}{D_{12}} + \frac{x_2 D_1}{D_{12}}} \quad (17)$$

$$\Delta_{12} = D_1 \frac{\frac{x_1 D_2}{D_{12}}}{1 + \frac{x_1 D_2}{D_{12}} + \frac{x_2 D_1}{D_{12}}} \quad (18)$$

$$\Delta_{21} = D_2 \frac{\frac{x_2 D_1}{D_{12}}}{1 + \frac{x_1 D_2}{D_{12}} + \frac{x_2 D_1}{D_{12}}} = \frac{x_2}{x_1} \Delta_{12} \quad (19)$$

In the limiting case of  $x_1 \rightarrow 1$ ;  $x_2 \rightarrow 0$  we obtain

$$\Delta_{11} = \frac{D_1 \left(1 + \frac{x_1 D_2}{D_{12}}\right)}{1 + \frac{x_1 D_2}{D_{12}} + \frac{x_2 D_1}{D_{12}}} \rightarrow D_1; \quad x_1 \rightarrow 1; \quad x_2 \rightarrow 0 \quad (20)$$

In the limiting case of  $x_2 \rightarrow 1$ ;  $x_1 \rightarrow 0$  we obtain

$$\Delta_{22} = \frac{D_2 \left( 1 + \frac{x_2 D_1}{D_{12}} \right)}{1 + \frac{x_2 D_1 + x_1 D_2}{D_{12}}} \rightarrow D_2; \quad x_2 \rightarrow 1; \quad x_1 \rightarrow 0 \quad (21)$$

#### 4. Analytic expressions for the thermodynamic correction factors $\Gamma_{ij}$

It is useful to develop simple analytic expressions for the thermodynamic correction factors. Analytic expressions for  $\Gamma_{ij}$  can be derived for the case in which the component loadings follow the mixed-gas

Langmuir isotherm

$$\theta_i = \frac{q_i}{q_{i,sat}} = \frac{c_i}{c_{i,sat}} = \frac{b_i p_i}{1 + \sum_{i=1}^n b_i p_i}; \quad i = 1, 2, \dots, n \quad (22)$$

Carrying out the differentiation in eq. (8) the elements of  $[\Gamma]$  are

$$\Gamma_{ij} = \frac{q_{i,sat}}{q_{j,sat}} \left( \delta_{ij} + \frac{\theta_i}{\theta_V} \right); \quad i, j = 1, \dots, n \quad (23)$$

where  $\delta_{ij}$  is the Kronecker delta. The fractional vacancy  $\theta_V$  is defined by:

$$\theta_V = (1 - \theta_1 - \theta_2 - \dots - \theta_n) \quad (24)$$

The mixed-gas Langmuir model (22) is thermodynamically inconsistent, except when the saturation adsorption capacities of all component species are equal [2]; in this case, the elements of  $[\Gamma]$  for binary mixtures further simplifies to yield

$$[\Gamma] = \frac{1}{1 - \theta_1 - \theta_2} \begin{bmatrix} 1 - \theta_2 & \theta_1 \\ \theta_2 & 1 - \theta_1 \end{bmatrix} \quad (25)$$

#### 5. Unary permeation fluxes and permeances

For pure component adsorption, the thermodynamic correction factor is

$$\Gamma_i = \frac{1}{1-\theta_i} \quad (26)$$

For the special case of unary diffusion of component  $i$  equation (3) simplifies to

$$N_i = -\rho D_i \Gamma_i \frac{dq_i}{dz} \quad (27)$$

The Maxwell-Stefan diffusivity,  $D_i$  is generally a function of the occupancy or loading. For a given guest molecule, the loading dependence is influenced by a variety of factors including pore size, pore topology and connectivity [3-6]. For structures in which correlation effects are particularly strong, such as MFI, MgMOF-74, NiMOF-74, FAU, the M-S diffusivities,  $D_i$ , often show a linear dependence on the number of vacant positions available, i.e.

$$D_i(\theta) = D_i \theta_V \quad (28)$$

For this scenario, combining equation (26), (27), and (28) yields

$$N_i = -\rho D_i \frac{dq_i}{dz} \quad (29)$$

The absence of the thermodynamic correction factor is remarkable; this is a result of the fact that the “correction” for the loading dependence of the M-S diffusivities cancels out with the “thermodynamic correction factor”. We will see below that such cancellation effects also manifest for mixture diffusion.

This simple expression (29) can be integrated across a porous crystalline layer of thickness  $\delta$  to obtain the following expression for the steady-state unary permeation fluxes

$$N_i = \frac{\rho D_i}{\delta} \Delta q_i \quad (30)$$

The unary permeance,  $\Pi_i^0$ , is obtained by dividing by the trans-membrane pressure difference

$$\Pi_i^0 \equiv \frac{N_i}{\Delta p_i} \quad (31)$$

In this work equations (30), and (31) are used to back out the values of the membrane transport coefficients  $\frac{\rho D_1}{\delta}$ , and  $\frac{\rho D_2}{\delta}$  from experimental data on unary permeation fluxes and permeances. These values of  $\frac{\rho D_1}{\delta}$ , and  $\frac{\rho D_2}{\delta}$  are then used for estimation of the permeances for binary mixtures, with additional assumptions regarding the values of the degrees of correlation.

## 6. Maxwell-Stefan modeling of binary mixture permeation

We now derive the appropriate set of equations to describe binary mixture permeation at *steady state* across a porous crystalline layer of thickness  $\delta$ . In our theoretical development, the resistance of the support layer is ignored. Furthermore, the micro-porous crystalline layer is considered to be defect-free, and inter-grain boundary resistances have been ignored.

At steady-state, the molar fluxes  $N_i$  obey

$$\frac{\partial N_i}{\partial z} = 0 \quad (32)$$

The steady state permeation fluxes  $N_i$  are obtained determined by solving the set of two coupled differential equations (9), subject to the boundary conditions

$$\begin{aligned} z = 0; & \quad p_i = p_{i0}; \quad c_i = c_{i0}; \quad q_i = q_{i0} \\ z = \delta; & \quad p_i = p_{i\delta}; \quad c_i = c_{i\delta}; \quad q_i = q_{i\delta} \end{aligned} \quad (33)$$

It is convenient to define the driving forces for trans-membrane transport

$$\Delta q_i = q_{i0} - q_{i\delta}; \quad \Delta c_i = c_{i0} - c_{i\delta}; \quad i = 1,2 \quad (34)$$

In our earlier works we have solved the set of equations (9), (33) and (34) using rigorous numerical porcedures that takes account of the loading dependence of the M-S diffusivities,  $D_1$ , and  $D_2$  [7-9].

Exact *analytic* solutions to the set of equations (9), (33) and (34) are possible when equation (28) for the loading dependence is invoked, along with the expression (25) for the thermodynamic correction factors [10]. The obtained expressions, corresponding to equation (44) of Krishna and Baur [10] are

$$N_1 = \frac{\rho D_1}{\delta} \left[ \frac{\left(1 + \frac{x_1 D_2}{D_{12}}\right) \Delta q_1 + \frac{x_1 D_2}{D_{12}} \Delta q_2}{1 + \frac{x_1 D_2}{D_{12}} + \frac{x_2 D_1}{D_{12}}} \right] \quad (35)$$

and

$$N_2 = \frac{\rho D_2}{\delta} \left[ \frac{\frac{x_2 D_1}{D_{12}} \Delta q_1 + \left(1 + \frac{x_2 D_1}{D_{12}}\right) \Delta q_2}{1 + \frac{x_1 D_2}{D_{12}} + \frac{x_2 D_1}{D_{12}}} \right] \quad (36)$$

The remarkable feature of the exact analytic solution is that the “correction” factor for the loading dependence of the M-S diffusivities, effectively cancels out the “thermodynamic correction factor”. Indeed this “canceling out” effect is precisely analogous to that witnessed for unary permeation in equation (30). The accuracy of these simple results presented in equations (35), and (36) have been established by detailed comparisons with numerical solutions in our earlier work [10]. While such cancelling out effects are the *combined result* of two *special* scenarios, described by equations (25) and (28). The final obtained equations are of more general applicability and not restricted to these chosen special scenarios; this have been verified by comparing with rigorous numerical solutions [11, 12] of the equations taking all factors into consideration.

Another important point to note is that though the derivation of the equations (35), and (36) uses the mixed gas Langmuir isotherm to obtain simplified analytic solutions, the pore loadings in equations (35), and (36) are determined with the rigorous IAST calculations of thermodynamic equilibrium in all cases considered in our paper.

Strictly speaking equations (35), and (36) are “exact” only for the case where the saturation capacities of the constituent species are identical; it provides a good approximation when the saturation capacities are different.

The component permeances in the binary mixture can be obtained from

$$\Pi_i \equiv \frac{N_i}{\Delta p_i}; \quad i = 1,2 \quad (37)$$

In the general case, for each component, the permeance determined from equation (37) is different from that determined from unary permeation experiments. These differences are largely traceable to the degrees of correlation. The permeation selectivity is defined as

$$S_{perm} = \frac{\Pi_1}{\Pi_2} = \frac{N_1/\Delta p_1}{N_2/\Delta p_2} \quad (38)$$

It is convenient to define the adsorption selectivity,  $S_{ads}$ , in terms of the component loadings and partial pressures in the upstream compartment.

$$S_{ads} = \frac{q_{10}/q_{20}}{p_{10}/p_{20}} \quad (39)$$

In further discussions it is useful to define the diffusion selectivity of the membrane,  $S_{diff}$ :

$$S_{diff} = \frac{S_{perm}}{S_{ads}} \quad (40)$$

If the downstream partial pressures  $p_{1\delta}$ , and  $p_{2\delta}$ , and component loadings  $q_{1\delta}$ , and  $q_{2\delta}$  are negligible in comparison with the corresponding upstream values, the following expression  $S_{diff}$  can be derived using equations (35), and (36)

$$S_{diff} = \frac{D_1}{D_2} \left[ \frac{\left(1 + \frac{x_1 D_2}{D_{12}}\right) q_{10} + \frac{x_1 D_2}{D_{12}} q_{20}}{\frac{x_2 D_1}{D_{12}} q_{10} + \left(1 + \frac{x_2 D_1}{D_{12}}\right) q_{20}} \right] \frac{q_{20}}{q_{10}} = \frac{D_1}{D_2} \left[ \frac{\left(1 + \frac{x_1 D_2}{D_{12}}\right) x_1 + \frac{x_1 D_2}{D_{12}} x_2}{\frac{x_2 D_1}{D_{12}} x_1 + \left(1 + \frac{x_2 D_1}{D_{12}}\right) x_2} \right] \frac{x_2}{x_1} \quad (41)$$

which simplifies to yield

$$S_{diff} = \frac{D_1}{D_2} \left[ \frac{\left(1 + \frac{x_1 D_2}{D_{12}}\right) x_1 + \frac{x_1 D_2}{D_{12}} x_2}{\frac{x_2 D_1}{D_{12}} x_1 + \left(1 + \frac{x_2 D_1}{D_{12}}\right) x_2} \right] \frac{x_2}{x_1} = \frac{D_1}{D_2} \frac{\left(1 + \frac{D_2}{D_{12}}\right)}{\left(1 + \frac{D_1}{D_{12}}\right)} \quad (42)$$

where we have invoked the relation

$$x_1 + x_2 = 1 \quad (43)$$

A good approximation to equation (42) is

$$S_{diff} \approx \frac{D_{1,self}}{D_{2,self}} \quad (44)$$

which is the ratio of self-diffusivities in the *mixture*. The detailed proof of equation (44) is available in the Supplementary Material accompanying our previous publication. Self-diffusivities are significantly to determine from MD simulations, and equation (44) serves as an adequate measure of diffusion selectivities in mixture diffusion.

## 7. Methodology adopted for interpretation of mixture permeation experiments

For interpretation of the experimental data on binary mixture permeation the following step-wise procedure is followed.

(1) The first step is to estimate the membrane transport coefficients  $\frac{\rho D_1}{\delta}$ , and  $\frac{\rho D_2}{\delta}$  from

experimental data on unary permeation fluxes and permeances using equations (30), and (31). For these calculations the pure component molar loadings obtained from dual-Langmuir Freundlich fits of pure component isotherms in most cases. For SAPO-34 membranes, the fits are using the Statistical isotherm [13]. The fit parameters used for this purpose are specified in the attached Figures accompanying this Supporting Material.



(2) The estimated values of the membrane transport coefficients  $\frac{\rho D_1}{\delta}$ , and  $\frac{\rho D_2}{\delta}$  are used for calculation of the steady-state fluxes  $N_1$  and  $N_2$  for binary mixtures for a range of upstream pressures using equations (35) and (36). In these calculations, the component loadings at the upstream face  $q_i$  of the membrane are calculated using the Ideal Adsorbed Solution Theory (IAST) of Myers and Prausnitz [1]. The pure component isotherm fits are the same as those used in step (1) above. For the estimation of the steady-state fluxes and permeances, additional information is required on the *degrees of correlations*,  $\frac{D_1}{D_{12}}$ , and  $\frac{D_2}{D_{12}}$ . There is only one adjustable parameter in view of the fact that the ratio  $\frac{D_1}{D_2}$  is simply the ratio of the membrane transport coefficients. In all cases, the chosen value of the degree of correlation, say  $\frac{D_1}{D_{12}}$  is on the basis of MD simulations for the specific guest/host combination. In every case, the rationale for the choice of the value of  $\frac{D_1}{D_{12}}$  is provided in the accompanying Figures, on a case-by-case basis.

## 8. Notation

$b_i$	parameter in the pure component Langmuir adsorption isotherm, $\text{Pa}^{-1}$
$[B]$	matrix of inverse M-S coefficients, defined by eq. (6), $\text{m}^{-2} \text{s}$
$c_i$	pore concentration of species $i$ , $c_i = q_i/V_p$ , $\text{mol m}^{-3}$
$c_{i,\text{sat}}$	saturation capacity of species $i$ , $\text{mol m}^{-3}$
$c_t$	total pore concentration in mixture, $\text{mol m}^{-3}$
$D_i$	Fick diffusivity of species $i$ , $\text{m}^2 \text{s}^{-1}$
$\bar{D}_i$	M-S diffusivity of species $i$ , $\text{m}^2 \text{s}^{-1}$
$D_{ij}$	M-S exchange coefficient, $\text{m}^2 \text{s}^{-1}$
$D_{i,\text{self}}$	self-diffusivity of species $i$ , $\text{m}^2 \text{s}^{-1}$
$f_i$	partial fugacity of species $i$ , Pa
$p_i$	partial pressure of species $i$ in upstream compartment, Pa
$p_t$	total system pressure in upstream compartment, Pa
$q_i$	molar loading of species $i$ , $\text{mol kg}^{-1}$
$q_{i,\text{sat}}$	molar loading of species $i$ at saturation, $\text{mol kg}^{-1}$
$q_t$	total molar loading of mixture, $\text{mol kg}^{-1}$
$n$	number of components in mixture, dimensionless
$N_i$	molar flux of species $i$ defined in terms of the membrane area, $\text{mol m}^{-2} \text{s}^{-1}$
$R$	gas constant, $8.314 \text{ J mol}^{-1} \text{ K}^{-1}$
$S_{\text{ads}}$	adsorption selectivity, dimensionless
$S_{\text{diff}}$	diffusion selectivity, dimensionless
$S_{\text{perm}}$	permeation selectivity, dimensionless
$T$	absolute temperature, K
$V_p$	pore volume, $\text{m}^3 \text{ kg}^{-1}$
$x_i$	mole fraction of species $i$ based on loading within pore, dimensionless
$z$	distance coordinate, m

### ***Greek letters***

$\delta_{ij}$	Kronecker delta, dimensionless
$\delta$	thickness of membrane, m
$[\Delta]$	matrix defined of M-S diffusivities, $\text{m}^2 \text{s}^{-1}$
$\Delta_{ij}$	elements of $[\Delta]$ , $\text{m}^2 \text{s}^{-1}$
$[\Gamma]$	matrix of thermodynamic factors, dimensionless
$\Gamma_{ij}$	element of $[\Gamma]$ , dimensionless
$\phi$	fractional pore volume, dimensionless
$\Pi_i$	permeance of species $i$ in mixture, $\text{mol m}^{-2} \text{s}^{-1} \text{Pa}^{-1}$
$\Pi_i^0$	permeance of pure component $i$ , $\text{mol m}^{-2} \text{s}^{-1} \text{Pa}^{-1}$
$\mu_i$	molar chemical potential, $\text{J mol}^{-1}$
$\theta_i$	fractional occupancy of component $i$ , dimensionless
$\theta_V$	fractional vacancy, dimensionless
$\Theta_i$	loading of species $i$ , molecules per unit cell or per cage
$\Theta_{i,\text{sat}}$	saturation loading of species $i$ , molecules per unit cell or per cage
$\Theta_t$	total molar loading of mixture, molecules per unit cell or per cage
$\rho$	framework density, $\text{kg m}^{-3}$

### ***Subscripts***

$i$	referring to component $i$
$t$	referring to total mixture

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## **10. Listing of Figures containing membrane permeation data and analysis**

The structural details, pore landscapes, surface area vs pore size distributions, isotherm fit parameters, and analysis of experimental data on membrane permeation, are presented in the accompanying Figures in the following order:

### **Cage-type structures with narrow windows**

SAPO-34

DDR

ZIF-8

### **1D micro-porous channels**

NiMOF-74

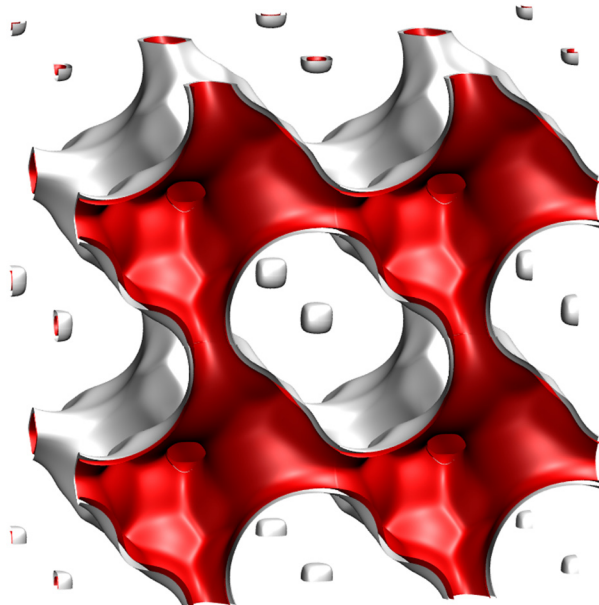
### **Intersecting channels**

MFI

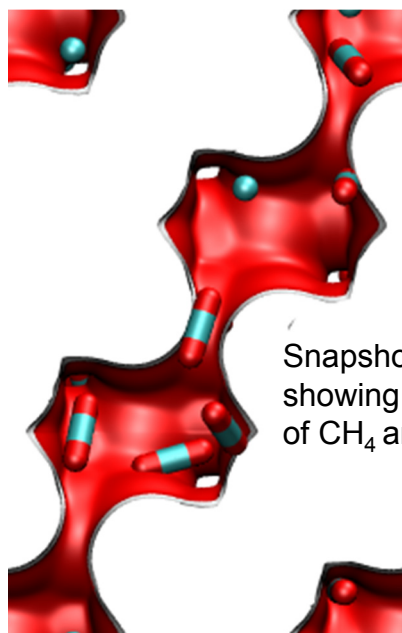
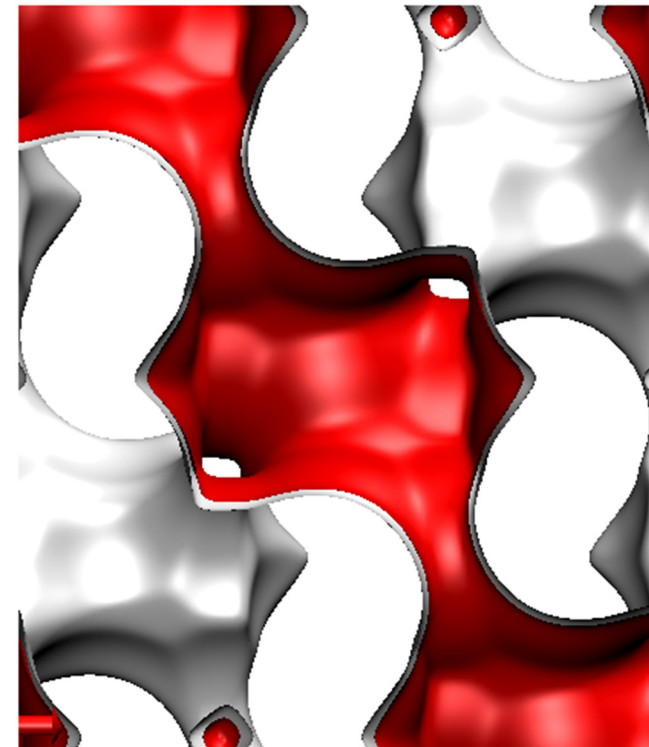
**SAPO-34**  
**(structural analog of CHA)**

# CHA landscape

Figure 2



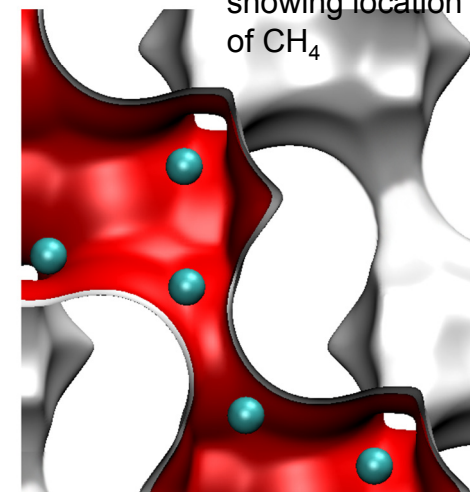
There are 6 cages per unit cell. The volume of one CHA cage is  $316.4 \text{ \AA}^3$ , slightly larger than that of a single cage of DDR ( $278 \text{ \AA}^3$ ), but significantly lower than FAU ( $786 \text{ \AA}^3$ ).



