# **Environmental** Science & Technology

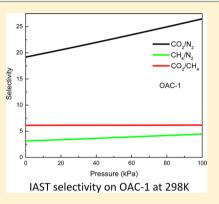
# Hydroquinone and Quinone-Grafted Porous Carbons for Highly Selective CO<sub>2</sub> Capture from Flue Gases and Natural Gas Upgrading

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**Supporting Information** 

**ABSTRACT:** Hydroquinone and quinone functional groups were grafted onto a hierarchical porous carbon framework via the Friedel–Crafts reaction to develop more efficient adsorbents for the selective capture and removal of carbon dioxide from flue gases and natural gas. The oxygen-doped porous carbons were characterized with scanning electron microscopy, transmission electron microscopy, X-ray powder diffraction, Fourier transform infrared spectroscopy, and Raman spectroscopy.  $CO_2$ ,  $CH_4$ , and  $N_2$  adsorption isotherms were measured and correlated with the Langmuir model. An ideal adsorbed solution theory (IAST) selectivity for the  $CO_2/N_2$  separation of 26.5 (298 K, 1 atm) was obtained on the hydroquinone-grafted carbon, which is 58.7% higher than that of the pristine porous carbon, and a  $CO_2/CH_4$  selectivity value of 4.6 (298 K, 1 atm) was obtained on the quinone-grafted carbon (OAC-2), which represents a 28.4% improvement over the pristine porous carbon. The highest  $CO_2$  adsorption capacity on the oxygen-doped carbon adsorbents is 3.46



mmol  $g^{-1}$  at 298 K and 1 atm. In addition, transient breakthrough simulations for  $CO_2/CH_4/N_2$  mixture separation were conducted to demonstrate the good separation performance of the oxygen-doped carbons in fixed bed adsorbers. Combining excellent adsorption separation properties and low heats of adsorption, the oxygen-doped carbons developed in this work appear to be very promising for flue gas treatment and natural gas upgrading.

#### **INTRODUCTION**

The continuously increasing  $CO_2$  level in the atmosphere has attracted the world's attention during last a few decades due to its potential environment damage. This motivates people to find various ways to capture  $CO_2$ , such as sequestrate carbon in the soil,<sup>1</sup> porous adsorbents,<sup>2,3</sup> and ocean fertilization.<sup>4</sup> Although a range of methods to capture and store  $CO_2$  has been investigated and developed, the liquid amine doped adsorbents are still the primary industrial approach.<sup>5,6</sup> However, this technology is not perfect due to some inherent drawbacks, such as amine degradation, evaporative losses, and high regeneration energy costs.<sup>7</sup> Low-quality natural gas such as biogas and land fill gas consists predominantly of CH<sub>4</sub> and has some impurities. The presence of  $CO_2$  and N<sub>2</sub> in natural gas could dramatically reduce the heating value and cause pipeline corrosion. Thus, separating impurities from CH<sub>4</sub> is required.

Porous materials have been proposed as good alternative adsorbents for the selective removal of carbon dioxide from flue gases and natural gas because they possess good chemical resistance and energy-efficient regeneration. Porous materials such as metal–organic frameworks (MOFs),<sup>8,9</sup> hyper-cross-linked polymers (HCPs),<sup>10,11</sup> covalent organic frameworks (COFs),<sup>12</sup> and amine-based cellulose<sup>13</sup> are being considered in industrial applications and are at various stages of development. Among those porous materials, activated carbons are one of the most studied materials. Activated carbons have been used for

many purposes such as wastewater treatment,<sup>14</sup> battery electrodes,<sup>15</sup> and supercapacitors.<sup>16</sup> Activated carbon continues to be a particularly intriguing alternative adsorbent because of its ease of synthesis, high CO<sub>2</sub> adsorption capacity, excellent selectivity and stability, and ease of regeneration.<sup>17–19</sup> To improve CO<sub>2</sub> capacity, we can either optimize the pore architecture by changing the synthetic parameters or modifying the surface properties through adding functional groups.<sup>20,21</sup> Secilla et al. demonstrated the importance of both microporosity and the presence of polar, basic functional groups in improving CO<sub>2</sub> adsorption.<sup>22</sup>

