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Shorter Communication

Investigating the validity of the Bosanquet formula for estimation of diffusivities in mesopores

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article info

ABSTRACT

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Keywords: Adsorption Self-diffusivity Mesopore Knudsen formula Dusty gas model Bosanquet interpolation Molecular Dynamics (MD) simulations were performed to determine the self-diffusivity, $D_{i,self}$, of H₂, CO₂, Ar, Kr, CH₄, C₂H₆, C₃H₈, and nC₄H₁₀ for a variety of pore concentrations within one-dimensional mesopores of 2 nm, 3 nm, 3.4 nm, 4 nm, 5.8 nm, 7.6 nm and 10 nm sizes. The MD simulated values of D_{iself} are compared with estimations of the commonly used Bosanquet formula $1/D_{i, self} = 1/D_{i, Kn} + 1/D_{i_ift}$, that combines molecule–wall and molecule–molecule interactions:, where $D_{i,Kn}$ is the Knudsen diffusivity, and $D_{ii,fl}$ is the self-diffusivity of species i in the fluid phase at the same molar concentration, c_i , as within the mesopores. For components with poor adsorption strength, such as H_2 , the MD simulated $D_{i, self}$ values are in good agreement with the estimations using the Bosanquet formula for the whole range of pore concentrations. For components with strong adsorption at the pore walls, the MD simulated values are significantly lower than the Bosanquet estimations when molecule–wall interactions are dominant. These deviations are traceable to the failure of the Knudsen prescription of diffuse reflectance on molecule–wall collisions, because adsorption at the pore walls introduces a bias in the molecular hops. For any given molecule, the Bosanquet estimations tend to be increasingly accurate when the pore diameters are increased.

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1. Introduction

The Bosanquet equation for self-diffusivity of species i for combining molecule–wall and molecule–molecule interactions

$$
1/D_{i,self} = 1/D_{i, Kn} + 1/D_{ii, fl}
$$
 (1)

was developed in a classified report dated September 27, 1944, and this interpolation formula only became known when it was later cited by [Pollard and Present \(1948\)](#page-4-0) in a classic paper providing a rigorous foundation to Eq. (1) using arguments based on the kinetic theory of gases.

The molecule–wall interactions are quantified by the formula that was first put forward by [Knudsen \(1909\),](#page-4-0) and subsequently refined by [Smoluchowski \(1910\):](#page-4-0)

$$
D_{i,Kn} = \frac{d_p}{3} \sqrt{\frac{8RT}{\pi M_i}}
$$
 (2)

The molecule–molecule interactions are characterized by the self-diffusivity in the *fluid* phase, $D_{i i f l}$. For fluids with low molar densities, c_i , the D_{iifl} can be estimated using a variety of methods ([Poling et al., 2001](#page-4-0)). One commonly used estimation method is that of Fuller, Schettler and Giddings (FSG) [\(Fuller et al., 1966\)](#page-4-0), developed for binary mixtures; $D_{ii,fl}$ can be determined by taking both species to be the identical to each other.

The Bosanquet formula is firmly entrenched in the chemical engineering literature and is widely used in practice to calculate $D_{i, self}$ in meso- and macro-porous materials [\(Spiegler, 1966;](#page-4-0) [Youngquist, 1970\)](#page-4-0), most notably in the calculation of catalyst effectiveness factors [\(Froment and Bischoff, 1979](#page-4-0)). The Bosanquet equation (1) can also be derived as a limiting case of Dusty Gas Model [\(Mason and Malinauskas, 1983](#page-4-0)) for binary mixture diffusion by taking the species in the mixture to be the identical to each other.

Our primary objective in this communication is to investigate the validity of Eq. (1) by carrying out Molecular Dynamics (MD) simulations of diffusivities in one-dimensional mesoporous channels with well-defined sizes. Specifically, MD simulations were performed to determine $D_{i,self}$ of H₂, CO₂, Ar, Kr, CH₄, C₂H₆, C₃H₈, and nC_4H_{10} for a variety of pore concentrations within onedimensional mesopores of 2 nm, 3 nm, 3.4 nm, 4 nm, 5.8 nm, 7.6 nm and 10 nm sizes. We aim to demonstrate that the Bosanquet formula severely over-estimates the self-diffusivities when the species i has a strong adsorption strength, and molecule–wall interactions are dominant.

For convenience and easy reference, the molecular simulation methodology, specification of force fields, and simulation data on

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diffusivities are provided in the Supplementary Material accompanying this publication.