Snapshots showing location of  $\text{CH}_4$  and  $\text{CO}_2$

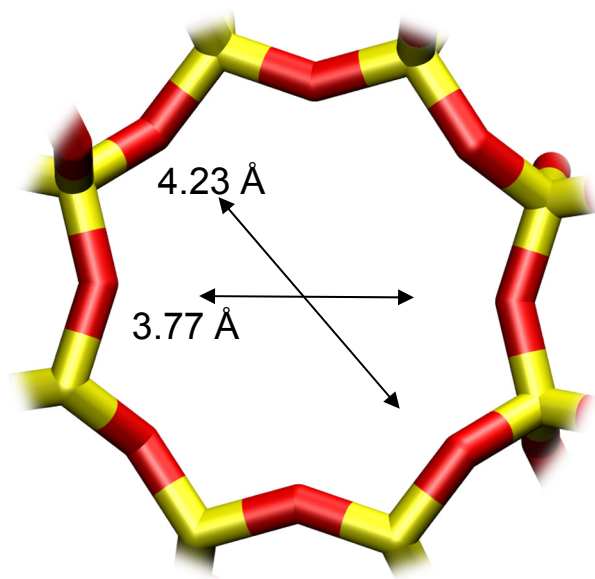


Snapshots showing location of  $\text{CH}_4$



Structural information from: C. Baerlocher, L.B. McCusker, Database of Zeolite Structures, International Zeolite Association, <http://www.iza-structure.org/databases/>

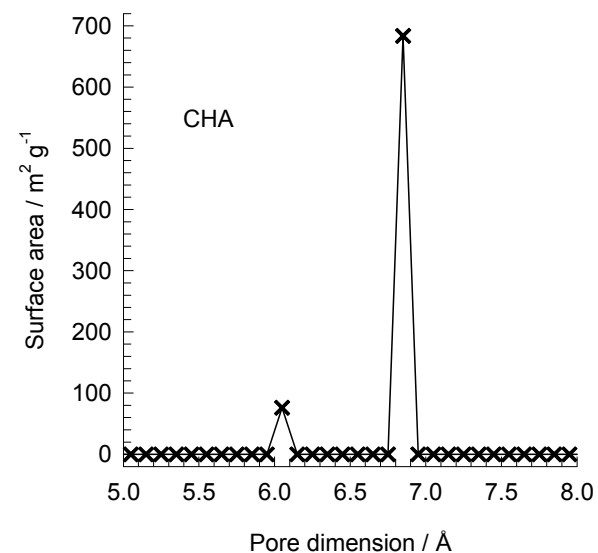
# CHA window and pore dimensions



CHA

The window dimensions calculated using the van der Waals diameter of framework atoms = 2.7 Å are indicated above by the arrows.

This plot of surface area versus pore dimension is determined using a combination of the DeLaunay triangulation method for pore dimension determination, and the procedure of Dürren for determination of the surface area.



	CHA
$a / \text{Å}$	15.075
$b / \text{Å}$	23.907
$c / \text{Å}$	13.803
Cell volume / $\text{Å}^3$	4974.574
conversion factor for [molec/uc] to [mol per kg Framework]	0.2312
conversion factor for [molec/uc] to [kmol/m <sup>3</sup> ]	0.8747
$\rho$ [kg/m <sup>3</sup> ]	1444.1
MW unit cell [g/mol(framework)]	4326.106
$\phi$ , fractional pore volume	0.382
open space / $\text{Å}^3/\text{uc}$	1898.4
Pore volume / $\text{cm}^3/\text{g}$	0.264
Surface area / $\text{m}^2/\text{g}$	758.0
DeLaunay diameter / $\text{Å}$	3.77



# SAPO-34 fits of experimental isotherms for pure components

Figure 4

Pure component isotherm fit data for SAPO-34 are fitted using the statistical isotherm

$$q_i = \frac{q_{i,sat}}{\Omega_i} \frac{b_i f_i + \sum_{m=2}^{\Omega_i} \frac{(b_i f_i)^m}{(m-1)!} \left[ \frac{1 - \frac{m}{\Omega_i + 1}}{1 - \frac{1}{\Omega_i + 1}} \right]^m}{1 + b_i f_i + \sum_{m=2}^{\Omega_i} \frac{(b_i f_i)^m}{(m)!} \left[ \frac{1 - \frac{m}{\Omega_i + 1}}{1 - \frac{1}{\Omega_i + 1}} \right]^m}$$

Molecule	$b_i$	$\Omega_i$	$q_{i,sat}$
CO <sub>2</sub>	$7.67 \times 10^{-5}$	6	8.2
CH <sub>4</sub>	$5.87 \times 10^{-6}$	6	8.2
N <sub>2</sub>	$1.26 \times 10^{-6}$	6	8.2
H <sub>2</sub>	$2.84 \times 10^{-7}$	9	12.3
O <sub>2</sub>	$1.2 \times 10^{-6}$	6	8.2
CO	$2.31 \times 10^{-6}$	6	8.2
Ar	$1.26 \times 10^{-6}$	6	8.2

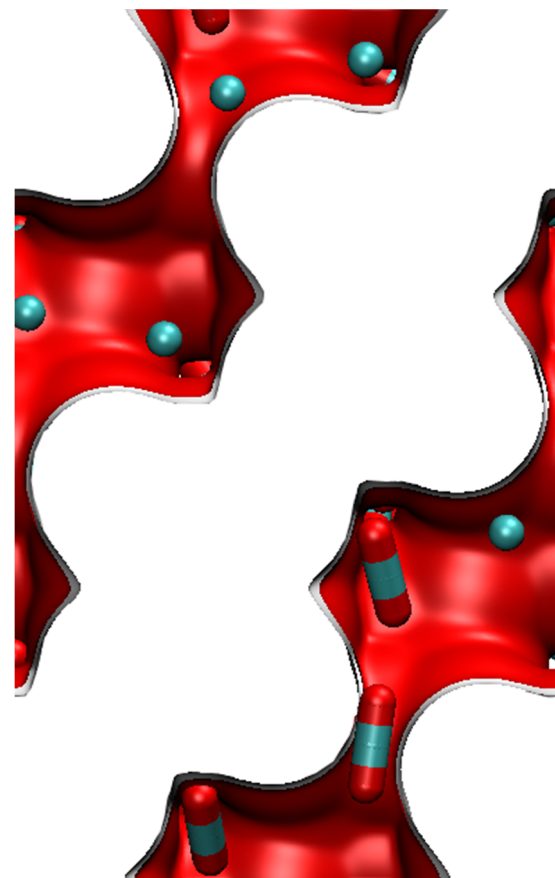
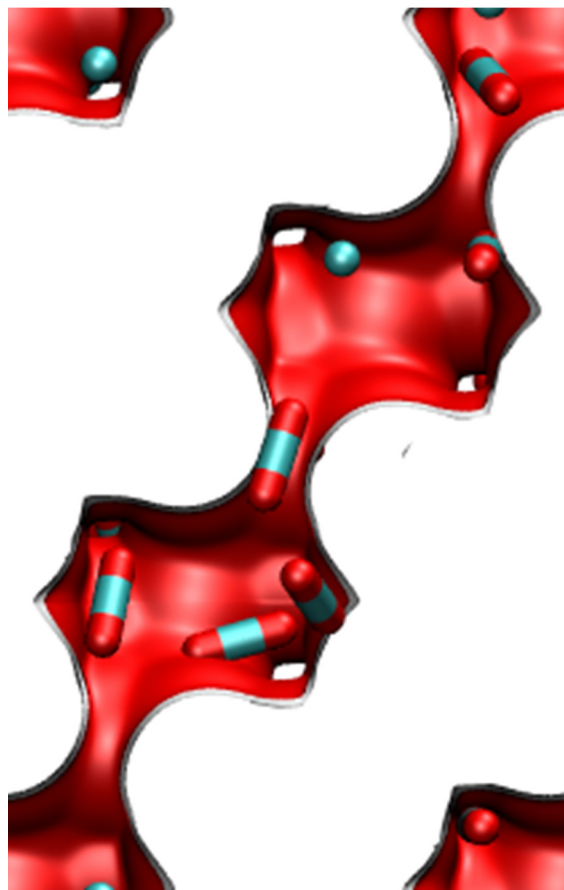
$b_i$  is expressed in Pa<sup>-1</sup>,  $\Omega_i$  in molecules per cage,  $q_{i,sat}$  in mol kg<sup>-1</sup>.

The original data is from:

S. Li, J.L. Falconer, R.D. Noble, R. Krishna, Interpreting unary, binary and ternary mixture permeation across a SAPO-34 membrane with loading-dependent Maxwell-Stefan diffusivities, J. Phys. Chem. C 111 (2007) 5075-5082.

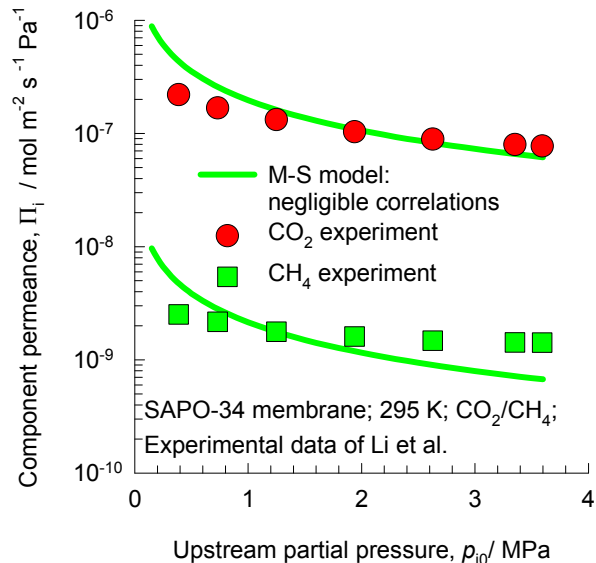
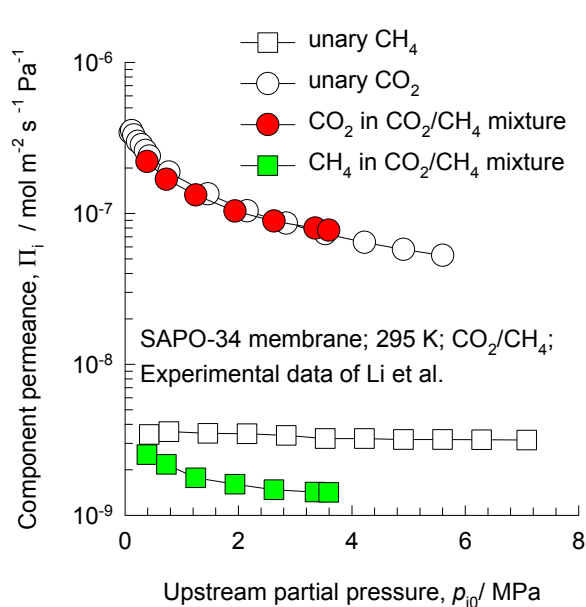
# SAPO-34: Snapshots of CO<sub>2</sub>/CH<sub>4</sub> mixtures

Figure 5



# SAPO-34: Maxwell-Stefan model calculations for CO<sub>2</sub>/CH<sub>4</sub> mixture permeation

Figure 6

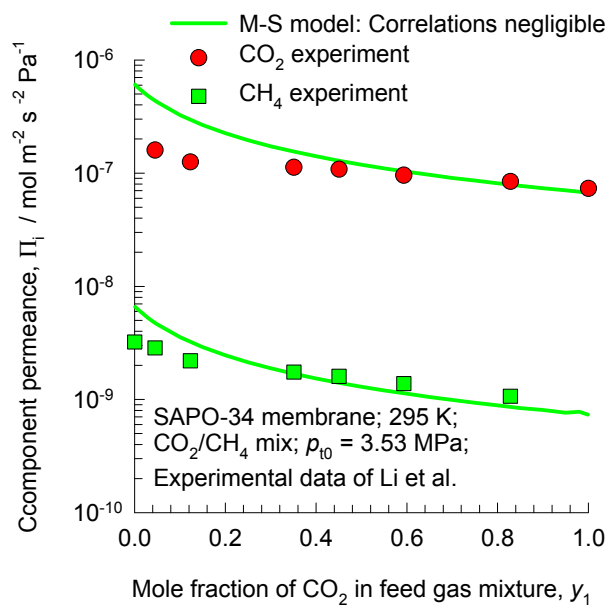
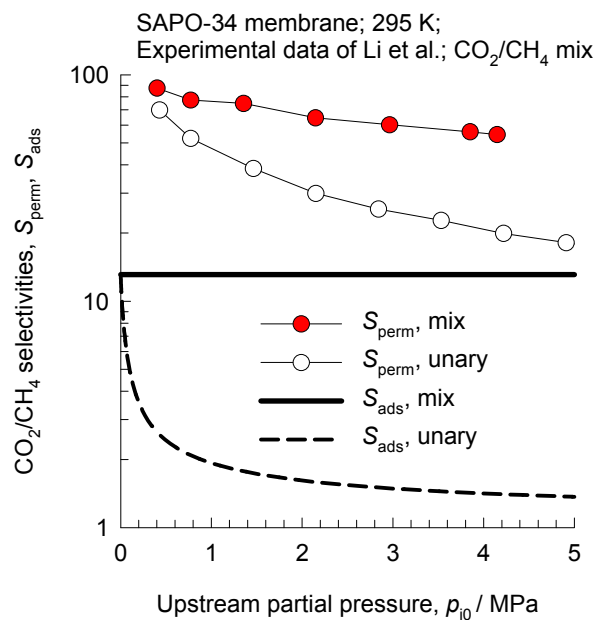


The inputs for the Maxwell-Stefan model are:

$$\frac{\rho D_1}{\delta} = 0.035 \text{ kg m}^{-2} \text{ s}^{-1}$$

$$\frac{\rho D_2}{\delta} = 0.005 \text{ kg m}^{-2} \text{ s}^{-1}$$

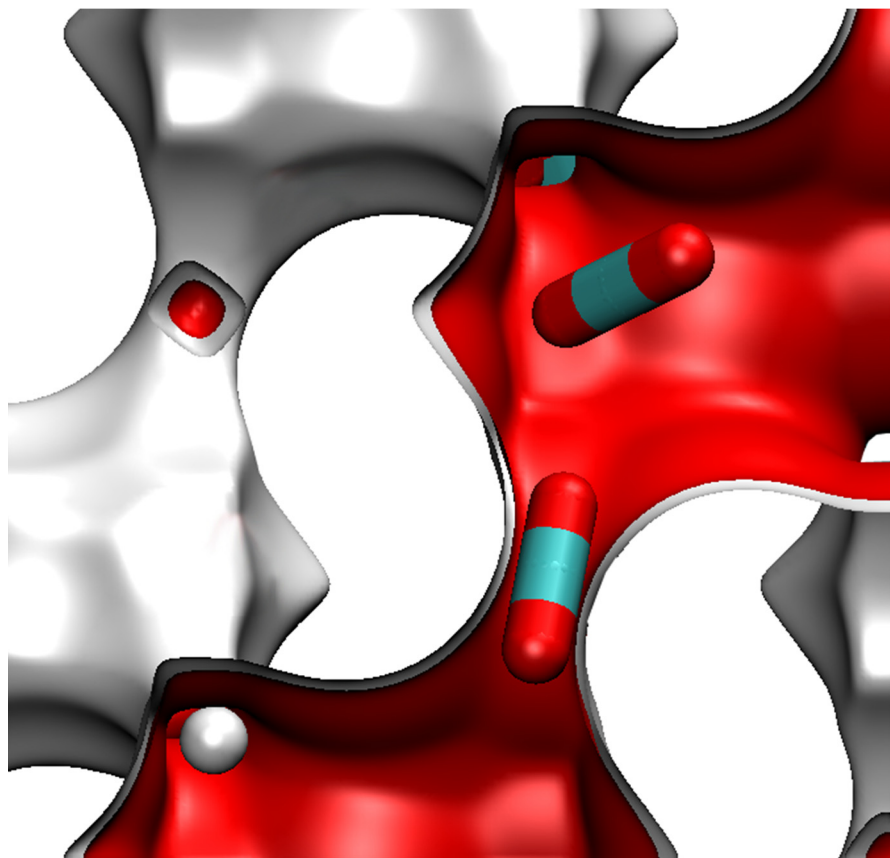
In the M-S model calculations, correlations are considered to be negligible



The same parameter inputs, as specified above, are equally good in estimation of component permeances when the composition of the feed mixture is varied, keeping the total upstream pressure = 3.53 MPa

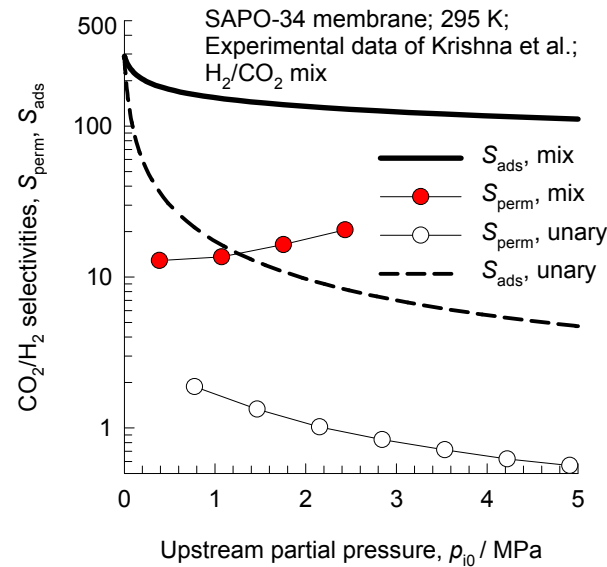
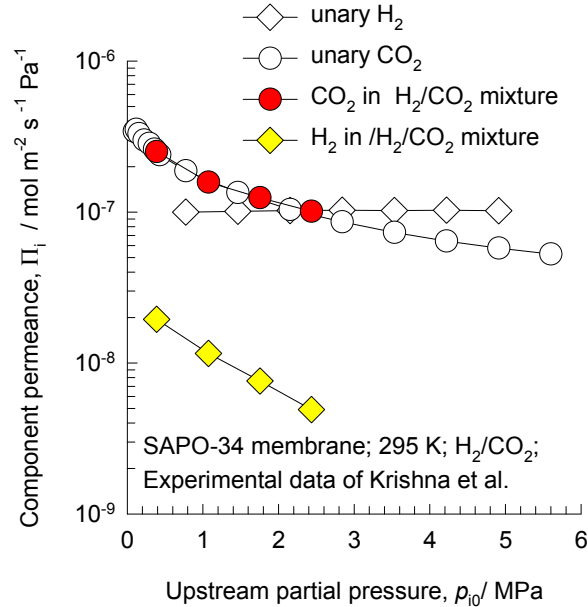
# SAPO-34: Snapshots of CO<sub>2</sub>/H<sub>2</sub> mixtures

Figure 7



# SAPO-34: Maxwell-Stefan model calculations for H<sub>2</sub>/CO<sub>2</sub> mixture permeation

Figure 8

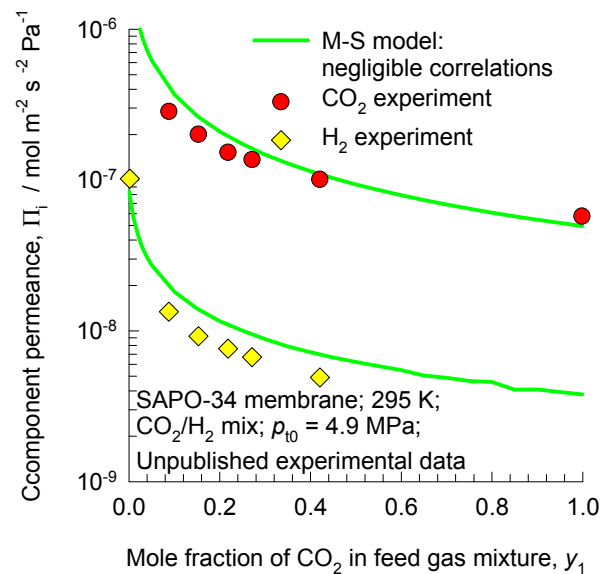
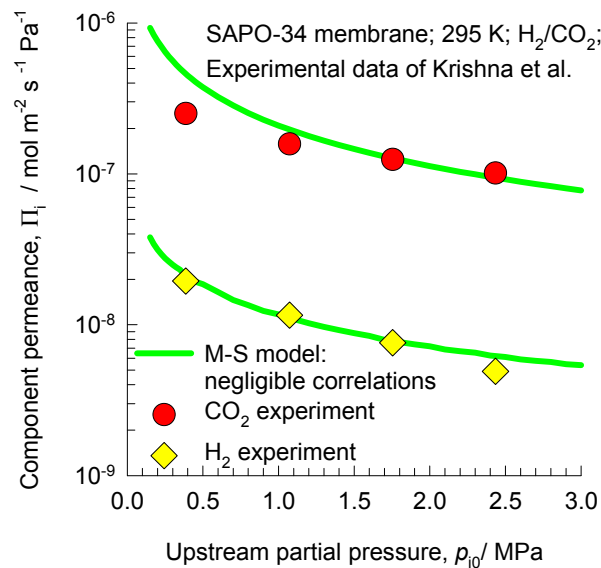