To improve the  $CO_2$  adsorption capacity and selectivity, incorporating functional groups into carbon materials is very important because they could enhance the affinity between adsorbate molecules and adsorbents. The common way is to introduce nitrogen-containing groups into the framework of activated carbon, due to its significant ability in improving interaction force. For example, Mahurin et al. demonstrated that an amidoxime-modified porous carbon displayed a  $CO_2$  capacity of 4.97 mmol g<sup>-1</sup> at 273 K and 1 atm, and enhanced  $CO_2/N_2$  selectivity.<sup>23</sup> Chandra et al. reported an N-doped

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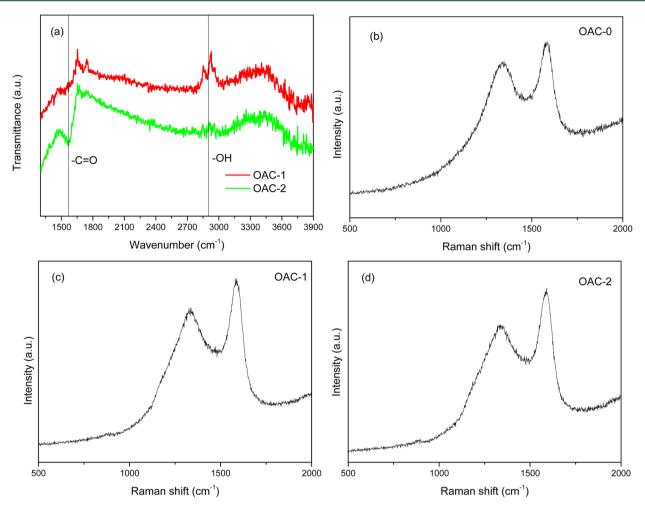


Figure 1. (a) FTIR spectra, (b) Raman pattern of OAC-0, (c) Raman pattern of OAC-1, and (d) Raman pattern of OAC-2.

carbon produced by the chemical activation of polypyrrole, which has a CO<sub>2</sub> capacity of 4.3 mmol  $g^{-1}$  at 298 K and 1 atm and good CO<sub>2</sub>/N<sub>2</sub> selectivty (approximately 34).<sup>24</sup> Although integrating nitrogen groups is the most studied method, other functional groups have been investigated recently; for example, Seema et al. reported on S-doped carbon and discovered a CO<sub>2</sub> capacity of 4.3 mmol g<sup>-1</sup> at 298 K and 1 atm.<sup>25</sup> Plaza et al. demonstrated the impact of oxidation on the CO<sub>2</sub> capacity of phenolic resin carbon and observed an enhancement in uptake in the oxidized carbon.<sup>26</sup> Thus, doping heteroatoms is an alternative method for enhancing CO<sub>2</sub> capture via electrostatics interactions. In 2013, Isikli et al. reported the influence of quinone grafting on carbon porous structure and supercapacitor performance.<sup>27</sup> The primary motivation of this work is to create porous carbon modified by hydroquinone and quinone groups and investigate the impact of the oxygen-containing groups on CO<sub>2</sub> adsorption and on the selective adsorption of gas mixture pairs containing CO<sub>2</sub>.

#### MATERIALS AND METHODS

Synthesis of OAC Material. Furan, dichloromethane (DCE), formaldehyde, anhydrous iron(III) chloride (FeCl<sub>3</sub>), methanol, 1,4,-benzoquinone (p-BQ), 2,3-dicloro-5,6-dicyano-1,4-benzoquinone (DDQ), and zinc chloride (ZnCl<sub>2</sub>) were purchased from Sigma-Aldrich. Hydrochloric acid (HCl) was purchased from Fisher Scientific. All chemicals were used as received without any further purification. To synthesize