2. Results and discussion

*Ð*_{ii, fl} / 10⁻⁸ m² s⁻¹

a

10

100

1000

2000

Firstly, consider the MD simulations of the self-diffusivity of a variety of species in the *fluid* phase, ${D}_{i i, f l}$, for a range of molar concentrations, c_i (see Figs. 1a and b). The fluid phase diffusivities decrease with increasing c_i . The decrease of \boldsymbol{D}_{iifl} with increasing c_i is linear at low concentrations, and the MD simulated values are in excellent agreement with estimations using the FSG method ([Fuller et al., 1966](#page-4-0)), also plotted in Fig. 1as continuous solid lines. For values of c_i > 10 kmol m⁻³, the $D_{ii,fl}$ predictions of the FSG method are not reliable. Similar agreement is obtained for the MD simulations of the fluid phase Maxwell–Stefan diffusivity for *binary* mixtures, ${D_{12,fl}}$, and the estimations of the FSG method (see Figs. 1c and d).

Let us now consider self-diffusivities $D_{i, self}$ within the 3.4 nm channels of the covalent organic framework BTP-COF [\(Krishna](#page-4-0) [and van Baten, 2011a](#page-4-0)); the MD simulated diffusivities of a variety of guest molecules are shown in [Fig. 2.](#page-2-0) For the Bosanquet

> $H₂$ Ar

calculations presented in [Fig. 2](#page-2-0) we use the MD simulated $D_{ii,fl}$, at the same molar concentrations as within the pores, rather than the FSG calculated ones, in order to ensure applicability for the entire range of pore concentrations. For H_2 , that has poor adsorption strength, the MD simulated values of $D_{i,self}$ are in reasonably good agreement with the calculations using Eq. (1) for the entire range of pore concentrations. For $CO₂$, Ar, CH₄, C₂H₆, and C_3H_8 the Bosanquet formula severely over-estimates the selfdiffusivities, when molecule–wall interactions are dominant, i.e. at low pore concentrations. The reasons for this departure must be attributed to the failure of the Knudsen formula for molecules with strong adsorption strength; this failure has been highlighted in several recent publications [\(Bhatia et al., 2011;](#page-4-0) [Krishna, 2009;](#page-4-0) [Krishna and van Baten, 2009a,](#page-4-0) [2009b](#page-4-0), [2011a,](#page-4-0) [2011b,](#page-4-0) [2011c\)](#page-4-0). Adsorption causes the molecules to bind to the wall, and perhaps hop to a neighboring adsorption site, rather than return to the bulk after collision; this introduces a bias in the molecular hops. This bias increases with increasing adsorption strength, causing a violation of the diffuse reflectance assumption that is invoked in deriving Eq. (2) [\(Bhatia et al., 2011;](#page-4-0) [Krishna, 2009;](#page-4-0) [Krishna and](#page-4-0) [van Baten, 2011c\)](#page-4-0). It has been demonstrated that the departures from the Knudsen formula correlates with the binding energy for

 $P_{\rm{li.\,fi}}$ / 10⁻⁸ m² s⁻¹

b

10

 $CH₄$ C_2H_6

100

300

Fig. 1. (a, b) MD simulations of fluid phase unary self-diffusivities $D_{ii,f}$ of (a) H₂, CO₂, Ar, and (b) CH4, C2H₆, C3H₈ at 300 K as a function of the fluids concentration c_i (c, d) MD simulations of fluid phase binary Maxwell–Stefan diffusivities $D_{12,~\it fl}$ of (c) CH4/H2, CH4/C2H6, CH4/C3H8, and (d) CO2/H2, CO2/CH4, CO2/Ne mixtures at 300 K as a function of the total fluids concentration c_t . The solid continuous lines are the estimations of D_{iifl} and D_{12fl} using the correlation of Fuller, Schettler and Giddings [\(Fuller](#page-4-0) [et al., 1966\)](#page-4-0).

Fig. 2. MD simulations of self-diffusivities $D_{i,\text{self}}$ of (a) H₂, (b) CO₂, (c) Ar, (d) CH₄, (e) C₂H₆, and (f) C₃H₈ at 300 K in BTP-COF, a covalent organic framework with 3.4 nm pore size, as a function of the pore concentration c_i . The MD simulations (circles) are compared with the estimations using the Bosanquet and Knudsen formulae.

adsorption of the molecules at the pore walls ([Krishna and van](#page-4-0) [Baten, 2011a\)](#page-4-0).

Results analogous to those presented in Fig. 2 apply to all guest–host combinations investigated; the data are available in the Supplementary Material accompanying this publication.