The inputs for the Maxwell-Stefan model are:

$$\frac{\rho D_1}{\delta} = 0.3 \text{ kg m}^{-2} \text{ s}^{-1}$$

$$\frac{\rho D_2}{\delta} = 0.035 \text{ kg m}^{-2} \text{ s}^{-1}$$

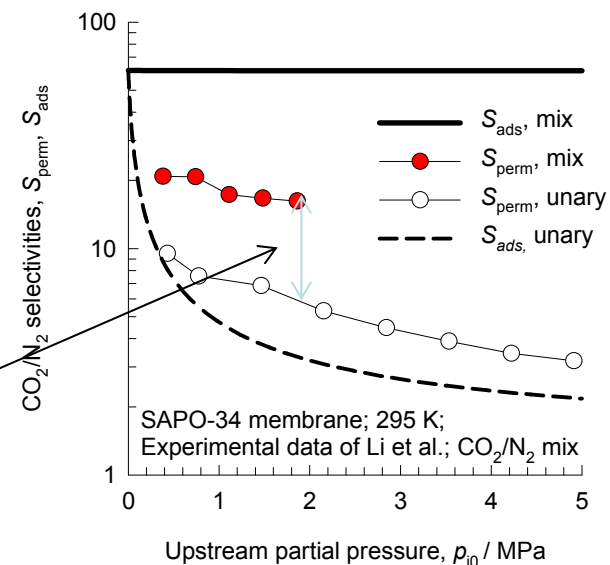
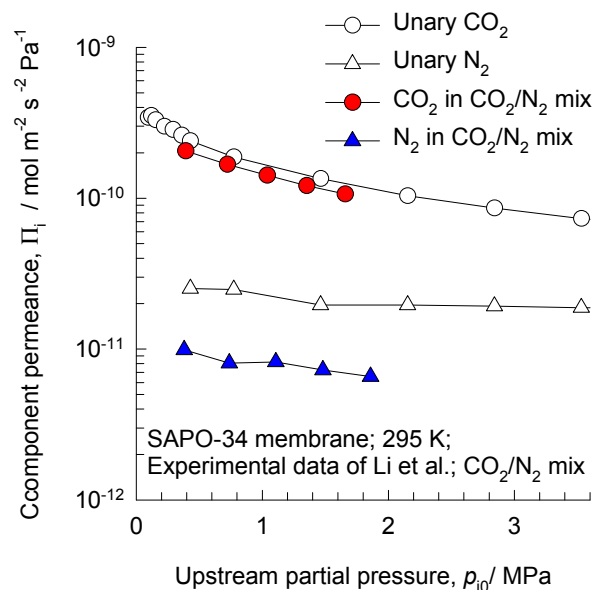
In the M-S model calculations, correlations are considered to be negligible



The same parameter inputs, as specified above, are equally good in estimation of component permeances when the composition of the feed mixture is varied, keeping the total upstream pressure = 4.9 MPa

# SAPO-34 CO<sub>2</sub>/N<sub>2</sub> mixture permeances, compared with unary permeances

Figure 9



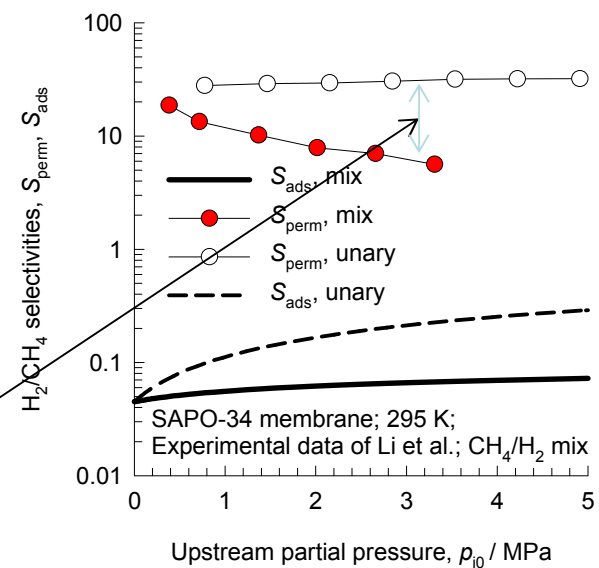
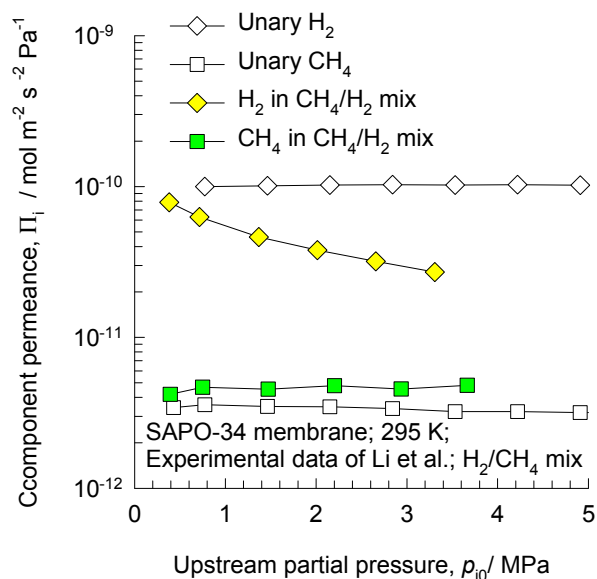
The increase in the permeation selectivity in the mixture is due primarily to the stronger adsorption selectivity. Correlation effects are not of major importance for this system.

The data are re-plotted using the information in:

S. Li, J.L. Falconer, R.D. Noble, R. Krishna, Interpreting unary, binary and ternary mixture permeation across a SAPO-34 membrane with loading-dependent Maxwell-Stefan diffusivities, *J. Phys. Chem. C* 111 (2007) 5075-5082

# SAPO-34 CH<sub>4</sub>/H<sub>2</sub> mixture permeances, compared with unary permeances

Figure 10



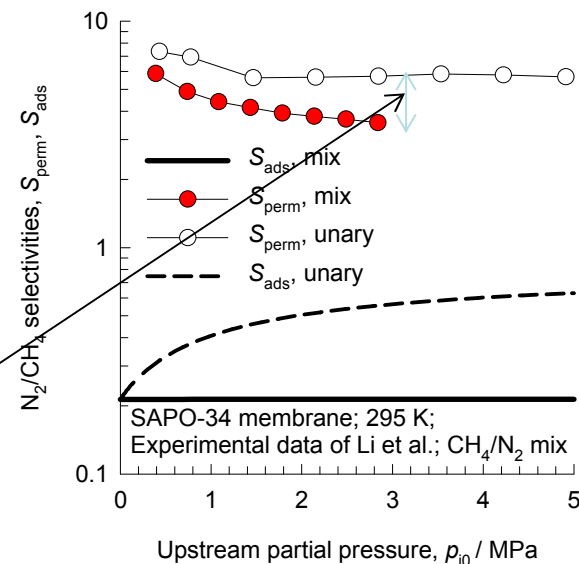
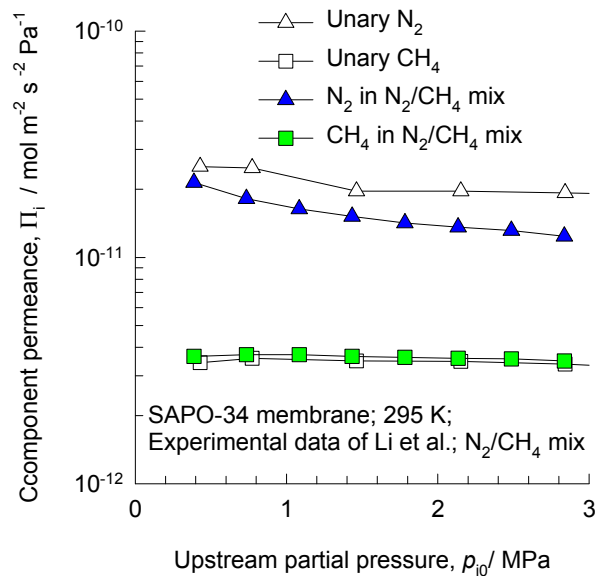
The reduction in the permeation selectivity in the mixture is due primarily to the decrease in adsorption selectivity. Correlation effects are not of major importance for this system.

The data are re-plotted using the information in:

S. Li, J.L. Falconer, R.D. Noble, R. Krishna, Interpreting unary, binary and ternary mixture permeation across a SAPO-34 membrane with loading-dependent Maxwell-Stefan diffusivities, J. Phys. Chem. C 111 (2007) 5075-5082

# SAPO-34 CH<sub>4</sub>/N<sub>2</sub> mixture permeances, compared with unary permeances

Figure 11



The reduction in the permeation selectivity in the mixture is due primarily to the decrease in adsorption selectivity. Correlation effects are not of major importance for this system.

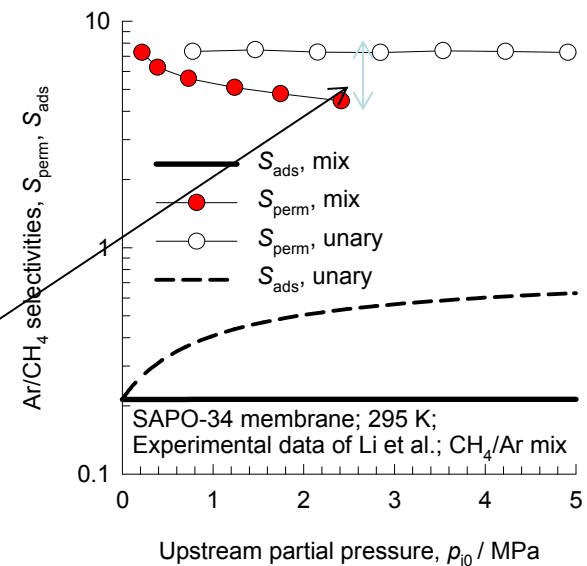
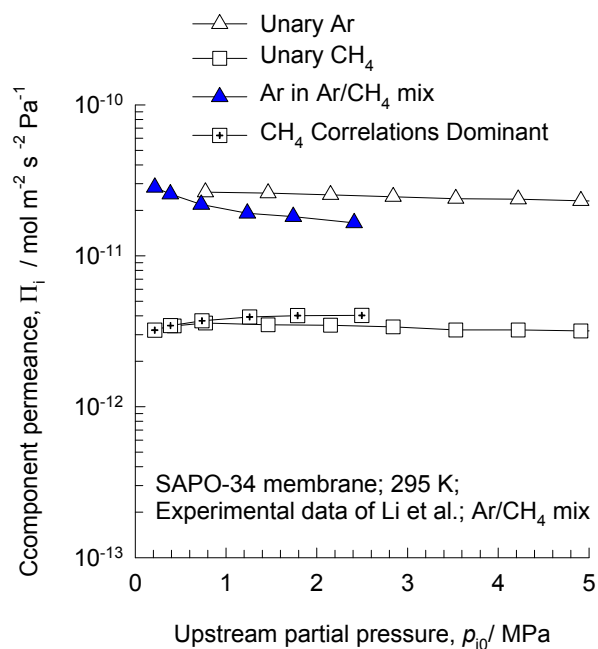
The data are re-plotted using the information in:

S. Li, J.L. Falconer, R.D. Noble, R. Krishna, Interpreting unary, binary and ternary mixture permeation across a SAPO-34 membrane with loading-dependent Maxwell-Stefan diffusivities, J. Phys. Chem. C 111 (2007) 5075-5082



# SAPO-34 CH<sub>4</sub>/Ar mixture permeances, compared with unary permeances

Figure 12



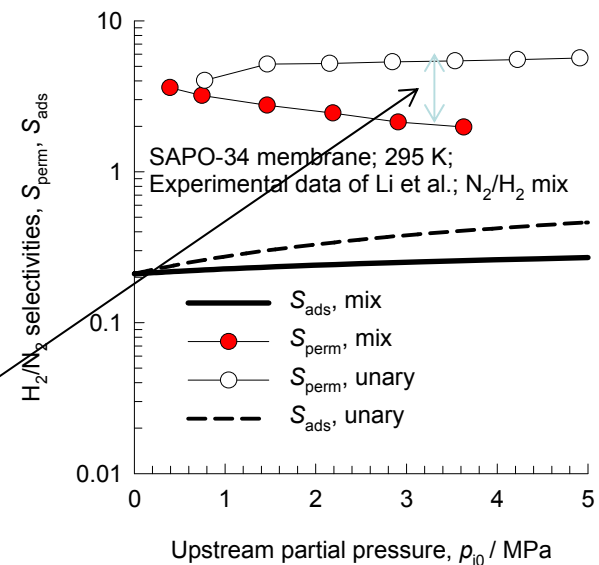
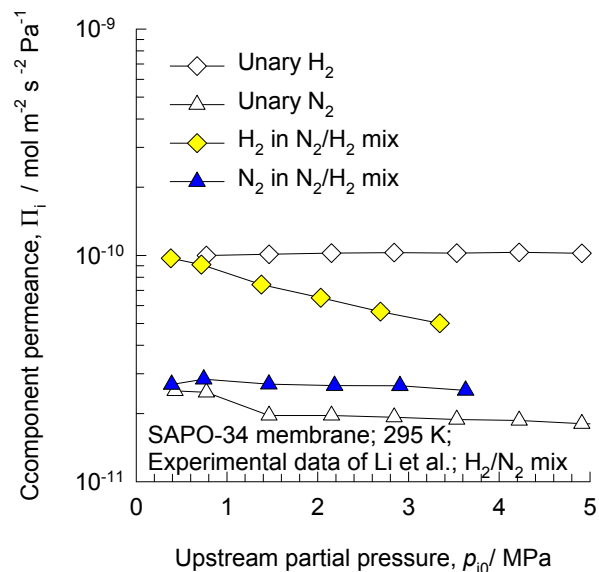
The reduction in the permeation selectivity in the mixture is due primarily to the decrease in adsorption selectivity. Correlation effects are not of major importance for this system.

The data are re-plotted using the information in:

S. Li, J.L. Falconer, R.D. Noble, R. Krishna, Interpreting unary, binary and ternary mixture permeation across a SAPO-34 membrane with loading-dependent Maxwell-Stefan diffusivities, J. Phys. Chem. C 111 (2007) 5075-5082

# SAPO-34 $N_2/H_2$ mixture permeances, compared with unary permeances

Figure 13



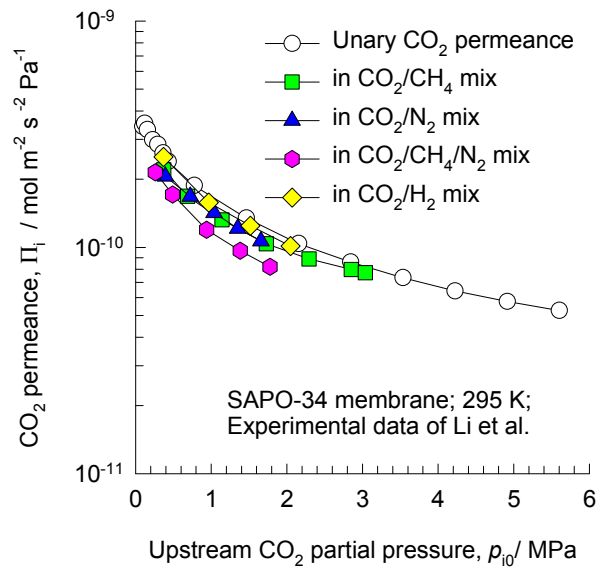
The reduction in the permeation selectivity in the mixture is due primarily to the decrease in adsorption selectivity. Correlation effects are not of major importance for this system.

The data are re-plotted using the information in:

S. Li, J.L. Falconer, R.D. Noble, R. Krishna, Interpreting unary, binary and ternary mixture permeation across a SAPO-34 membrane with loading-dependent Maxwell-Stefan diffusivities, *J. Phys. Chem. C* 111 (2007) 5075-5082

# SAPO-34 CO<sub>2</sub> permeance: influence of partner species

Figure 14



The permeance of CO<sub>2</sub> is practically unchanged as the partner molecules are changed. Correlations effects are not of importance for the more strongly adsorbed species.

The data are re-plotted using the information in:

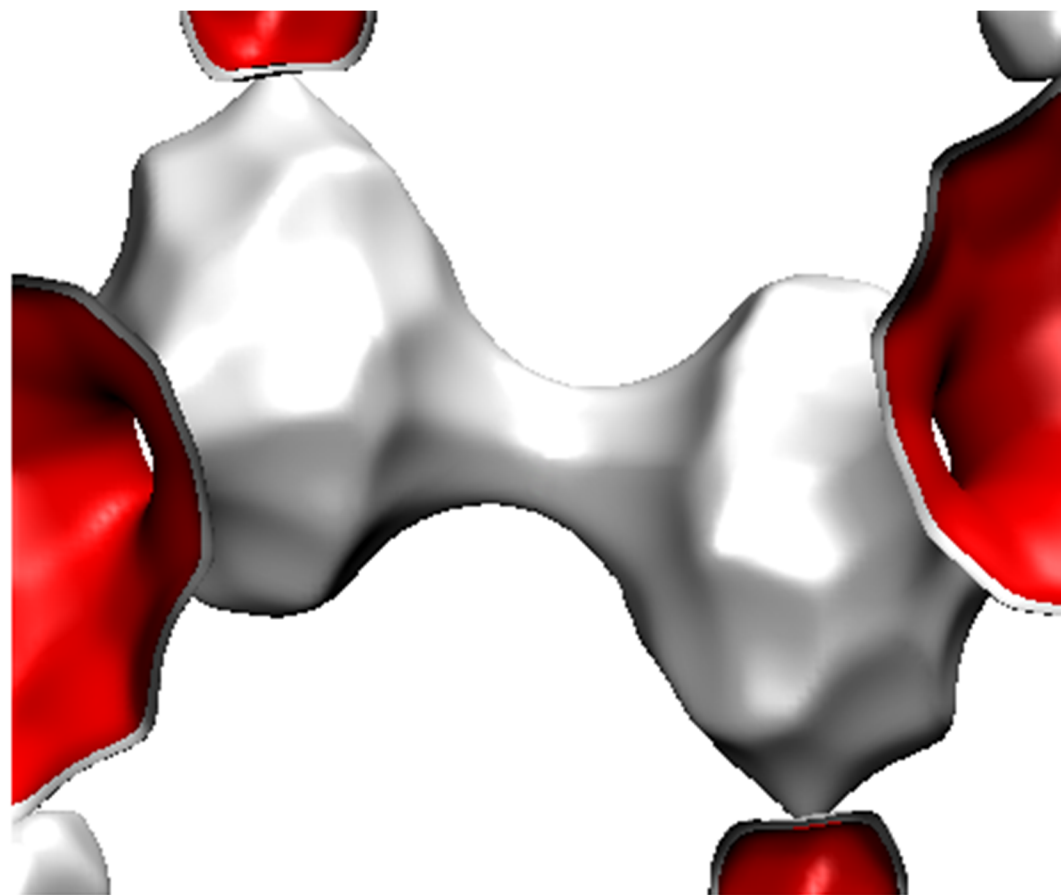
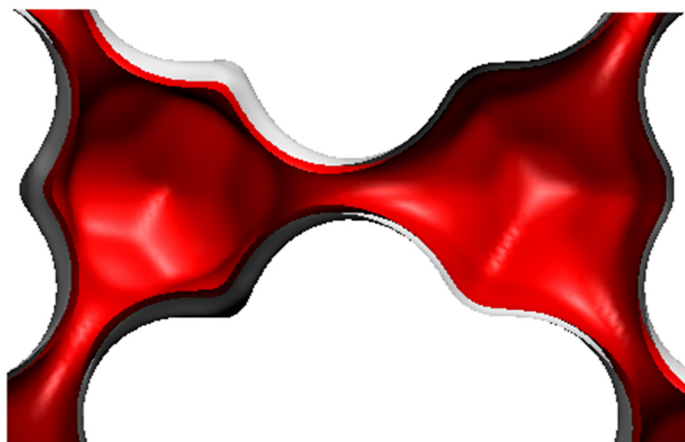
S. Li, J.L. Falconer, R.D. Noble, R. Krishna, Interpreting unary, binary and ternary mixture permeation across a SAPO-34 membrane with loading-dependent Maxwell-Stefan diffusivities, *J. Phys. Chem. C* 111 (2007) 5075-5082

**DDR**  
**(all-silica)**

# DDR landscape

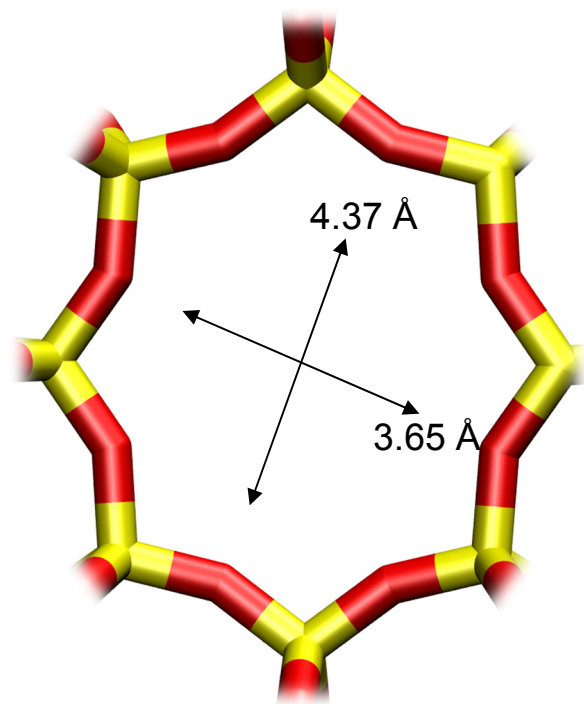
Figure 16

There are 12 cages per unit cell.  
The volume of one DDR cage is  $278 \text{ \AA}^3$ , significantly smaller than that of a single cage of FAU ( $786 \text{ \AA}^3$ ), or ZIF-8 ( $1168 \text{ \AA}^3$ ).