polyfuran (PF), we dispersed 3 g of furan in 50 mL dichloromethane (DCE) at room temperature using sonication. A total of 8 g of FeCl<sub>3</sub> dissolved in 50 mL DCE was added to this solution. The resulting solution was stirred and allowed to polymerize for 24 h at room temperature. After polymerization, the precipitate was washed and dried at 80 °C for 24 h in a vacuum. The calcination of OAC-0 was carried out using a 1:1 ratio of PF and zinc chloride at 600 °C for 1 h under N2 atmosphere at 600  $^\circ\text{C}$  with a heating rate of 3  $^\circ\text{C}/\text{min}.$  After activation, the product was washed with HCl solution until a neutral pH was observed and then washed with DI water and dried under a vacuum at 80 °C for 24 h. Modified carbons were prepared following the reported process.<sup>27</sup> OAC-1 was prepared by a typical Friedel-Crafts reaction: OAC-0, p-BQ, FeCl<sub>3</sub>, and DCE were mixed into a round-bottom flask. The mixture was stirred vigorously at room temperature under nitrogen protection overnight. Part of the precipitation was collected by filtration, washed, dried under a vacuum at 80 °C for 24 h, and denoted as OAC-1. We then added DDQ into the mixture and stirred it for 3 h to complete the oxidation of the hydroquinone intermediate at a room temperature. The final product was washed and dried under a vacuum at 80 °C for 24 h and denoted as OAC-2.

**Material Characterization.** The XRD pattern was recorded using a Rigaku Miniflex-II X-ray diffractometer with Cu K $\alpha$  ( $\lambda$  = 1.5406 Å) radiation, a 30 kV/15 mA current, and a k $\beta$ -filter. Fourier transformed infrared (FT-IR) spectra were

# (a) (b) (c)

Figure 2. TEM images of (a) OAC-0, (b) OAC-1, and (c) OAC-2.

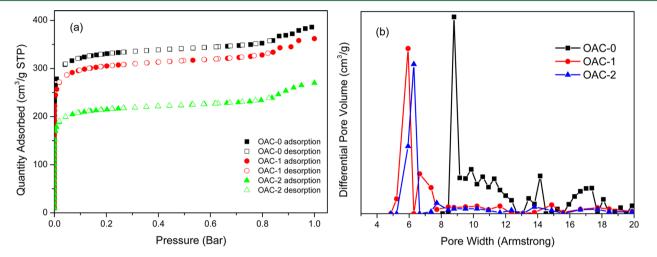


Figure 3. (a)  $N_2$  adsorption isotherms and (b) pore size distributions for OACs.

| Table 1 | . Textual | Properties | of ( | Oxygen-D | oped | Porous | Carbons |  |
|---------|-----------|------------|------|----------|------|--------|---------|--|
|         |           |            |      |          |      |        |         |  |

| sample | $S_{\rm BET}^{a} ({\rm m}^2/{\rm g})$ | $S_{\text{Lang}}^{b}$ (m <sup>2</sup> /g) | $V_{\rm p}^{\ c} \ ({\rm cm}^3/{\rm g})$ | $S_{\rm micro}^{\ \ d} ({\rm m}^2/{\rm g})$ | $V_{\rm micro}^{\ \ d} ({\rm m}^3/{\rm g})$ | $D_{\rm pore}^{e}$ (nm) | oxygen content $^{f}$ (wt %) |
|--------|---------------------------------------|---|--|---|---|-------------------------|------------------------------|
| OAC-0  | 1006                                  | 1567                                      | 0.57                                     | 1304  | 0.51  | 1.1                     | 1.57                         |
| OAC-1  | 925                                   | 1461                                      | 0.53                                     | 1156  | 0.45  | 0.6                     | 7.51                         |
| OAC-2  | 653                                   | 1046                                      | 0.38                                     | 854   | 0.34  | 0.7                     | 3.38                         |
|        |                                       |   |  |   | 1   |                         |                              |

<sup>*a*</sup>The specific surface area was determined by the BET equation  $(p/p_0 = 0.05-0.3)$ . <sup>*b*</sup>The specific surface area was determined by the Langmuir equation  $(p/p_0 = 0.05-0.98)$ . <sup>*c*</sup>Total pore volume at  $p/p_0 = 0.98$ . <sup>*d*</sup>The micropore volume and micropore surface area were estimated as d < 2 nm using the equilibrium model for slit pores. <sup>*c*</sup>The maxima of the pore size distribution was calculated by the NLDFT method. <sup>*f*</sup>The oxygen content was determined by energy-dispersive X-ray (EDS) spectroscopy.

recorded with a PerkinElmer FTIR. Scanning electron microscopy (SEM) images and readings of the oxygen content of the materials were taken using an SEM S-3400NII equipped with energy-dispersive X-ray analysis. Transmission electron microscopy (TEM) images were taken with a Hitachi H-7650. Raman spectra were collected using a Reinshaw Raman microscope with a 632.8 nm (1.96 eV) laser excitation.