The curvature of the surfaces of the pore walls should be expected to have an influence on the failure of the diffuse reflectance prescription of Knudsen. We should anticipate that with increasing pore diameters, i.e. with decreasing curvature of the walls, the bias in the hops introduced by adsorption would be lessened. In order to confirm this expectation, let us compare the self-diffusivities of CH_4 in cylindrical silica pores of increasing

diameters (see [Fig. 3](#page-3-0)). We note that the degree of over-estimation of the self-diffusivities using the Bosanquet formula decreases as the pore diameter increases. For quantification of this finding, we collected the data on the self-diffusivities of Ar and $CH₄$ at a fixed loading of c_i = 0.1 kmol m⁻³ in pores of different diameters. [Fig. 4](#page-3-0) presents a plot of the ratio of the self-diffusivities $D_{i, self}$ for Ar and $CH₄$ obtained from MD to that calculated by the Bosanquet formula plotted as a function of the pore diameter. As anticipated in the foregoing discussions, we observe that the departures of the Bosanquet formula from MD simulation results tends to decrease with increasing pore diameters. Since it has already been established that the molecule–molecule interaction term in

Fig. 3. MD simulations of self-diffusivities D_{iself} of CH₄ at 300 K in cylindrical silica pores of (a) 3 nm, (b) 5.8 nm, (c) 7.6 nm, and (d) 10 nm sizes, as a function of the pore concentration c_i . The MD simulations (circles) are compared with the estimations using the Bosanquet and Knudsen formulae .

Fig. 4. Ratio of the self-diffusivities $D_{i, self}$ for Ar and CH₄ (at c_i = 0.1 kmol m⁻³) in mesopores obtained from MD to that calculated by the Bosanquet formula plotted as a function of the pore diameter.

Eq. (1), is correctly estimated (cf. [Fig. 1\)](#page-1-0), the observed deviations in the Bosanquet formula are solely attributable to the shortcomings of the Knudsen prescription.

Bhatia and co-workers ([Bhatia, 2010](#page-4-0); [Bhatia and Nicholson,](#page-4-0) [2011\)](#page-4-0) have re-analyzed a variety of experimental data for diffusion in mesopores that confirm the conclusions drawn in the foregoing. Further experimental confirmation is provided below.

[Gruener and Huber \(2008\)](#page-4-0) have conducted careful experiments to determine the diffusivities of poorly adsorbing gases He and Ar in silicon nanochannels. As anticipated, the experimental data are in good agreement with the estimations using Knudsen prescription.

[Huang et al. \(2011\)](#page-4-0) have reported the diffusivities of n-heptane and toluene in SBA-15 materials of two different pore sizes. The molar masses of the two species are nearly equal, 0.1 and 0.092 kg mol $^{-1}$, respectively. However, the diffusivities of toluene are found to be significantly lower than that of n-heptane, because of its significantly higher adsorption strength [\(Krishna](#page-4-0) [and van Baten, 2011b\)](#page-4-0).

In the experimental study of [Katsanos et al. \(2005\)](#page-4-0), the measured diffusivities of nC5, nC6 and nC7 in α -alumina (d_p =21.6 nm) and γ -alumina (d_p = 10.6 nm) were found to be in the range of 10–27% of the values anticipated by the Knudsen formula.

[Goguet et al. \(2011\)](#page-4-0) have found that the interpretation of their TAP pulse response experiments, designed to be conducted within the Knudsen diffusion regime, required the inclusion of the influence of adsorption on molecule–wall interactions.

In his re-analysis of the diffusivity data of Reyes et al. (1997), Ruthven (2010) has concluded that the Knudsen formula holds with good accuracy. This conclusion is however based on the assumption of a uniform pore diameter of 7 nm for the silica sample used in the Reyes experiments. It remains to be ascertained whether the existence of a distribution of pore sizes of the silica sample will have a bearing on the conclusions drawn by Ruthven.

3. Conclusions

Molecular Dynamics simulations have been used to investigate the validity of the Bosanquet formula for calculation of the selfdiffusivities of a variety of guest molecules within one-dimensional channels in the 2 nm–10 nm range. For molecules, such as H_2 , that adsorb poorly, the Bosanquet formula is of reasonable accuracy for a wide range of pore concentrations. Significant departures of the Bosanquet formula are found with strongly adsorbing molecules at pore concentrations where molecule–wall collisions are dominant, already underlined in several recent investigations (Bhatia et al., 2011; Krishna and van Baten, 2009a, 2009b, 2011a, 2011b, 2011c). An important new finding pertains to the influence of the pore diameter. With increasing pore diameter, the bias in the molecular hops introduced by adsorption has a decreasing influence. As a consequence, the Bosanquet formula becomes increasingly accurate as the pore diameter increases.

The results of the current investigation also cast doubts on the validity of the Dusty Gas Model, when applied to mixtures in which at least one of the species has strong adsorption characteristics. As an example, the recent work on the modeling of intra-particle diffusion and reaction for steam reforming and methanol synthesis processes (Rout et al., 2008; Solsvik and Jakobsen, 2008) requires further scrutiny to examine the extent to which the failure of the Knudsen prescription influences the overall reactor performance.

Nomenclature

T temperature, K

Subscripts

i referring to component i

 f referring to fluid phase

Kn referring to Knudsen

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Appendix A. Supplementary Information

Supplementary data associated with this article can be found in the online version at <doi:10.1016/j.ces.2011.11.026>.

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