Structural information from: C. Baerlocher, L.B. McCusker, Database of Zeolite Structures, International Zeolite Association, <http://www.iza-structure.org/databases/>

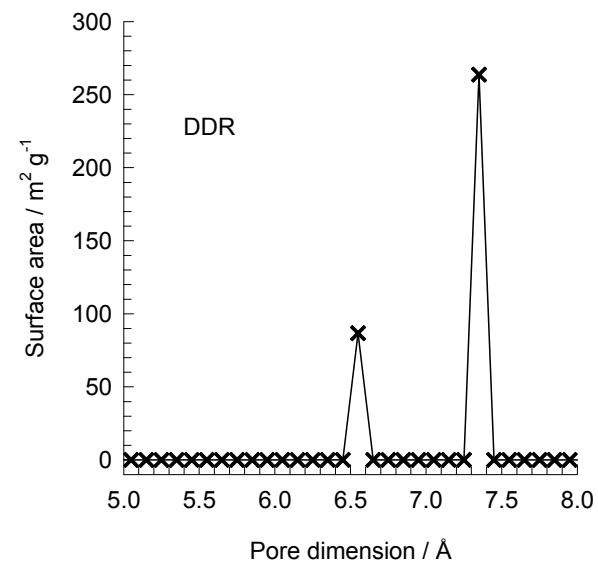
# DDR window and pore dimensions



DDR

The window dimensions calculated using the van der Waals diameter of framework atoms = 2.7 Å are indicated above by the arrows.

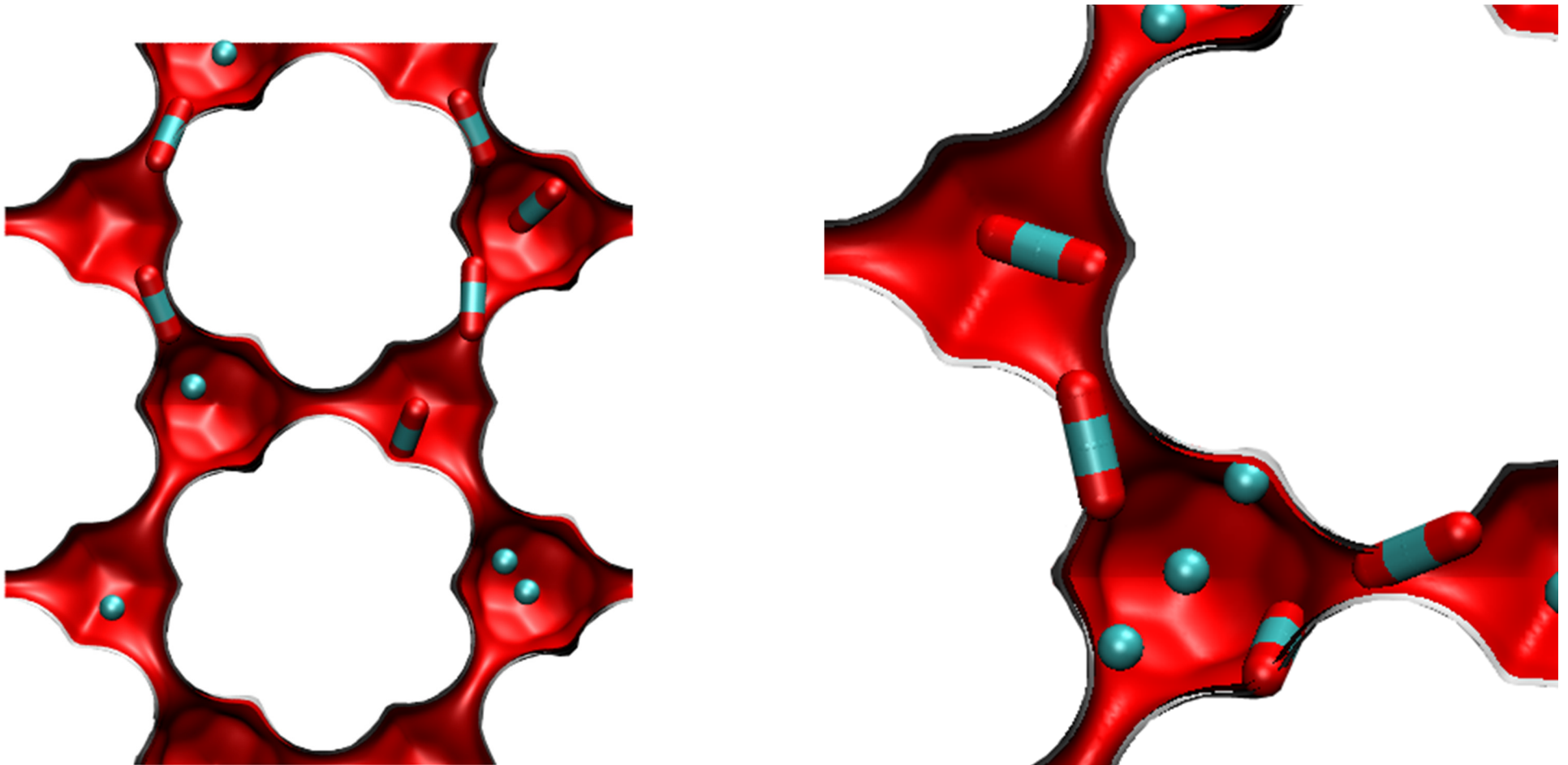
This plot of surface area versus pore dimension is determined using a combination of the DeLaunay triangulation method for pore dimension determination, and the procedure of Dürren for determination of the surface area.



	DDR
$a / \text{Å}$	24.006
$b / \text{Å}$	13.86
$c / \text{Å}$	40.892
Cell volume / $\text{Å}^3$	13605.72
conversion factor for [molec/uc] to [mol per kg Framework]	0.0693
conversion factor for [molec/uc] to [kmol/m <sup>3</sup> ]	0.4981
$\rho$ [kg/m <sup>3</sup> ]	1759.991
MW unit cell [g/mol(framework)]	14420.35
$\phi$ , fractional pore volume	0.245
open space / $\text{Å}^3/\text{uc}$	3333.5
Pore volume / cm <sup>3</sup> /g	0.139
Surface area / m <sup>2</sup> /g	350.0
DeLaunay diameter / $\text{Å}$	3.65

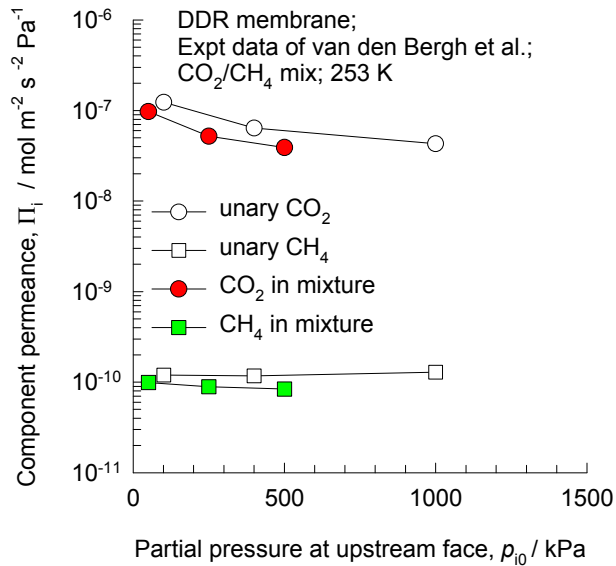
# DDR snapshots of CO<sub>2</sub>/CH<sub>4</sub> mixture

Figure 18

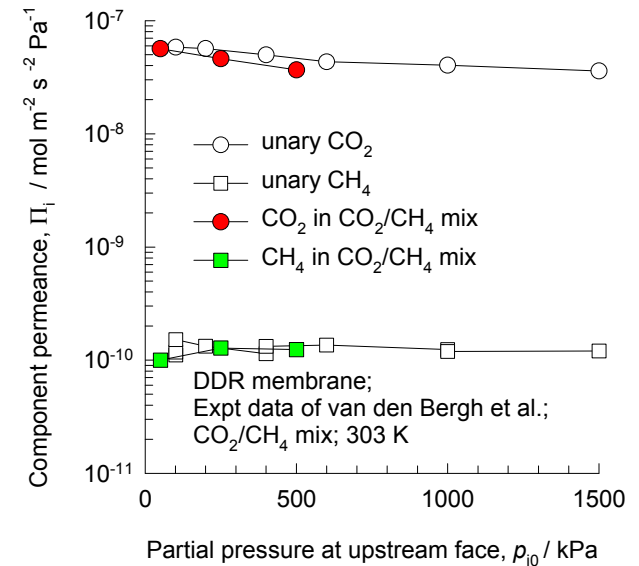


# DDR $\text{CO}_2/\text{CH}_4$ mixture permeances, compared with unary permeances

Figure 19



The permeance of each species remains practically unchanged in the mixture. This indicates that correlation effects are not significant.



The experimental data are from:

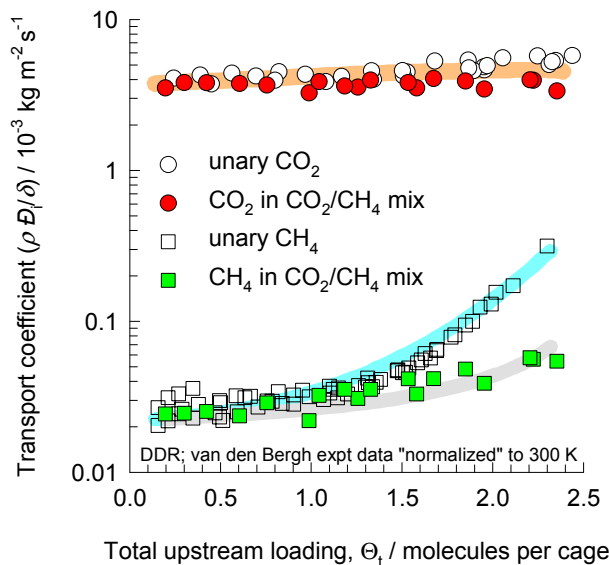
J. van den Bergh, W. Zhu, J. Gascon, J.A. Moulijn, F. Kapteijn, Separation and Permeation Characteristics of a DD3R Zeolite Membrane, *J. Membr. Sci.* 316 (2008) 35-45.

J. van den Bergh, W. Zhu, J.C. Groen, F. Kapteijn, J.A. Moulijn, K. Yajima, K. Nakayama, T. Tomita, S. Yoshida, Natural Gas Purification with a DDR Zeolite Membrane; Permeation Modelling with Maxwell-Stefan Equations, *Stud. Surf. Sci. Catal.* 170 (2007) 1021-1027.

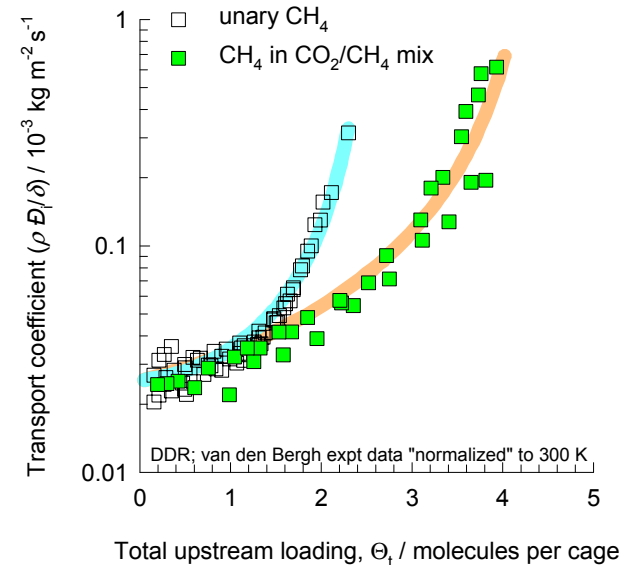


# DDR Transport coefficients: unary vs mixture

Figure 20



A more detailed analysis of the permeation experiments reveals that the transport of CH<sub>4</sub> in the mixture is hindered due to the preponderance of CO<sub>2</sub> molecules lodged at the window sites.



The experimental data are from:

J. van den Bergh, W. Zhu, J. Gascon, J.A. Moulijn, F. Kapteijn, Separation and Permeation Characteristics of a DD3R Zeolite Membrane, *J. Membr. Sci.* 316 (2008) 35-45.

J. van den Bergh, W. Zhu, J.C. Groen, F. Kapteijn, J.A. Moulijn, K. Yajima, K. Nakayama, T. Tomita, S. Yoshida, Natural Gas Purification with a DDR Zeolite Membrane; Permeation Modelling with Maxwell-Stefan Equations, *Stud. Surf. Sci. Catal.* 170 (2007) 1021-1027.

These data were modeled using the Maxwell-Stefan diffusion equations, along with Reed & Ehrlich model for the concentration dependence of the M-S diffusivities. The modeling work is reported in our previous works:

R. Krishna, J.M. van Baten, Onsager coefficients for binary mixture diffusion in nanopores, *Chem. Eng. Sci.* 63 (2008) 3120-3140.

R. Krishna, J.M. van Baten, Segregation effects in adsorption of CO<sub>2</sub> containing mixtures and their consequences for separation selectivities in cage-type zeolites, *Sep. Purif. Technol.* 61 (2008) 414-423.

R. Krishna, J.M. van Baten, A molecular dynamics investigation of a variety of influences of temperature on diffusion in zeolites, *Microporous Mesoporous Mater.* 125 (2009) 126-134.

C. Chmelik, R. Krishna, Hindering effects in CO<sub>2</sub>/CH<sub>4</sub> mixture diffusion in ZIF-8 crystals, *J. Membr. Sci.* 397-398 (2012) 87-91.

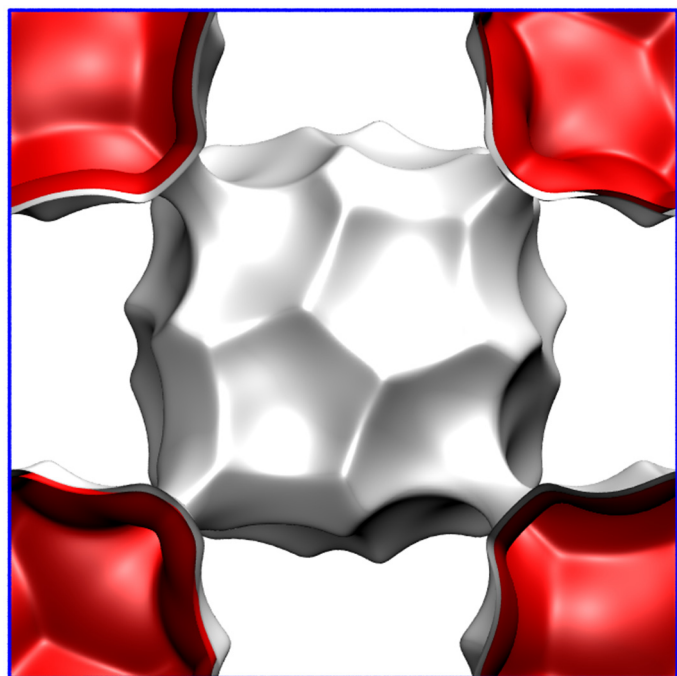
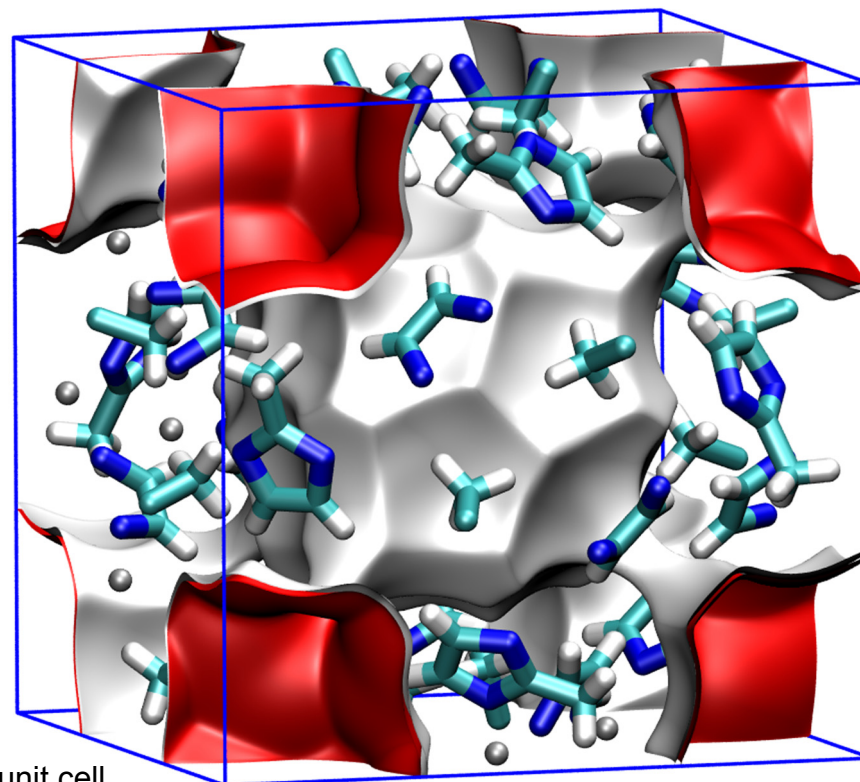
**ZIF-8**

# ZIF-8 pore landscapes

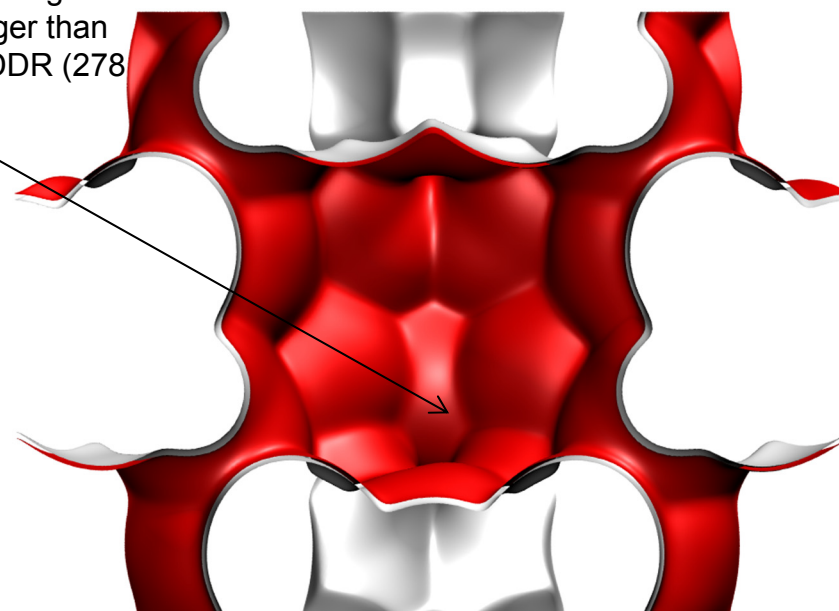
The ZIF-8 =  $\text{Zn}(\text{methylimidazole})_2$  structure was taken from

R. Banerjee, A. Phan, B. Wang, C. Knobler, H. Furukawa, M. O'Keeffe, O.M. Yaghi, High-Throughput Synthesis of Zeolitic Imidazolate Frameworks and Application to  $\text{CO}_2$  Capture, *Science* 319 (2008) 939-943.