Adsorption Measurements. The pure component adsorption isotherms of  $CO_2$ ,  $CH_4$ , and  $N_2$  at three temperatures (273, 298, and 323 K) and gas pressure up to 800 mmHg were measured volumetrically in the Micromeritics ASAP 2020 adsorption apparatus. The degassing procedure was repeated in all samples before measurements at 180 °C for 24 h. The BET (Brunauer–Emmett–Teller) surface area was calculated between 0.05 and 0.3 relative pressure, and pore size distributions were calculated using NLDFT (nonlocal density functional theory) from  $N_2$  adsorption–desorption isotherms measured at 77K.

#### RESULTS AND DISCUSSION

Material Synthesis and Characterization. The FTIR spectra of OAC-1 and OAC-2 are shown in Figure 1a. After the synthetic modification, the adsorption band of C=O becomes more intense at  $1612 \text{ cm}^{-1}$ , and the adsorption band of OH at 2896 cm<sup>-1</sup> is almost invisible on OAC-2, indicating that the conversion from C–OH to C=O is completed. Also, the FTIR pattern of OAC-2 does not show additional adsorption at 2200 cm<sup>-1</sup>, allowing us to determine that the washing process after the grafting is effective. Because the stretching bond at this position could be attributed to the nitrile group (CN) present in DDQ or its reduced form,<sup>27</sup> the defective nature in hightemperature-activated carbon can be confirmed by Raman spectra (Figure 1 b-d). The Raman patterns evidence that two peaks with broad bands at around 1358 and 1588  $\rm cm^{-1}$  can be found, which correspond to the D and G bands, respectively. The D band corresponds to the breath modes of the rings or  $\kappa$ -

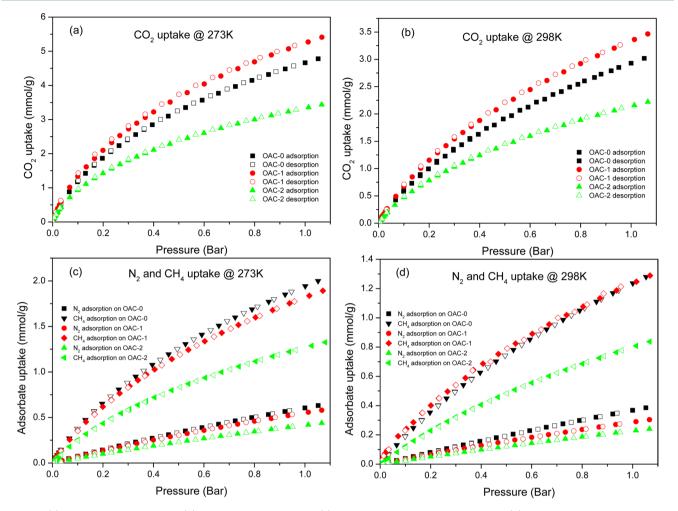


Figure 4. (a) CO<sub>2</sub> adsorption at 273 K, (b) CO<sub>2</sub> adsorption at 298 K, (c) N<sub>2</sub> and CH<sub>4</sub> adsorption at 273 K, and (d) N<sub>2</sub> and CH<sub>4</sub> adsorption at 298 K of OAC-0, OAC-1, and OAC-2.

Table 2. Adsorption Properties of Oxygen-Doped Porous Carbons

|        | CO <sub>2</sub> adsorption<br>(1 bar, mmol/g) |       | CH <sub>4</sub> adsorption (1 bar, mmol/g) |       |      | orption<br>mmol/g) |
|--------|---|-------|--|-------|------|--------------------|
| sample | 0 °C  | 25 °C | 0 °C                                       | 25 °C | 0 °C | 25 °C              |
| OAC-0  | 4.78  | 3.02  | 2.00                                       | 1.28  | 0.63 | 0.39               |
| OAC-1  | 5.41  | 3.46  | 1.89                                       | 1.29  | 0.58 | 0.30               |
| OAC-2  | 3.43  | 2.22  | 1.32                                       | 0.84  | 0.44 | 0.24               |

point phonons of A1g symmetry, while the G band represents the in-plane bond-stretching motion of the pairs of C sp2 atoms (the E2g phonons).<sup>28,29</sup> The intensity ratio between D and G bands ( $I_G/I_D$ ) could exactly reflect the degree of graphitization in carbon materials. The ratio is 0.58, 0.55, and 0.53 for OAC-0, OAC-1, and OAC-2, respectively. The decreasing values indicate that OAC-1 and OAC-2 possess the more defective nature and lower degree of graphitization. This could be attributed to a disordered structure caused by grafting modifications. The XRD patterns are shown in Figure S1 in the Supporting Information; the three labeled peaks at  $2\theta$ values of 32, 34, and 36° could be contributed from the Zn species. The carbon precursors were carbonized at 873 K, while the Zn metal boiling point is about 1181 K; thus, the Zn metal could not be vaporized completely under the carbonization conditions. The empirical parameter (R) was used to evaluated the degree of graphitization, which is defined as the ratio of height of the  $(0 \ 0 \ 2)$  Bragg peak to the background.<sup>30,31</sup> Therefore, the graphitization of OAC-1 and OAC-2 is higher than that of OAC-0 due to its more apparent  $(0 \ 0 \ 2)$  peak.

To establish the impact of the decorated groups on the pore structure, we acquired SEM images for all OACs (Figure S2 in the Supporting Information). A comparison of the microscopy images shows that the addition of the decorated groups does not alter the primary structure of the carbon. TEM images are shown in Figure 2; the micropores are randomly and uniformly distributed on all OACs. The texture properties of the OACs were further examined using nitrogen adsorption isotherms at 77 K. Figure 3 summarizes all of the isotherms and their corresponding pore size distribution using the NLDFT method. It is clear that all isotherms displayed a typical type I shape according to IUPAC classification.<sup>32</sup> The steep adsorption at the low pressure indicates that all of the pores are micropores (<2 nm). As shown in Table 1, which includes the specific surface areas and pore volumes for each OAC, the BET surface area is 1006, 925, and 653  $m^2 g^{-1}$  for OAC-0, OAC-1, and OAC-2, respectively; the pore volume is 0.57, 0.53, and 0.38 cm<sup>3</sup> g<sup>-1</sup> for OAC-0, OAC-1, and OAC-2, respectively. We noted that both the specific surface area and the pore volume significantly decrease after grafting functional groups because the addition groups block partial pore volume. Thus, the pore size is also reduced from 1.1 to 0.6 nm.

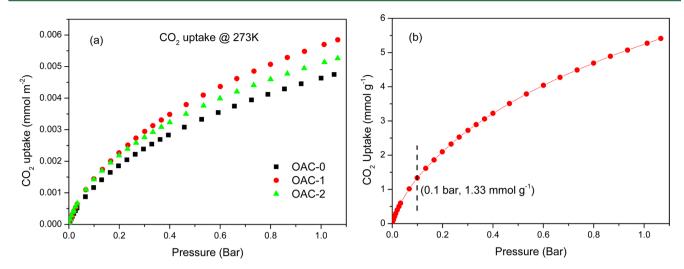


Figure 5. (a) CO<sub>2</sub> adsorption at 273 K of OACs normalized with respect to the surface areas. (b) CO<sub>2</sub> adsorption of OAC-1 at 273 K.