*The original structural data (cif file) contains solvent molecules; these were removed and the solvent-free structures were simulated.*

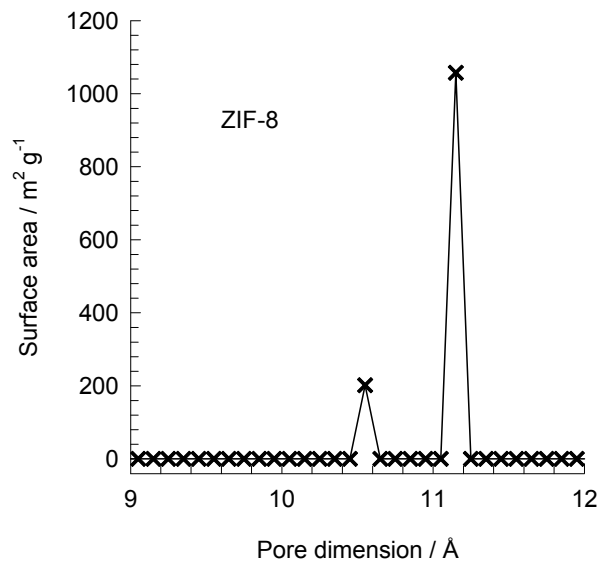


There are 2 cages per unit cell. The volume of one ZIF-8 cage is  $1168 \text{ \AA}^3$ , significantly larger than that of a single cage of DDR ( $278 \text{ \AA}^3$ ), or FAU ( $786 \text{ \AA}^3$ ).



# ZIF-8 dimensions

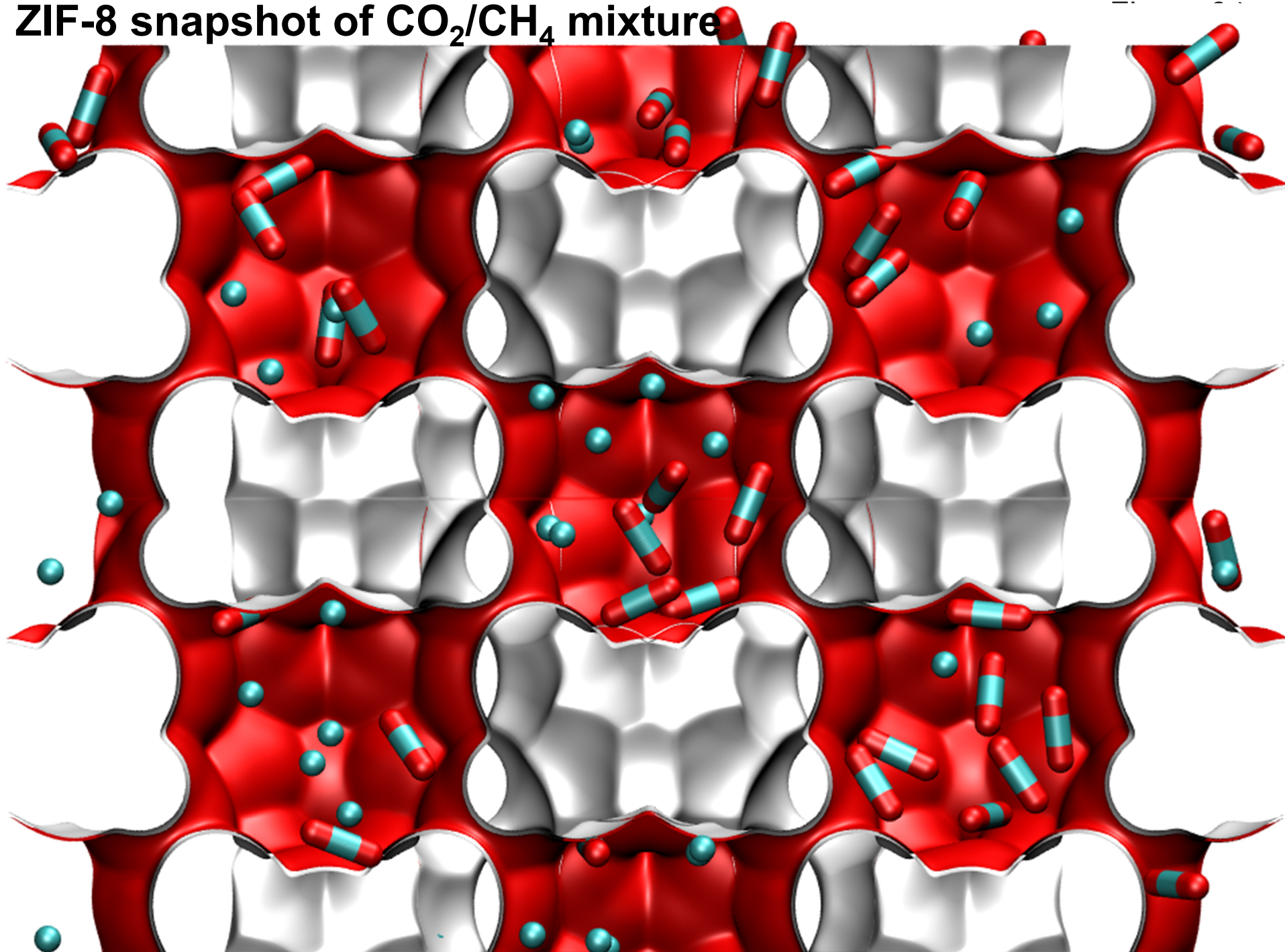
Figure 23



This plot of surface area versus pore dimension is determined using a combination of the DeLaunay triangulation method for pore dimension determination, and the procedure of Düren for determination of the surface area.

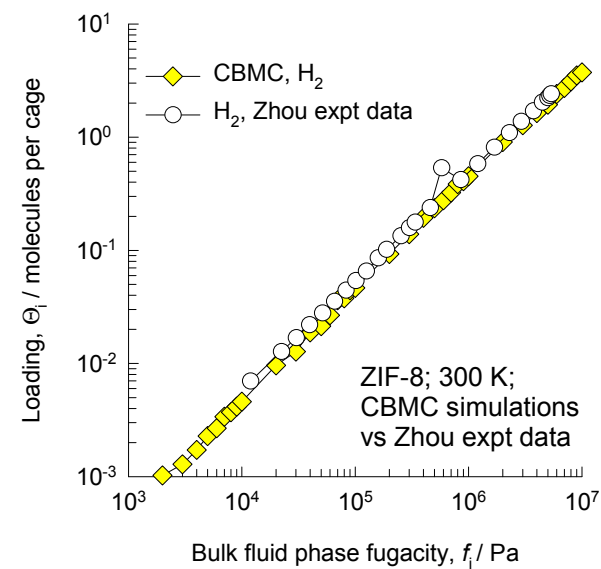
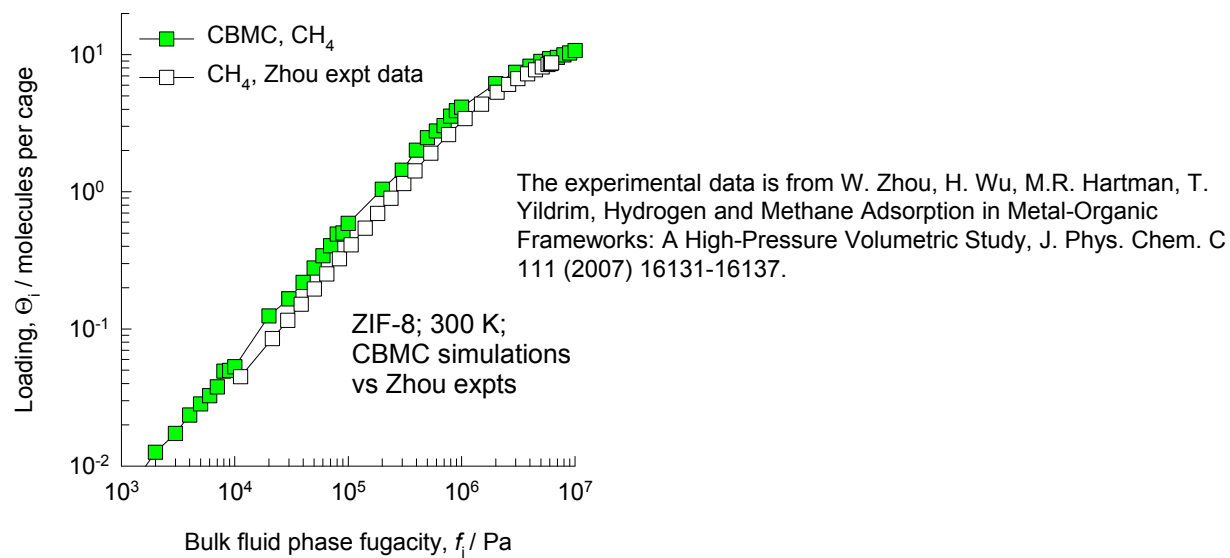
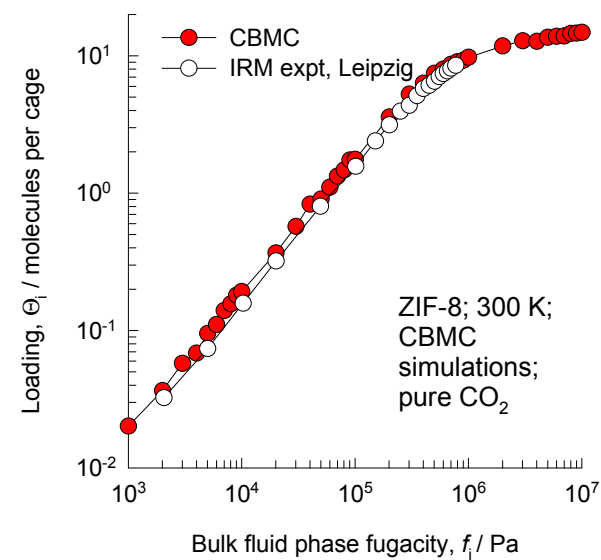
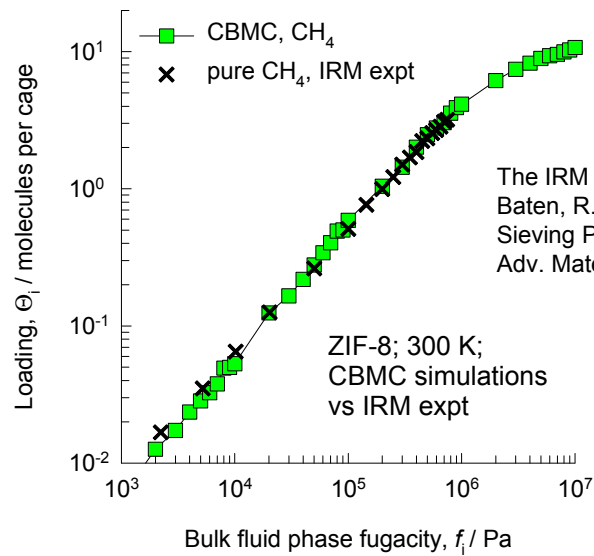
	ZIF-8
$a / \text{Å}$	16.991
$b / \text{Å}$	16.991
$c / \text{Å}$	16.991
Cell volume / $\text{Å}^3$	4905.201
conversion factor for [molec/uc] to [mol per kg Framework]	0.3663
conversion factor for [molec/uc] to [kmol/m <sup>3</sup> ]	0.7106
$\rho$ [kg/m <sup>3</sup> ]	924.253
MW unit cell [g/mol(framework)]	2730.182
$\phi$ , fractional pore volume	0.476
open space / $\text{Å}^3/\text{uc}$	2337.0
Pore volume / cm <sup>3</sup> /g	0.515
Surface area / m <sup>2</sup> /g	1164.7
DeLaunay diameter / $\text{Å}$	3.26

# ZIF-8 snapshot of CO<sub>2</sub>/CH<sub>4</sub> mixture



# ZIF-8 pure component isotherms; comparison with experiments

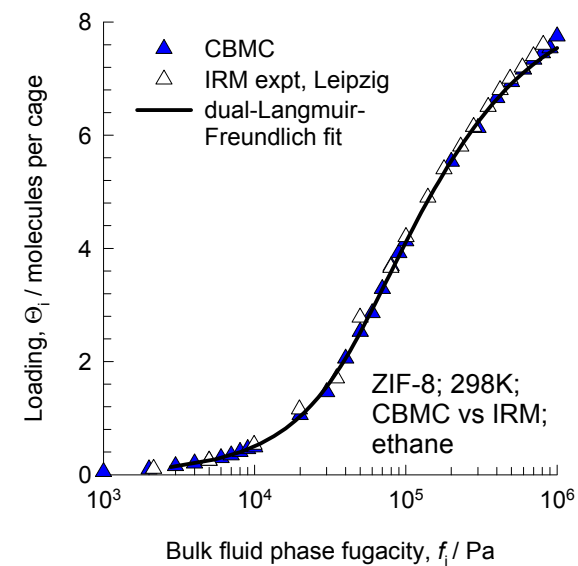
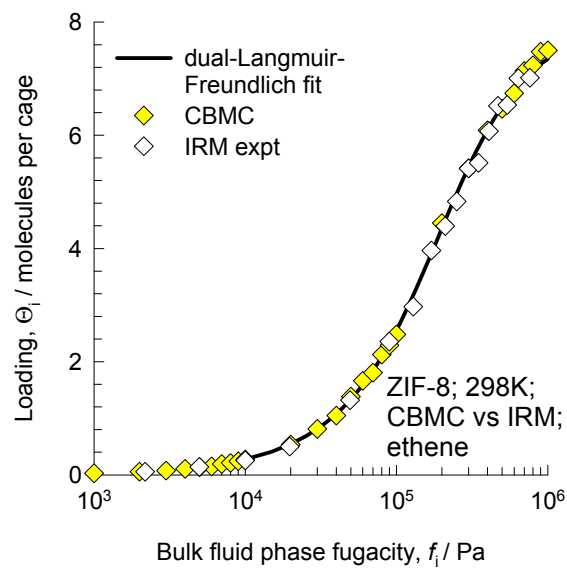
Figure 25





# ZIF-8 pure component isotherms; comparison with experiments

Figure 26

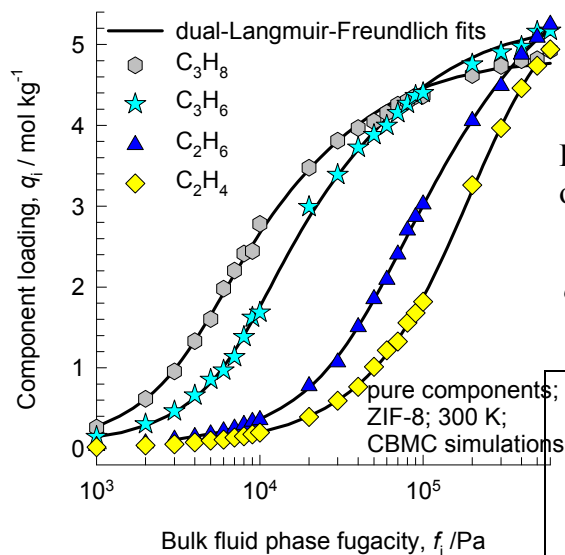


The IRM experimental data is from H. Bux, C. Chmelik, R. Krishna, J. Caro, Ethene/Ethane Separation by the MOF Membrane ZIF-8: Molecular Correlation of Permeation, Adsorption, Diffusion, J. Membr. Sci. 369 (2011) 284-289.

# ZIF-8

## Fitting parameters for pure component isotherms determined from CBMC simulations

Figure 27



Dual-site Langmuir-Freundlich parameters for pure components in ZIF-8 at 300 K. These fits are based on CBMC simulations.

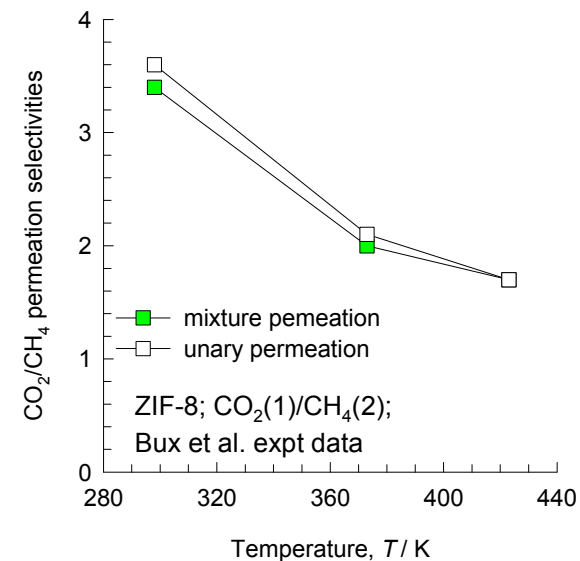
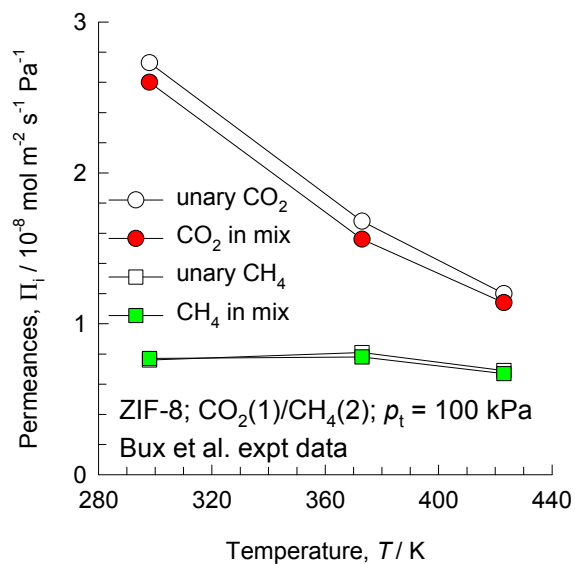
$$q_i = q_{i,A,sat} \frac{b_{i,A} f_i^{V_A}}{1 + b_{i,A} f_i^{V_A}} + q_{i,B,sat} \frac{b_{i,B} f_i^{V_B}}{1 + b_{i,B} f_i^{V_B}}$$

	Site A			Site B		
	$q_{i,A,sat}$ mol/kg	$b_{i,A}$ $\text{Pa}^{-V_i}$	$V_{i,A}$ dimensionless	$q_{i,B,sat}$ mol/kg	$b_{i,B}$ $\text{Pa}^{-V_i}$	$V_{i,B}$ dimensionless
CO <sub>2</sub>	1.5	$2.85 \times 10^{-15}$	2.7	9.8	$1.43 \times 10^{-6}$	1
CH <sub>4</sub>	9.5	$4.37 \times 10^{-7}$	1	3.7	$6.7 \times 10^{-9}$	1
H <sub>2</sub>	19	$2.37 \times 10^{-9}$	0.87	20	$1.61 \times 10^{-8}$	1
ethene	0.9	$2.04 \times 10^{-5}$	1	4.9	$6.97 \times 10^{-9}$	1.54
ethane	4.8	$7.4 \times 10^{-6}$	1	1.3	$2.89 \times 10^{-10}$	2
propene	1.5	$2.45 \times 10^{-10}$	2.4	3.8	$3.63 \times 10^{-5}$	1
propane	1.55	$1.22 \times 10^{-8}$	2.13	3.3	$7.57 \times 10^{-5}$	1



# ZIF-8 CO<sub>2</sub>/CH<sub>4</sub> unary and mixture permeances

Figure 28

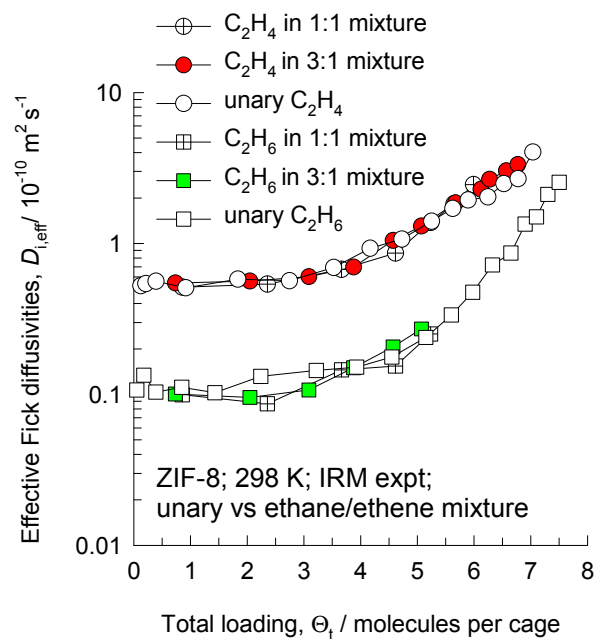


The data on membrane permeances are obtained from:

H. Bux, C. Chmelik, J.M. Van Baten, R. Krishna, J. Caro, Novel MOF-Membrane for Molecular Sieving Predicted by IR-Diffusion Studies and Molecular Modeling, *Adv. Mater.* 22 (2010) 4741-4743.