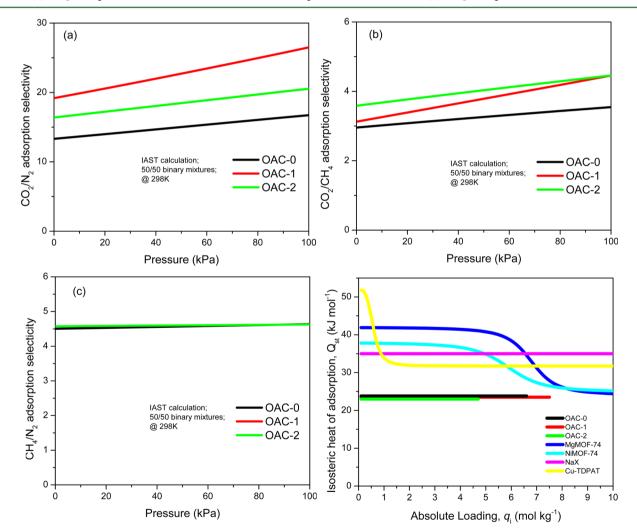


Figure 6. IAST-predicted adsorption selectivities for equimolar binary mixtures of (a)  $CO_2/CH_4$ , (b)  $CH_4/N_2$ , and (c)  $CO_2/N_2$ . (d) Isosteric heats of adsorption for  $CO_2$  on the OACs, MgMOF-74, NiMOF-74, Cu-TDPA, and NaX.

Adsorption Isotherms of CO<sub>2</sub>, CH<sub>4</sub>, and N<sub>2</sub>. The adsorption isotherms of CO<sub>2</sub>, N<sub>2</sub>, and CH<sub>4</sub> were measured at 273, 298, and 323 K and the pressure measured up to 1 bar, as shown in Figure 4, Table 2, and Figures S3 and S4 in the Supporting Information. All isotherms showed excellent

reversibility without any hysteresis, indicating that a vacuum can remove all of the adsorbate molecules during the desorption process. Thus, OACs can be easily regenerated without any other heat energy input. All adsorption isotherms were fitted by the Langmuir–Freundlich model. The

Article

Fractional Uptake (%)

Fractional Uptake (%)

0.0

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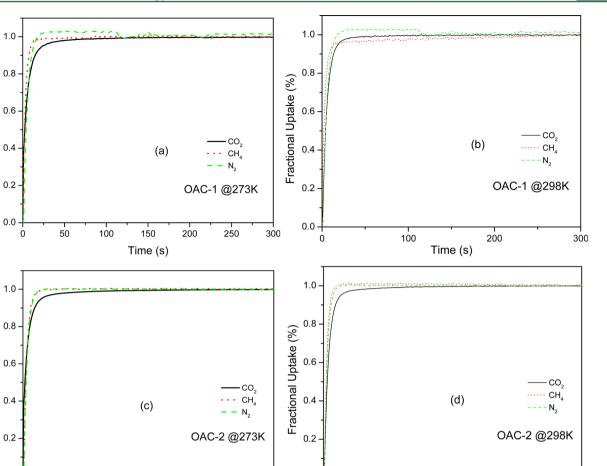


Figure 7. CO<sub>2</sub>, CH<sub>4</sub>, and N<sub>2</sub> fractional uptake on (a) OAC-1 at 273 K, (b) OAC-1 at 298 K, (c) OAC-2 at 273 K, and (d) OAC-2 at 298 K.

300

0.0

0

Table 3. Summary of Diffusion Time Constants of  $CO_2$ ,  $CH_4$ , and  $N_2$  on the Oxygen-Doped Porous Carbons at 273 and 298 K

100

Time (s)

200

| sample | T (k) | $\begin{array}{c} \text{CO}_2 \ D_c / r_c^2 \\ (10^{-3} \ s^{-1}) \end{array}$ | $ \begin{array}{c} \mathrm{CH_4} \ D_{\mathrm{c}} / r_{\mathrm{c}}^2 \\ (10^{-2} \ \mathrm{s}^{-1}) \end{array} $ |     |
|--------|-------|--|---|-----|
| OAC-1  | 273   | 7.2  | 1.2   | 6.2 |
|        | 298   | 7.3  | 1.3   | 6.2 |
| OAC-2  | 273   | 7.7  | 7.1   | 6.2 |
|        | 298   | 7.8  | 8.2   | 7.3 |