# ZIF-8 ethane/ethane diffusion at 298 K: IRM experimental data

Figure 29

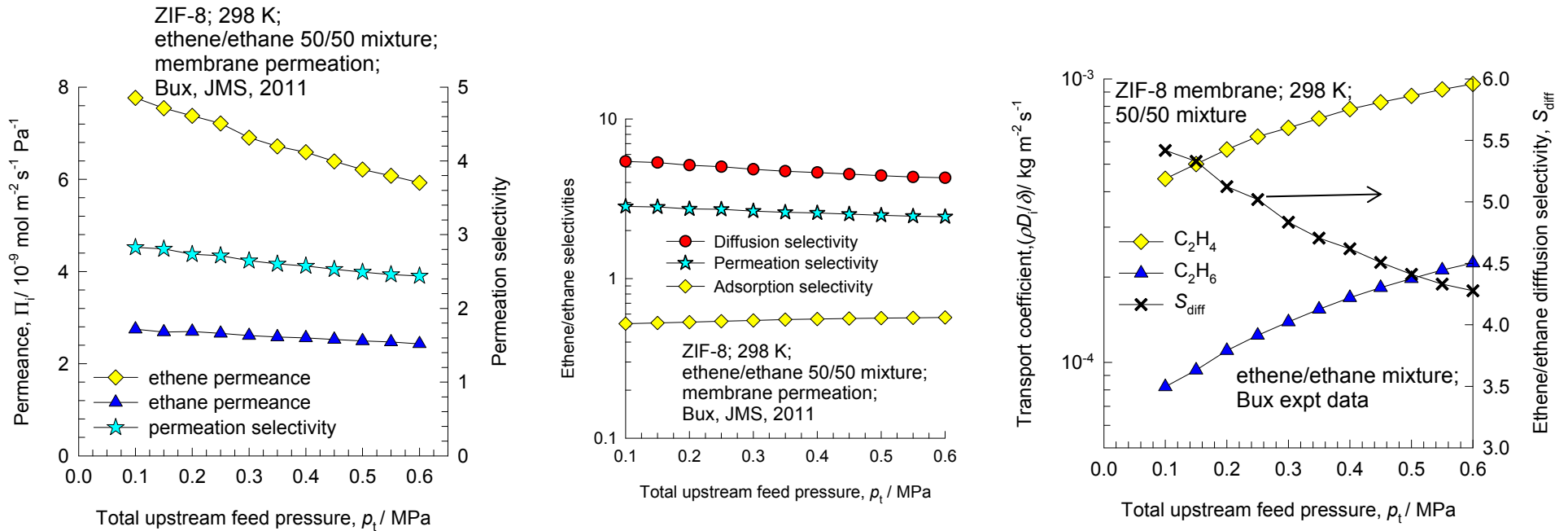


The IRM experimental data are re-plotted using the information in:

H. Bux, C. Chmelik, R. Krishna, J. Caro, Ethene/Ethane Separation by the MOF Membrane ZIF-8: Molecular Correlation of Permeation, Adsorption, Diffusion, *J. Membr. Sci.* 369 (2011) 284-289.

# ZIF-8 ethene/ethane mixture permeances and transport coefficients

Figure 30



The data on membrane permeances are obtained from:

H. Bux, C. Chmelik, R. Krishna, J. Caro, Ethene/Ethane Separation by the MOF Membrane ZIF-8: Molecular Correlation of Permeation, Adsorption, Diffusion, J. Membr. Sci. 369 (2011) 284-289.

The calculations of  $S_{\text{ads}}$  are on the basis of IAST, with dual-Langmuir-Freundlich fits of pure component isotherms.

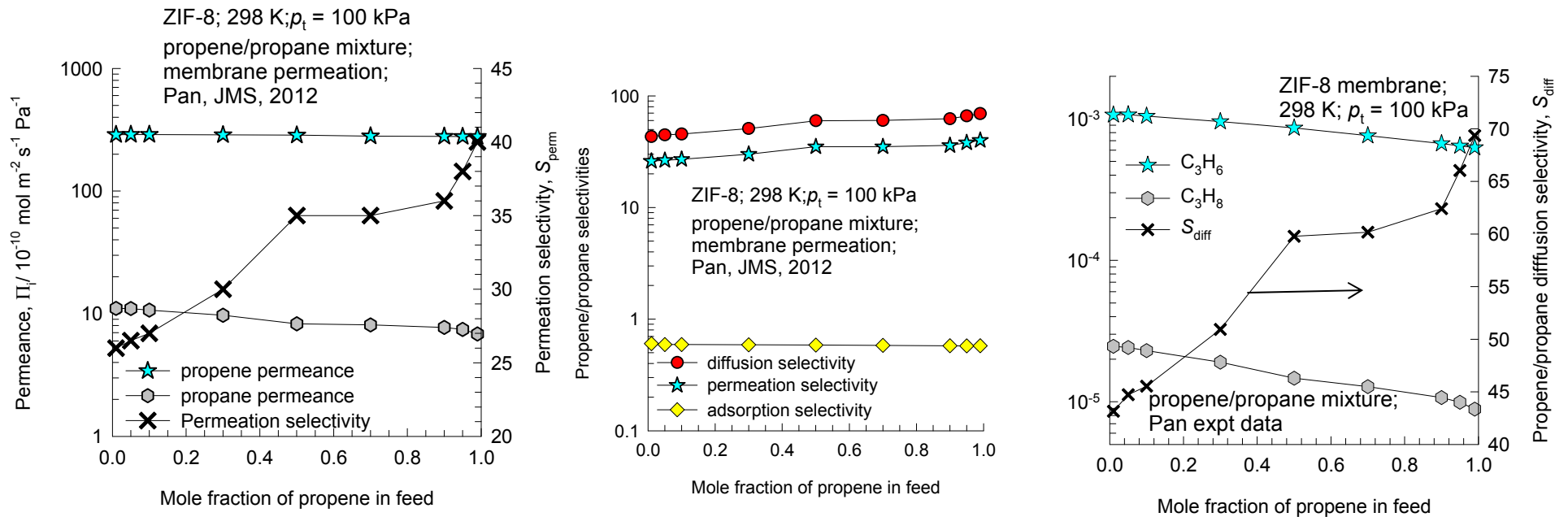
The transport coefficients are calculated from

$$\frac{\rho D_i}{\delta} = \frac{\Pi_i}{\frac{q_i}{p_i}}$$

with the molar loadings  $q_i$  at the upstream face determined from IAST calculations.

# ZIF-8 propene/propane mixture permeances and transport coefficients

Figure 31



The data on membrane permeances are obtained from:

Y. Pan, T. Li, G. Lestari, Z. Lai, Effective Separation of Propylene/Propane Binary Mixtures by ZIF-8 Membranes, J. Membr. Sci. 390-391 (2012) 93-98.

The calculations of  $S_{\text{ads}}$  are on the basis of IAST, with dual-Langmuir-Freundlich fits of pure component isotherms.

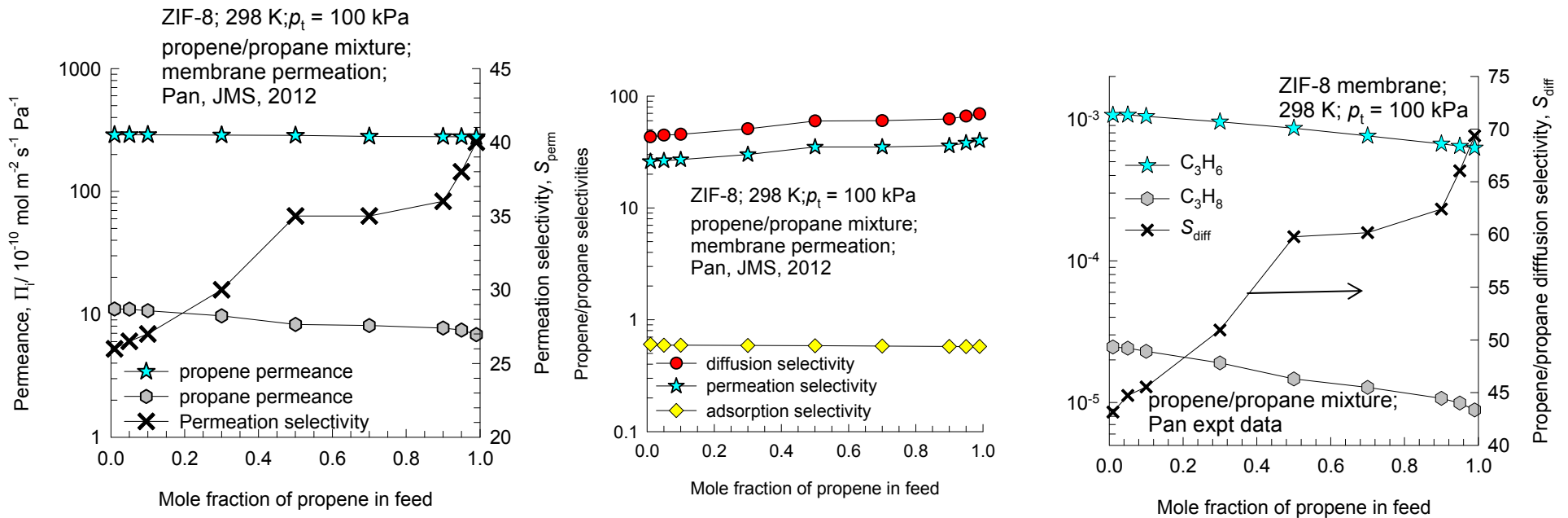
The transport coefficients are calculated from

$$\frac{\rho D_i}{\delta} = \frac{\Pi_i}{\frac{q_i}{p_i}}$$

with the molar loadings  $q_i$  at the upstream face determined from IAST calculations.

# ZIF-8 propene/propane mixture permeances and transport coefficients

Figure 32



The data on membrane permeances are obtained from:

Y. Pan, T. Li, G. Lestari, Z. Lai, Effective Separation of Propylene/Propane Binary Mixtures by ZIF-8 Membranes, J. Membr. Sci. 390-391 (2012) 93-98.

The calculations of  $S_{ads}$  are on the basis of IAST, with dual-Langmuir-Freundlich fits of pure component isotherms.

The transport coefficients are calculated from

$$\frac{\rho D_i}{\delta} = \frac{\Pi_i}{\frac{q_i}{p_i}}$$

with the molar loadings  $q_i$  at the upstream face determined from IAST calculations.

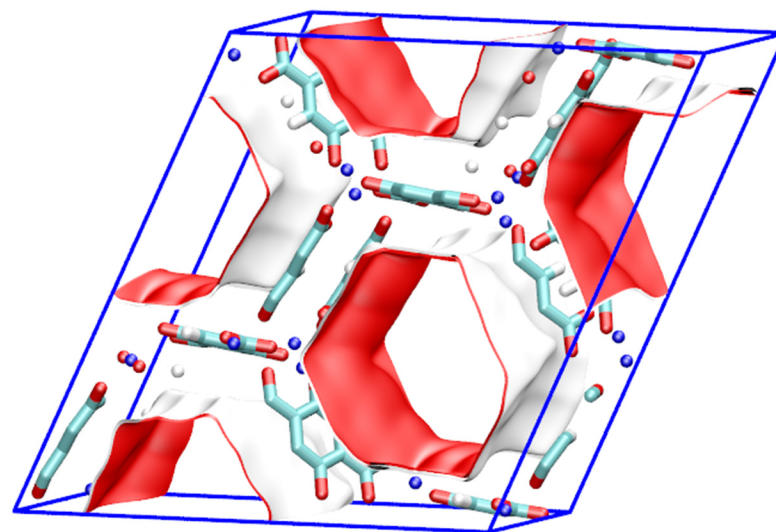
# NiMOF-74

# NiMOF-74 pore landscapes

Figure 34

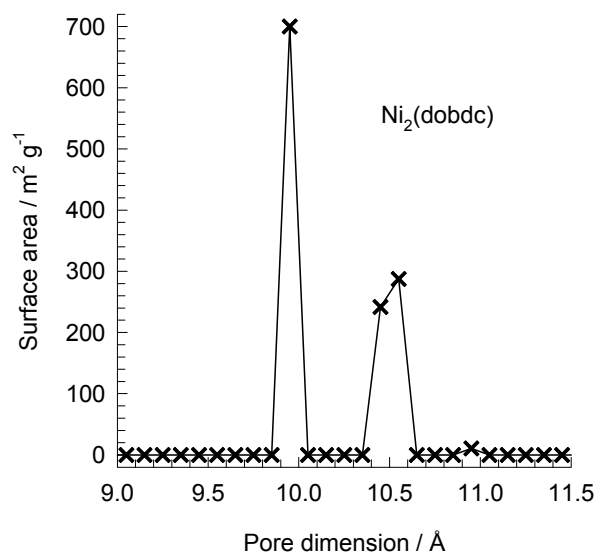
The structural information on NiMOF-74 ( $= \text{Ni}_2(\text{dobdc}) = \text{Ni}(\text{dobdc}) = \text{CPO-27-Ni}$ ) with  $\text{dobdc} = (\text{dobdc}^{4-} = 2,5\text{-dioxido-1,4-benzenedicarboxylate})$  were obtained from

- A.Ö. Yazaydın, R.Q. Snurr, T.H. Park, K. Koh, J. Liu, M.D. LeVan, A.I. Benin, P. Jakubczak, M. Lanuza, D.B. Galloway, J.J. Low, R.R. Willis, Screening of Metal-Organic Frameworks for Carbon Dioxide Capture from Flue Gas using a Combined Experimental and Modeling Approach, *J. Am. Chem. Soc.* 131 (2009) 18198-18199.
- D. Britt, H. Furukawa, B. Wang, T.G. Glover, O.M. Yaghi, Highly efficient separation of carbon dioxide by a metal-organic framework replete with open metal sites, *Proc. Natl. Acad. Sci. U.S.A.* 106 (2009) 20637-20640.
- N.L. Rosi, J. Kim, M. Eddaoudi, B. Chen, M. O'Keeffe, O.M. Yaghi, Rod Packings and Metal-Organic Frameworks Constructed from Rod-Shaped Secondary Building Units, *J. Am. Chem. Soc.* 127 (2005) 1504-1518.
- P.D.C. Dietzel, B. Panella, M. Hirscher, R. Blom, H. Fjellvåg, Hydrogen adsorption in a nickel based coordination polymer with open metal sites in the cylindrical cavities of the desolvated framework, *Chem. Commun.* (2006) 959-961.
- P.D.C. Dietzel, V. Besikiotis, R. Blom, Application of metal-organic frameworks with coordinatively unsaturated metal sites in storage and separation of methane and carbon dioxide, *J. Mater. Chem.* 19 (2009) 7362-7370.
- S.R. Caskey, A.G. Wong-Foy, A.J. Matzger, Dramatic Tuning of Carbon Dioxide Uptake via Metal Substitution in a Coordination Polymer with Cylindrical Pores, *J. Am. Chem. Soc.* 130 (2008) 10870-10871.



# NiMOF-74 pore dimensions

Figure 35



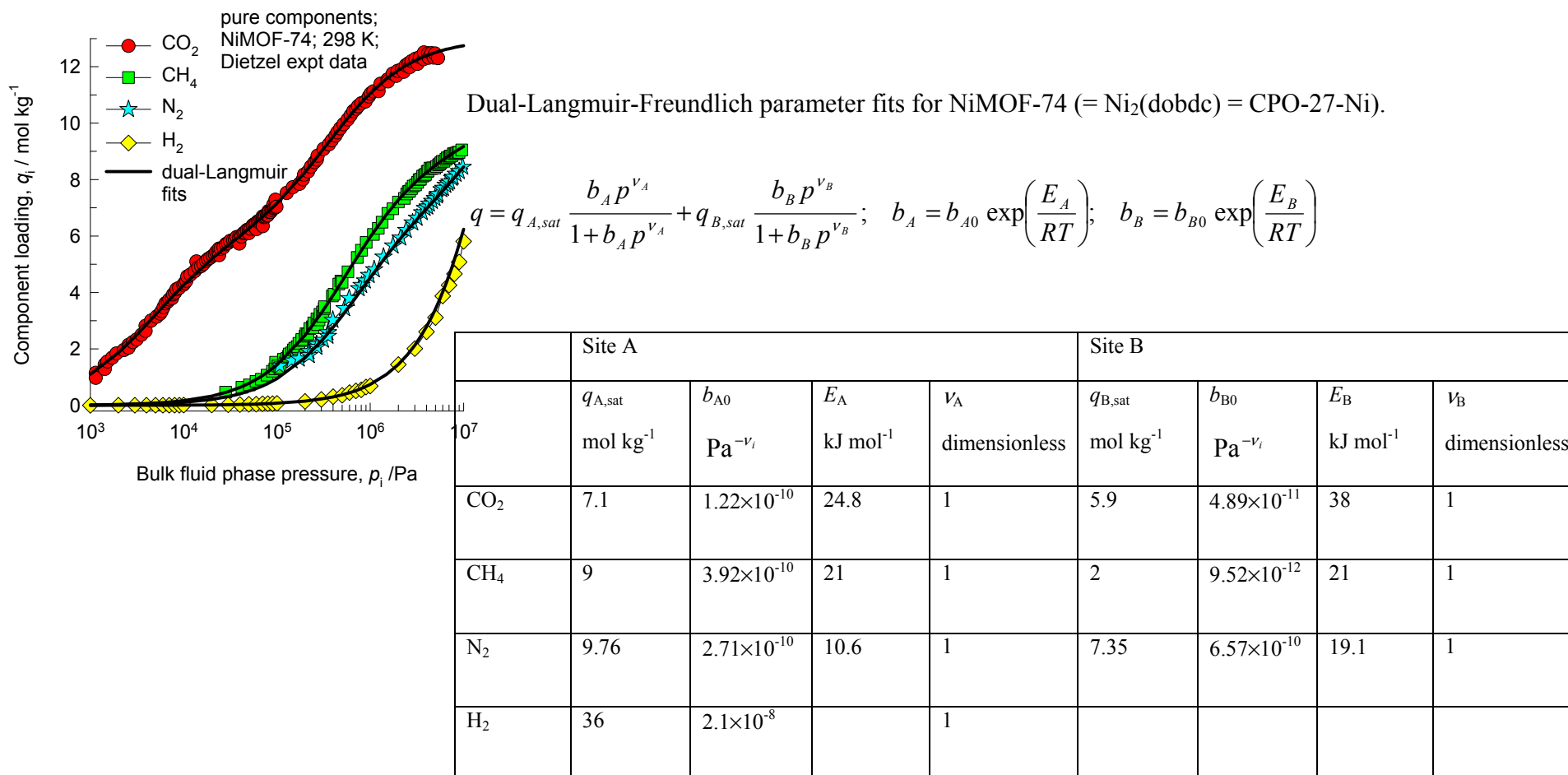
This plot of surface area versus pore dimension is determined using a combination of the DeLaunay triangulation method for pore dimension determination, and the procedure of Düren for determination of the surface area.

	NiMOF-74
$a / \text{Å}$	25.7856
$b / \text{Å}$	25.7856
$c / \text{Å}$	6.7701
Cell volume / $\text{Å}^3$	3898.344
conversion factor for [molec/uc] to [mol per kg Framework]	0.3568
conversion factor for [molec/uc] to [kmol/m <sup>3</sup> ]	0.6133
$\rho$ [kg/m <sup>3</sup> ]	1193.811
MW unit cell [g/mol/framework]	2802.592
$\phi$ , fractional pore volume	0.695
open space / $\text{Å}^3/\text{uc}$	2707.6
Pore volume / cm <sup>3</sup> /g	0.582
Surface area / m <sup>2</sup> /g	1239.0
DeLaunay diameter / $\text{Å}$	9.80



# NiMOF-74 fits of experimental isotherms for CO<sub>2</sub>, CH<sub>4</sub>, N<sub>2</sub>, and H<sub>2</sub>

Figure 36



The isotherm fit parameters for CO<sub>2</sub>, CH<sub>4</sub>, and N<sub>2</sub> are on the basis of experimental data for variety of temperatures from:

P.D.C. Dietzel, V. Besikiotis, R. Blom, Application of metal–organic frameworks with coordinatively unsaturated metal sites in storage and separation of methane and carbon dioxide, *J. Mater. Chem.* 19 (2009) 7362-7370.

The isotherm fit parameters for H<sub>2</sub> are on the basis of experimental data from:

Yaghi, O. M. Hydrogen Storage in Metal Organic Frameworks, [www.hydrogen.energy.gov/pdfs/review11/st049\\_yaghi\\_2011\\_p.pdf](http://www.hydrogen.energy.gov/pdfs/review11/st049_yaghi_2011_p.pdf), University of California Los Angeles, California, 2011.

These H<sub>2</sub> isotherm data were supplemented with CBMC simulations.

# NiMOF-74: Analysis of unary permeance data of Lee et al. (2012)

Figure 37



The data on membrane permeances are obtained from:

Lee, D. J.; Li, Q.; Kim, H.; Lee, K. Preparation of Ni-MOF-74 membrane for CO<sub>2</sub> separation by layer-by-layer seeding technique, *Microporous Mesoporous Mater.* **2012**, *163*, 169-177.

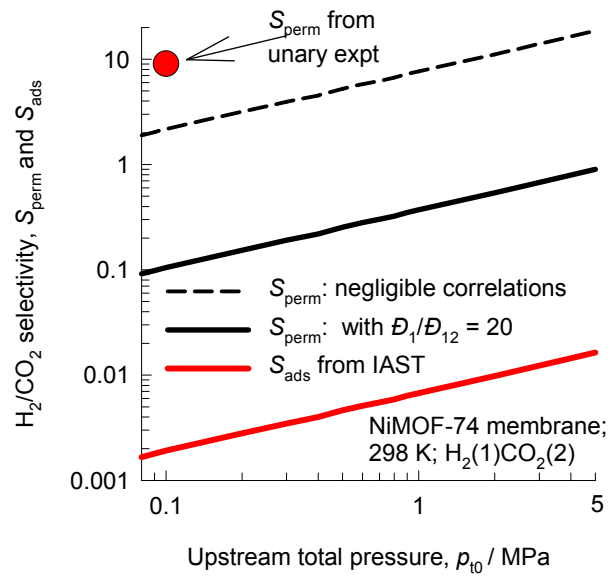
The transport coefficients are calculated from

$$\frac{\rho D_i}{\delta} = \frac{\Pi_i}{\Gamma_i \frac{q_i}{p_i}}$$

with the molar loadings  $q_i$  and thermodynamic correction factor  $\Gamma_i$  at the upstream face determined from pure component isotherm fits evaluated at  $p_i = 100 \text{ kPa}$

# NiMOF-74: Maxwell-Stefan model calculations for H<sub>2</sub>/CO<sub>2</sub> mixture permeation

Figure 38



The inputs for the Maxwell-Stefan model are:

$$\frac{\rho D_1}{\delta} = 17 \text{ kg m}^{-2} \text{ s}^{-1}$$

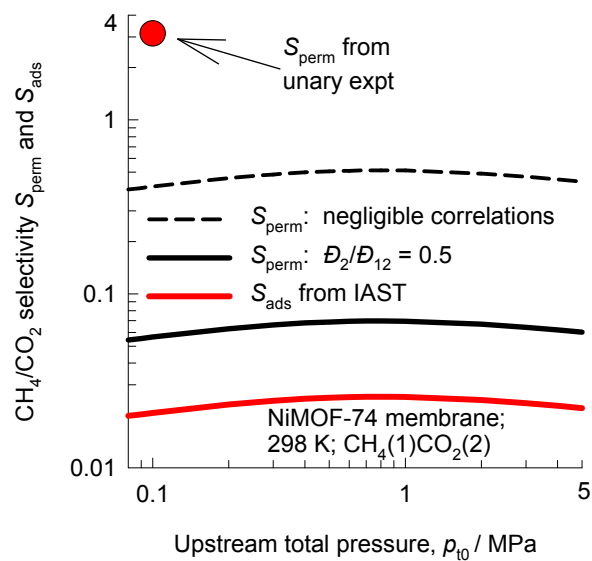
$$\frac{\rho D_2}{\delta} = 0.015 \text{ kg m}^{-2} \text{ s}^{-1}$$

$$\frac{D_1}{D_{12}} = 20$$

$$\frac{D_2}{D_{12}} = 0.17647$$

# NiMOF-74: Maxwell-Stefan model calculations for CH<sub>4</sub>/CO<sub>2</sub> mixture permeation

Figure 39



The inputs for the Maxwell-Stefan model are:

$$\frac{\rho D_1}{\delta} = 0.3 \text{ kg m}^{-2} \text{ s}^{-1}$$

$$\frac{\rho D_2}{\delta} = 0.015 \text{ kg m}^{-2} \text{ s}^{-1}$$

$$\frac{D_1}{D_{12}} = 10$$

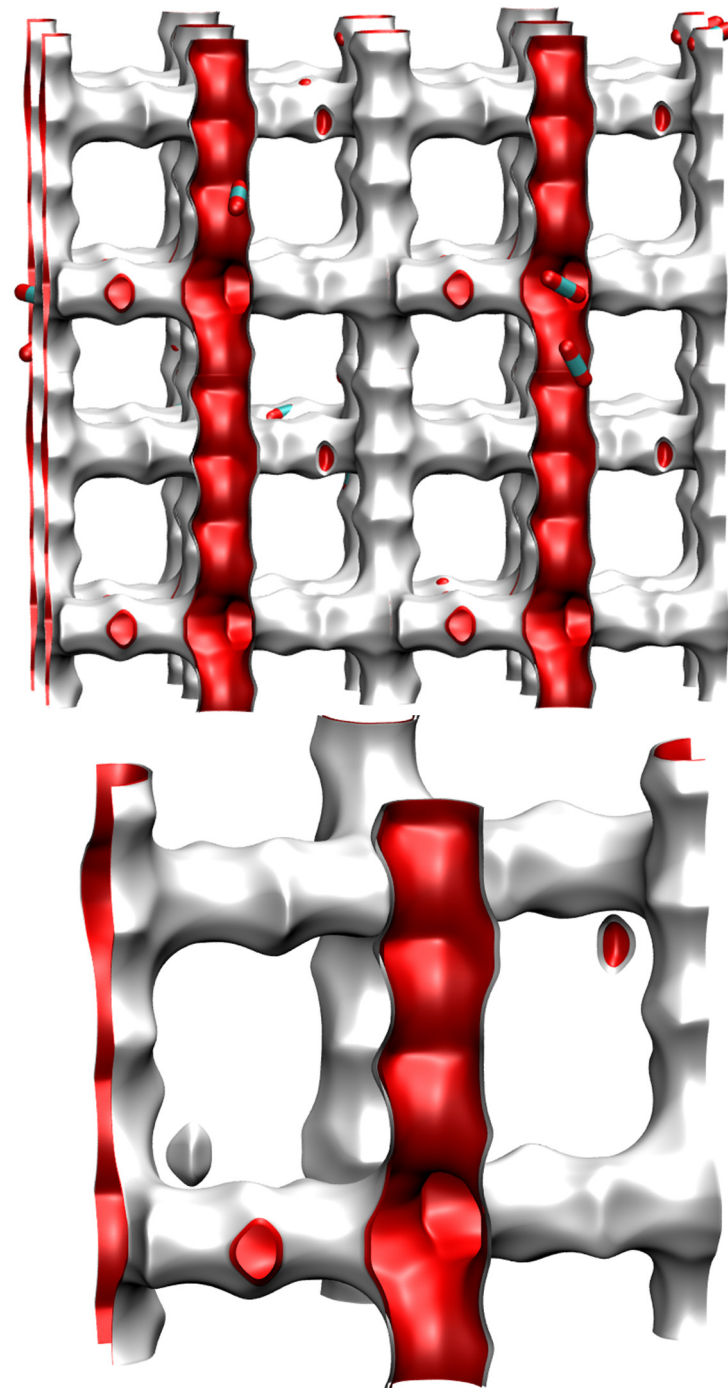
$$\frac{D_2}{D_{12}} = 0.5$$

**MFI**

# MFI pore landscape

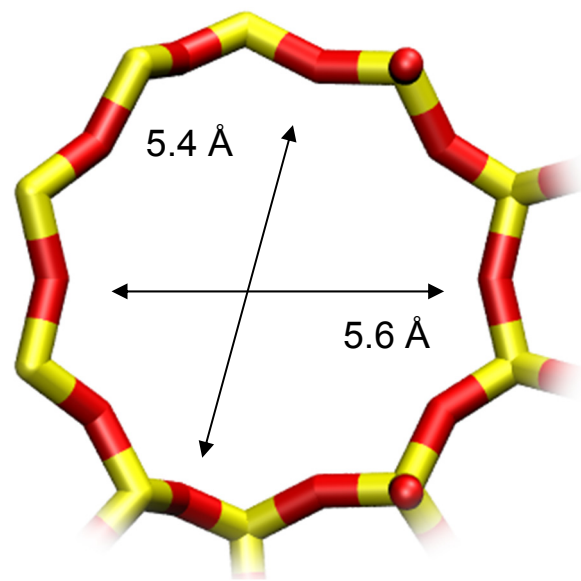
	MFI
$a / \text{\AA}$	20.022
$b / \text{\AA}$	19.899
$c / \text{\AA}$	13.383
Cell volume / $\text{\AA}^3$	5332.025
conversion factor for [molec/uc] to [mol per kg Framework]	0.1734
conversion factor for [molec/uc] to [kmol/m <sup>3</sup> ]	1.0477
$\rho$ [kg/m <sup>3</sup> ]	1796.386
MW unit cell [g/mol(framework)]	5768.141
$\phi$ , fractional pore volume	0.297
open space / $\text{\AA}^3/\text{uc}$	1584.9
Pore volume / cm <sup>3</sup> /g	0.165
Surface area / m <sup>2</sup> /g	487.0
DeLaunay diameter / $\text{\AA}$	5.16

Structural information from: C. Baerlocher, L.B. McCusker,  
Database of Zeolite Structures, International Zeolite Association,  
<http://www.iza-structure.org/databases/>

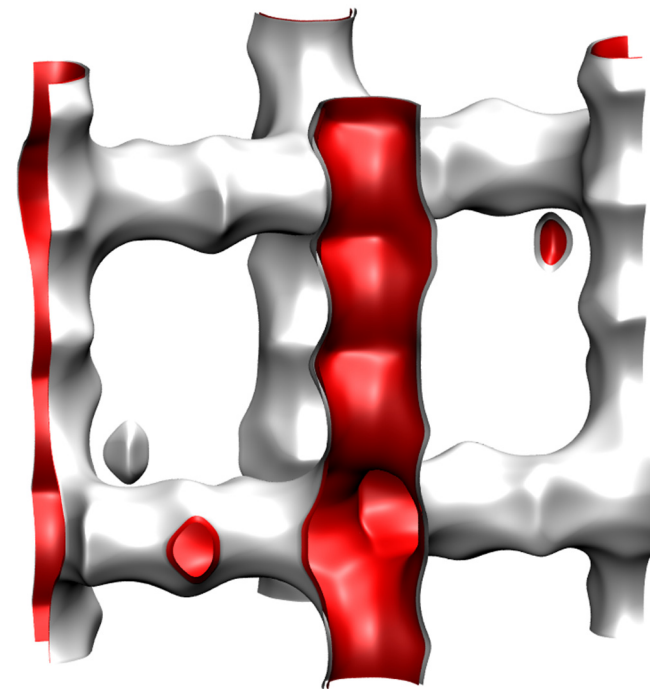
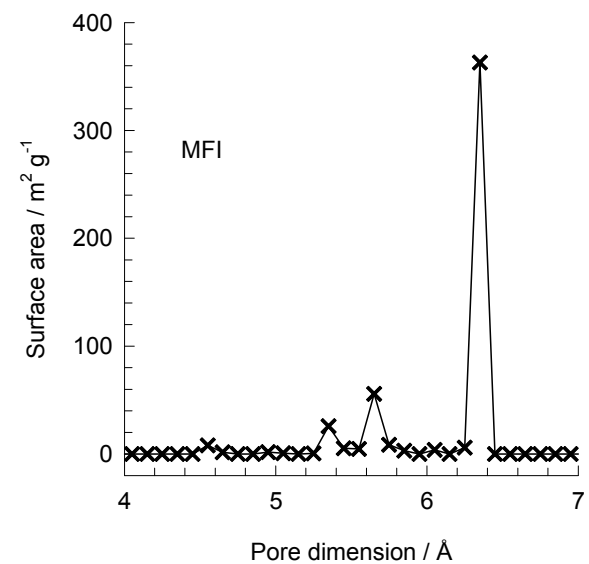
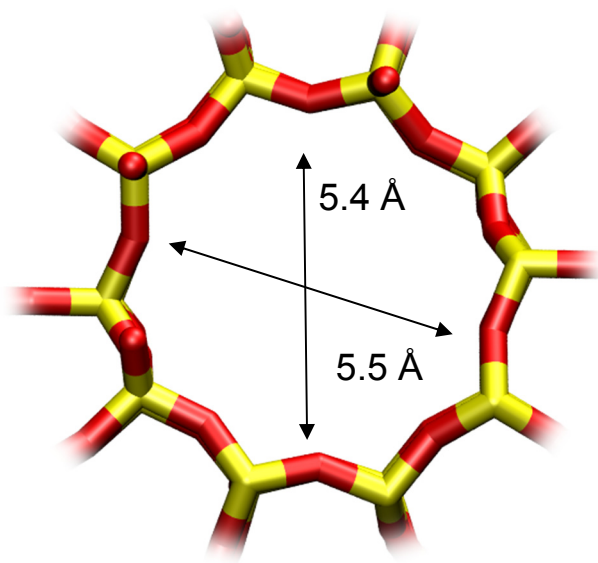


# MFI pore dimensions

10 ring channel  
of MFI viewed  
along [100]



10 ring channel  
of MFI viewed  
along [010]



# MFI pure component isotherm fits

Figure 43

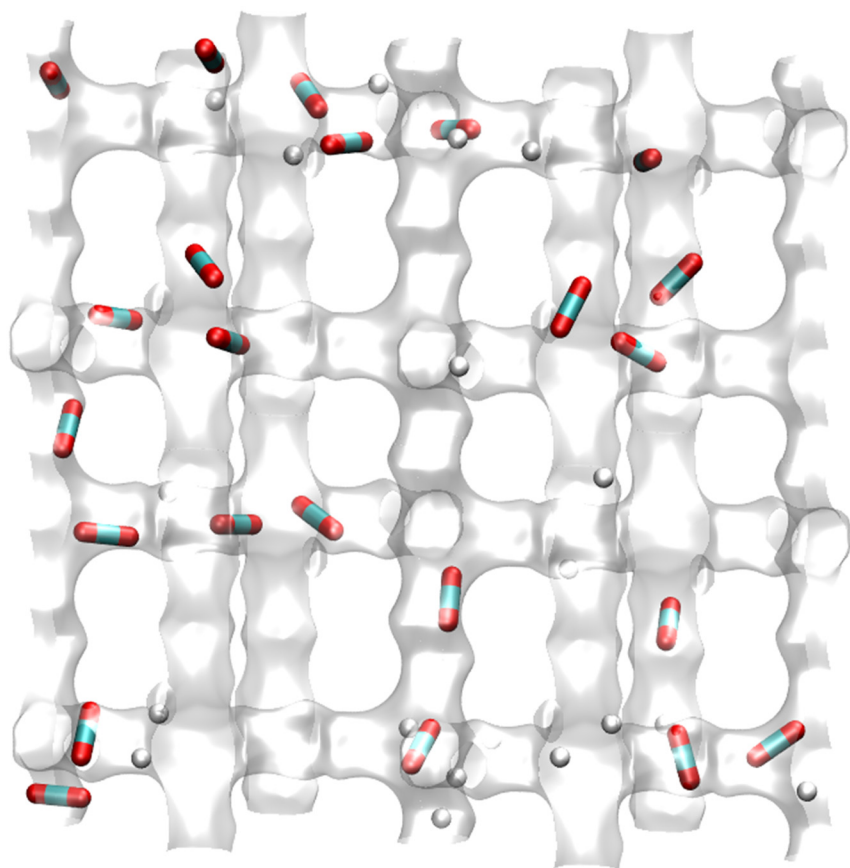
Dual-site Langmuir-Freundlich parameters for pure component isotherms in MFI at 300 K. Note that the saturation capacities are specified in molecules per unit cell. Multiply these by 0.173366 to obtain the values in mol per kg framework. All the isotherm data are obtained by Configurational-Bias Monte Carlo (CBMC) simulations. These have been verified to be in good agreement with all available experimental isotherm data.

	Site A			Site B		
	$\Theta_{i,A,sat}$ molecules uc <sup>-1</sup>	$b_{i,A}$ Pa <sup>-<math>\nu_i</math></sup>	$\nu_{i,A}$ dimensionless	$\Theta_{i,B,sat}$ molecules uc <sup>-1</sup>	$b_{i,B}$ Pa <sup>-<math>\nu_i</math></sup>	$\nu_{i,B}$ dimensionless
CO <sub>2</sub>	19	$6.12 \times 10^{-6}$	1	11	$1.73 \times 10^{-8}$	1
CH <sub>4</sub>	7	$5 \times 10^{-9}$	1	16	$3.1 \times 10^{-6}$	1
H <sub>2</sub>	30	$3.57 \times 10^{-8}$	1	42	$1.39 \times 10^{-9}$	1
C <sub>2</sub> H <sub>6</sub>	3.3	$4.08 \times 10^{-7}$	1	13	$7.74 \times 10^{-5}$	1
C <sub>3</sub> H <sub>8</sub>	1.4	$3.35 \times 10^{-4}$	0.67	10.7	$6.34 \times 10^{-4}$	1.06
nC <sub>4</sub> H <sub>10</sub>	1.5	$2.24 \times 10^{-3}$	0.57	8.7	$9.75 \times 10^{-3}$	1.12



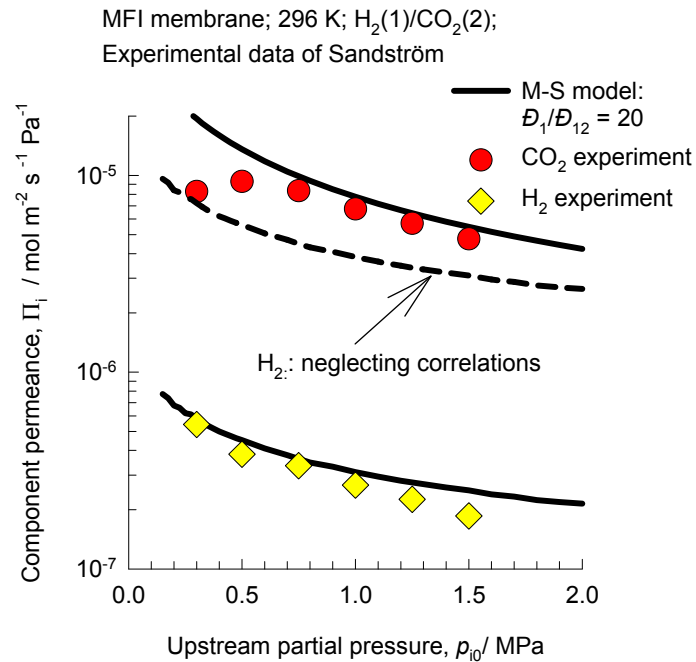
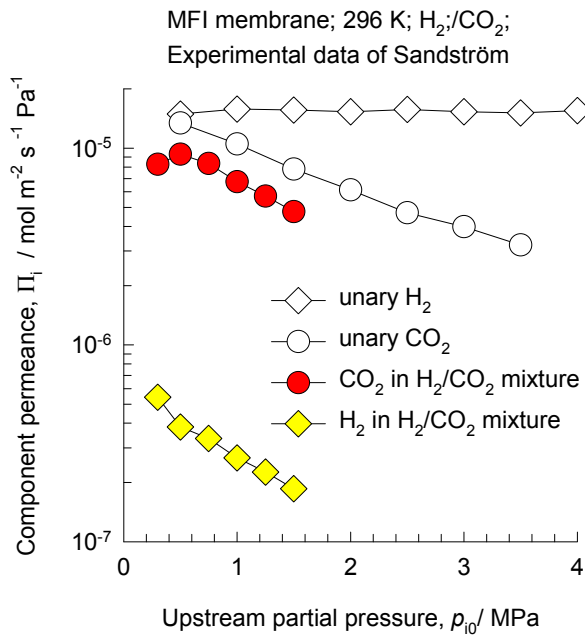
# MFI: Snapshots of H<sub>2</sub>/CO<sub>2</sub> mixture

Figure 44

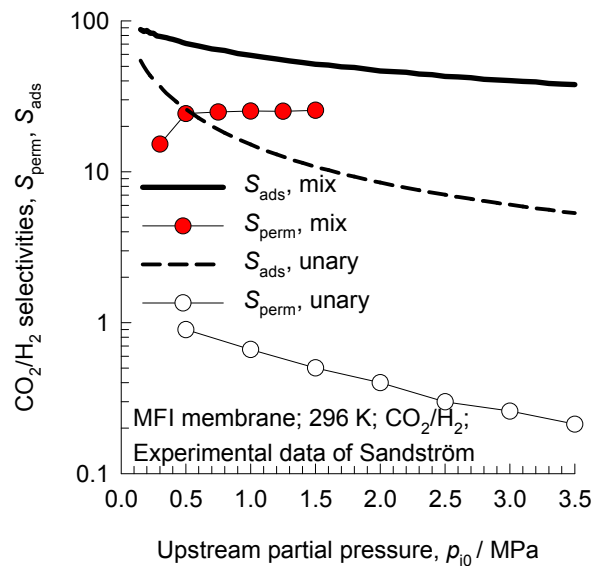


# MFI: Maxwell-Stefan model calculations for H<sub>2</sub>/CO<sub>2</sub> mixture permeation

Figure 45



The experimental data are re-plotted using the information contained in:  
L. Sandström, E. Sjöberg, J. Hedlund, Very high flux MFI membrane for CO<sub>2</sub> separation, J. Membr. Sci. 380 (2011) 232-240.



The inputs for the Maxwell-Stefan model are:

$$\frac{\rho D_1}{\delta} = 80 \text{ kg m}^{-2} \text{ s}^{-1}$$

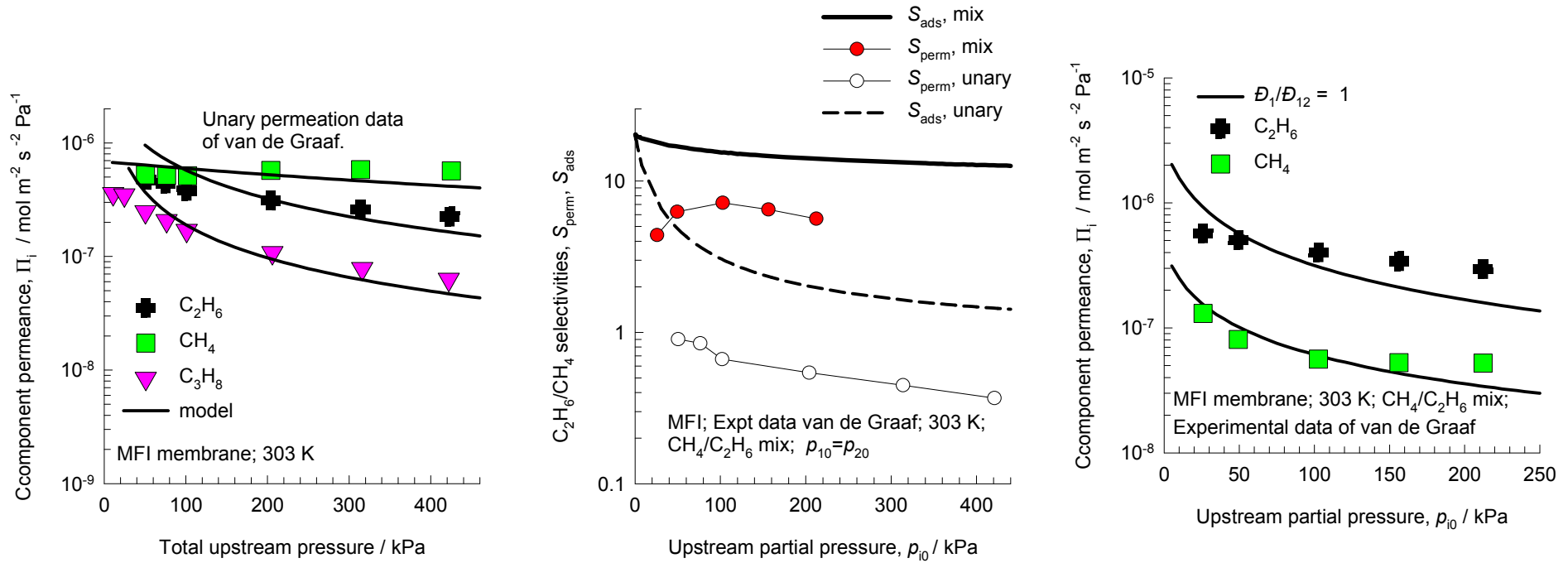
$$\frac{\rho D_2}{\delta} = 2.7 \text{ kg m}^{-2} \text{ s}^{-1}$$

$$\frac{D_1}{D_{12}} = 20$$

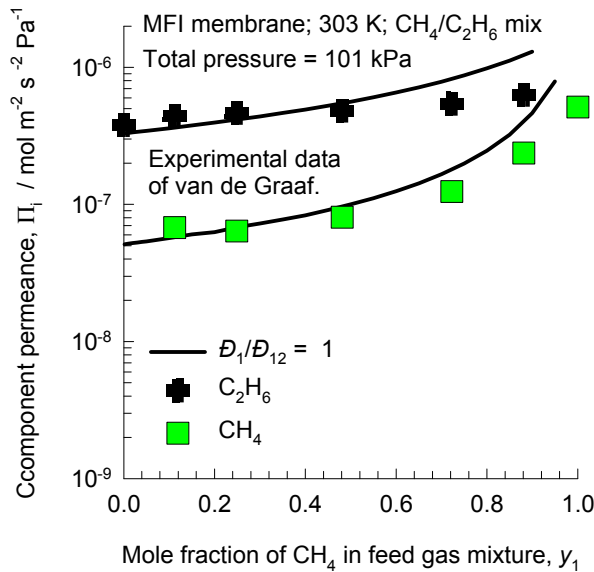
$$\frac{D_2}{D_{12}} = 0.675$$

# MFI: Maxwell-Stefan model calculations for CH<sub>4</sub>/C<sub>2</sub>H<sub>6</sub> mixture permeation

Figure 46



The experimental data are re-plotted using the information in: J.M. van de Graaf, F. Kapteijn, J.A. Moulijn, Modeling permeation of binary mixtures through zeolite membranes, A.I.Ch.E.J. 45 (1999) 497-511.



The inputs for the Maxwell-Stefan model are:

$$\frac{\rho D_1}{\delta} = 0.08 \text{ kg m}^{-2} \text{ s}^{-1}$$

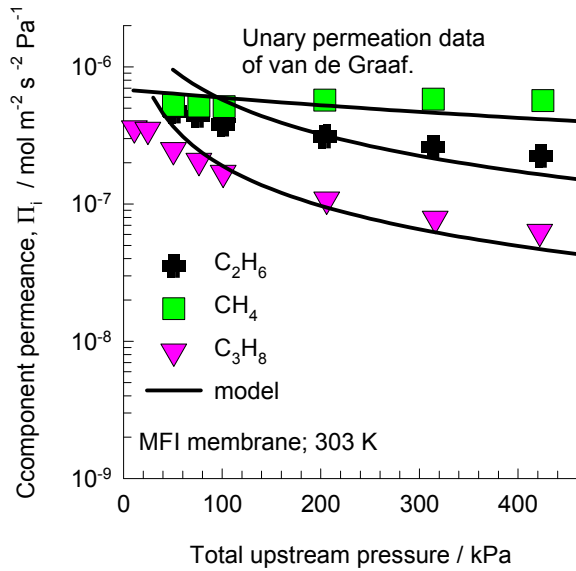
$$\frac{\rho D_2}{\delta} = 0.016 \text{ kg m}^{-2} \text{ s}^{-1}$$

$$\frac{D_1}{D_{12}} = 1$$

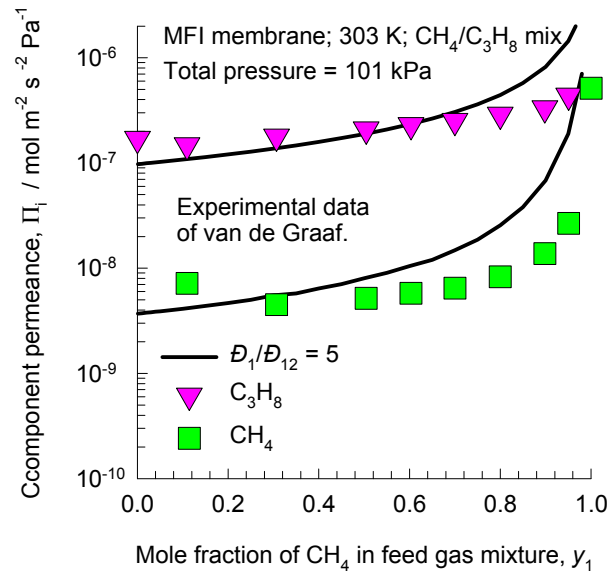
$$\frac{D_2}{D_{12}} = 0.2$$

# MFI: Maxwell-Stefan model calculations for CH<sub>4</sub>/C<sub>3</sub>H<sub>8</sub> mixture permeation

Figure 47



The experimental data are re-plotted using the information in: J.M. van de Graaf, F. Kapteijn, J.A. Moulijn, Modeling permeation of binary mixtures through zeolite membranes, A.I.Ch.E.J. 45 (1999) 497-511.



The inputs for the Maxwell-Stefan model are:

$$\frac{\rho D_1}{\delta} = 0.08 \text{ kg m}^{-2} \text{ s}^{-1}$$

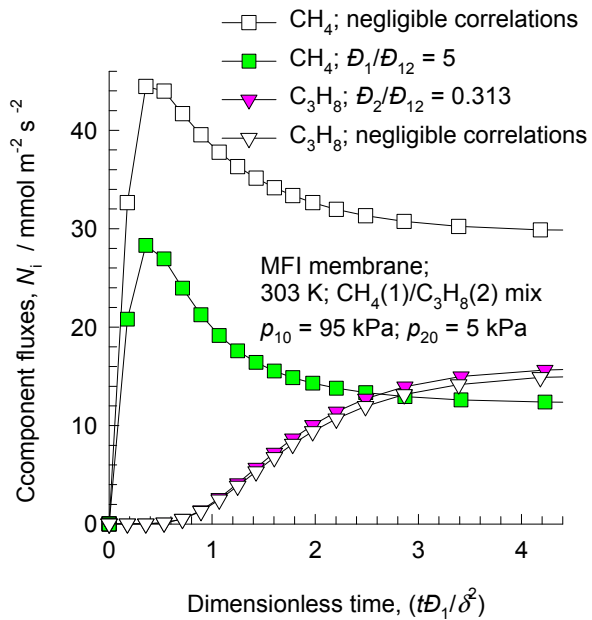
$$\frac{\rho D_2}{\delta} = 0.005 \text{ kg m}^{-2} \text{ s}^{-1}$$

$$\frac{D_1}{D_{12}} = 5$$

$$\frac{D_2}{D_{12}} = 0.313$$

# MFI: Maxwell-Stefan model calculations for transient CH<sub>4</sub>/C<sub>3</sub>H<sub>8</sub> mixture permeation

Figure 48



The inputs for the Maxwell-Stefan model are:

$$\frac{\rho D_1}{\delta} = 0.08 \text{ kg m}^{-2} \text{ s}^{-1}$$

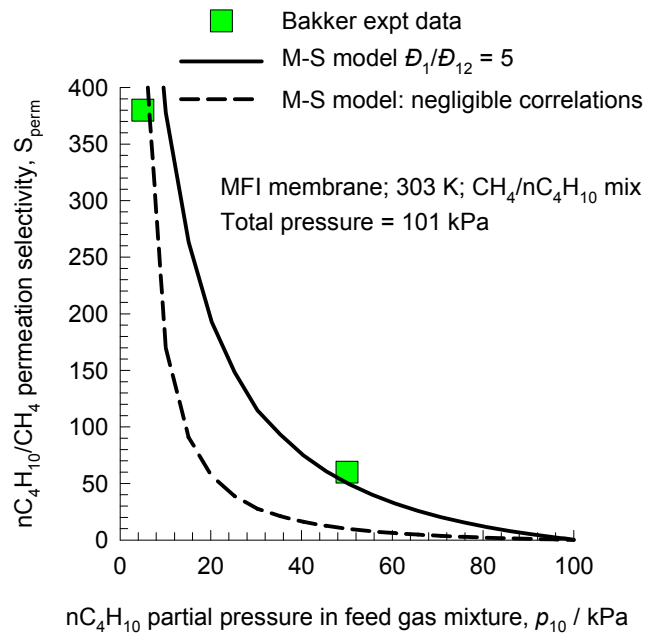
$$\frac{\rho D_2}{\delta} = 0.005 \text{ kg m}^{-2} \text{ s}^{-1}$$

$$\frac{D_1}{D_{12}} = 5$$

$$\frac{D_2}{D_{12}} = 0.313$$

# MFI: Maxwell-Stefan model calculations for CH<sub>4</sub>/nC<sub>4</sub>H<sub>10</sub> mixture permeation

Figure 49



The inputs for the Maxwell-Stefan model are:

The experimental data are re-plotted using the information in:  
W.J.W. Bakker, Structured systems in gas separation, Ph.D. Thesis, Delft University of Technology, Delft, 1999.

$$\frac{\rho D_1}{\delta} = 0.08 \text{ kg m}^{-2} \text{ s}^{-1}$$

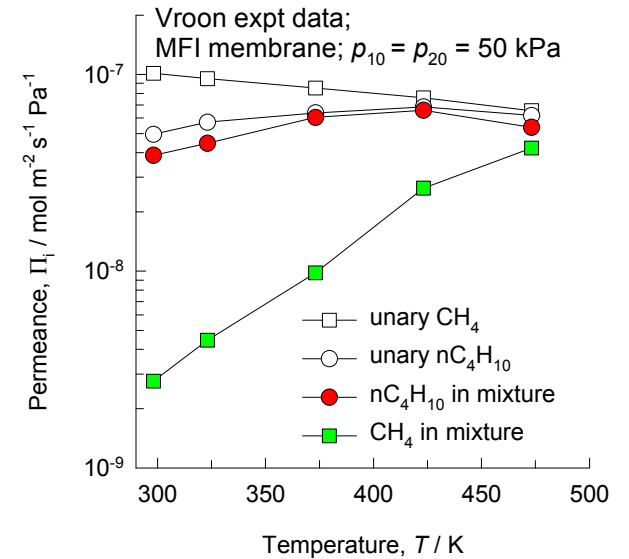
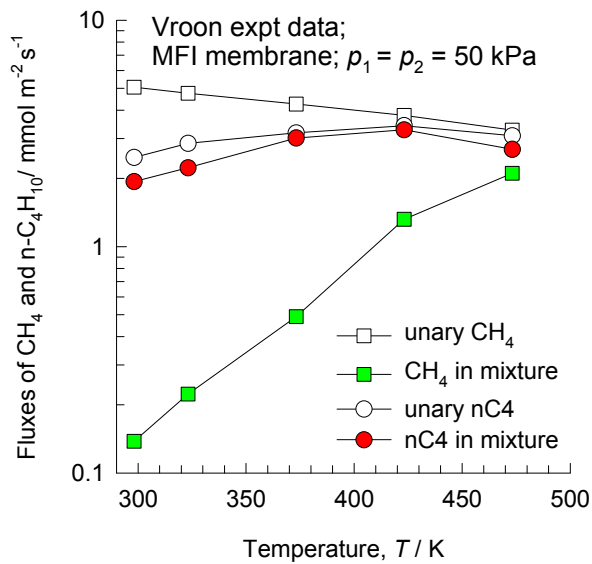
$$\frac{\rho D_2}{\delta} = 0.003 \text{ kg m}^{-2} \text{ s}^{-1}$$

$$\frac{D_1}{D_{12}} = 5$$

$$\frac{D_2}{D_{12}} = 0.188$$

# MFI $\text{CH}_4/\text{nC}_4\text{H}_{10}$ mixture permeation fluxes, compared with unary fluxes

Figure 50

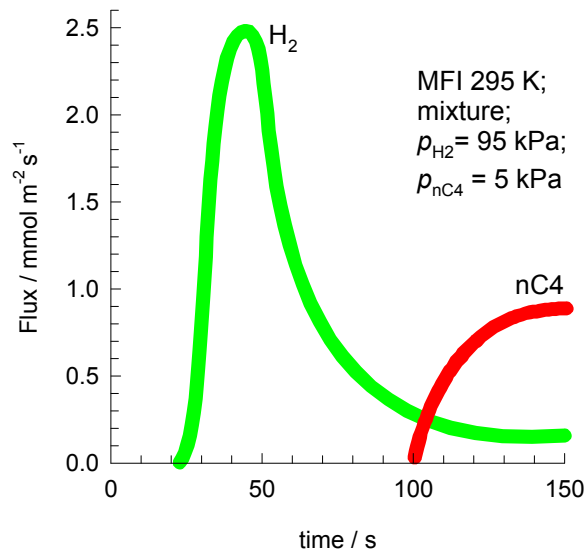


The experimental data are re-plotted using the information in:

Z.A.E.P. Vroon, K. Keizer, M.J. Gilde, H. Verweij, A.J. Burggraaf, Transport properties of alkanes through ceramic thin zeolite MFI membranes, *J. Membr. Sci.* 113 (1996) 293-300.

# MFI $nC_4H_{10}/H_2$ transient mixture permeation

Figure 51



The experimental data are re-plotted using the information contained in:

W.J.W. Bakker, Structured systems in gas separation, Ph.D. Thesis, Delft University of Technology, Delft, 1999.

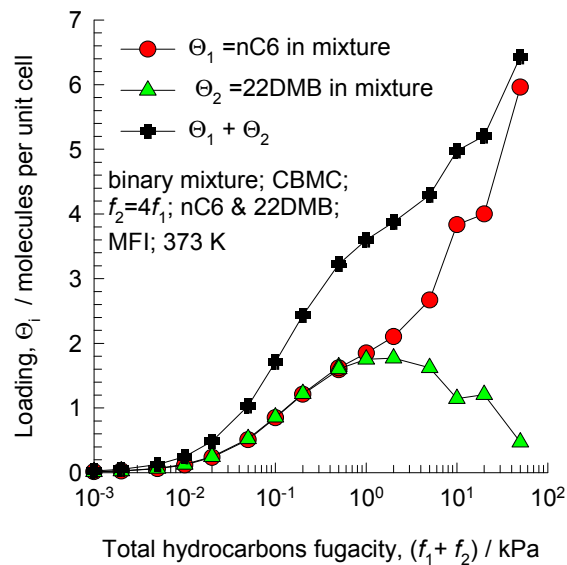
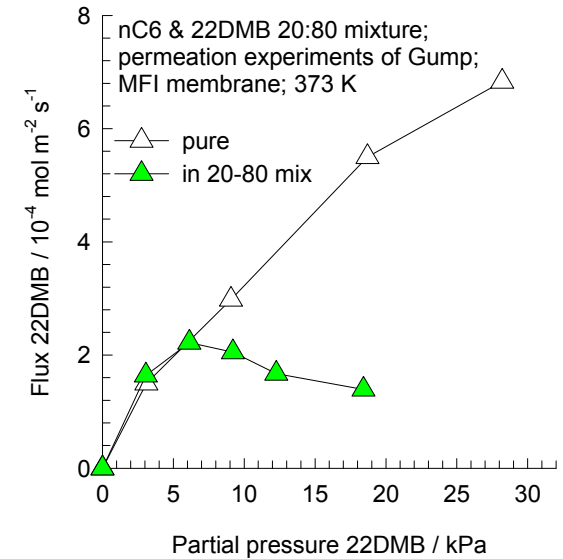
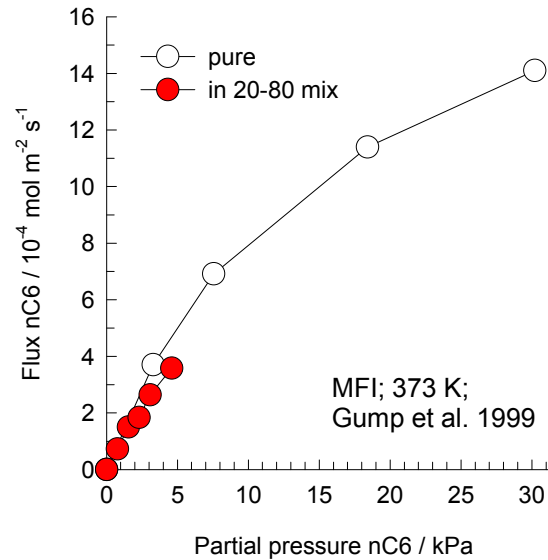
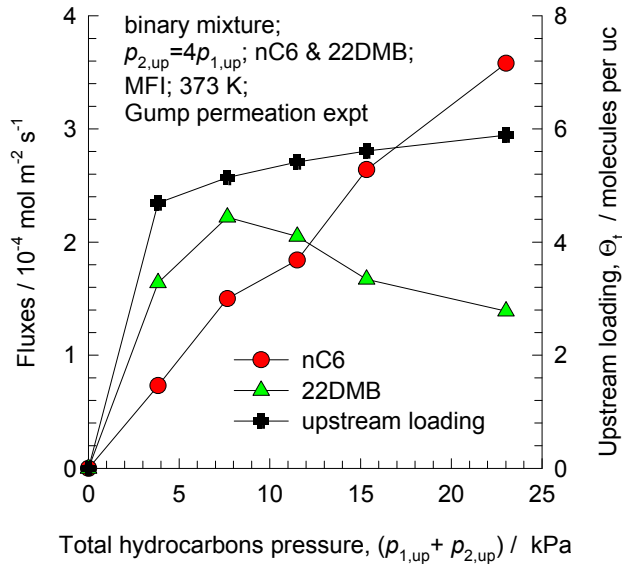
The theoretical analysis of this transient mixture permeation behavior is provided in our earlier works:

R. Krishna, R. Baur, Modelling issues in zeolite based separation processes, Sep. Purif. Technol. 33 (2003) 213-254.



# MFI nC6/22DMB mixture permeation

Figure 52



The experimental data are re-plotted using the information contained in:  
 C.J. Gump, R.D. Noble, J.L. Falconer, Separation of hexane isomers through nonzeolite pores in ZSM-5 zeolite membranes, *Ind. Eng. Chem. Res.* 38 (1999) 2775-2781.

The experimentally observed maximum in the flux of 22DMB is attributable to a corresponding maximum in the component loading as observed in CBMC simulations of mixture adsorption (shown on left). Further details of data interpretation is available in our previous works:

R. Krishna, J.M. van Baten, Screening of zeolite adsorbents for separation of hexane isomers: A molecular simulation study, *Sep. Purif. Technol.* 55 (2007) 246-255.

R. Krishna, Describing the diffusion of guest molecules inside porous structures, *J. Phys. Chem. C* 113 (2009) 19756-19781.

Figure 53

# Snapshots showing the location of nC6/22DMB within MFI