calculation detail and parameters are summarized in Tables S1, S2, and S3 in the Supporting Information. It is clear that OACs adsorb CO<sub>2</sub> to a much greater degree than N<sub>2</sub> and CH<sub>4</sub> because of its higher quadrupole moment and polarizability. CH<sub>4</sub> is preferentially adsorbed over N<sub>2</sub> because the polarizability of CH<sub>4</sub> is higher than that of N<sub>2</sub>.<sup>33,34</sup> The highest CO<sub>2</sub> capacity was obtained by OAC-1, measured to be 5.41 and 3.46 mmol g<sup>-1</sup> at 273 and 298 K, respectively. This result is superior to other recently reported activated carbons.<sup>7,35,36</sup> Compared with that of the original material (OAC-0), the CO<sub>2</sub> capacity of OAC-1 increased, whereas the capacity of OAC-2 decreased from 4.78 to 3.43 mmol g<sup>-1</sup> at 1 bar and 273 K. To clarify the effect of oxygen-containing groups decorated on CO<sub>2</sub> uptake capacity, we normalized the CO<sub>2</sub> uptake values with respect to the surface areas (Figure 5a), which clearly showed the enhancement of CO<sub>2</sub> uptake with an increase of O content.

Furthermore, the extremely high  $CO_2$  uptake (1.33 mmol g<sup>-1</sup> at 0.1 bar and 273 K) for OAC-1 confirms the stronger interaction of  $CO_2$  molecules with more O-decorated carbon framework at low pressures. It is noteworthy that the partial pressure of  $CO_2$  in flue gas is approximately 0.1–0.15 bar.

Time (s)

200

300

100

**IAST Selectivities and Interaction Energy.** The selectivity of the preferential adsorption of component 1 over component 2 in a mixture containing both, perhaps in the presence of other components as well, can be formally defined as

$$S_{\rm ads} = \frac{q_1/q_2}{p_1/p_2}$$

In the above equation,  $q_1$  and  $q_2$  are the absolute component loadings of the adsorbed phase in the mixture. These component loadings are also termed the "uptake capacities". In all of the calculations presented below, the calculations of  $q_1$ and  $q_2$  are based on the use of the ideal adsorbed solution theory of Myers and Prausnitz.<sup>37</sup> The accuracy of the IAST calculations for the estimation of the component loadings for several binary mixtures in a wide variety of zeolites and MOFs has been established by comparison with configurational-bias Monte Carlo (CBMC) simulations of mixture adsorption.<sup>38–40</sup> Further details are provided in the Supporting Information. The selectivities of all equimolar binary combination of CO<sub>2</sub>,

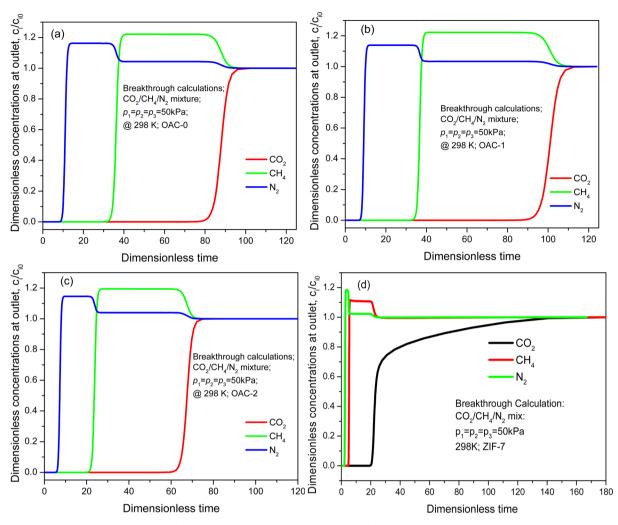


Figure 8. Breakthrough simulation data of (a) OAC-0, (b) OAC-1, (c) OAC-2, and (d) ZIF-7 at 298 K and 1 bar for an equimolar  $CO_2/CH_4/N_2$  mixture ( $c_i$ , concentration of outlet gas;  $c_0$ , concentration of inlet gas).

 $CH_4$ , and  $N_2$  at 298 K are present in Figure 6a–c. For  $CO_2/N_2$ mixtures, the highest adsorption selectivities are obtained by the OAC-1 of 26.5 at 100 kPa, which is significantly higher than that of other adsorbents such as ordered mesoporous carbon (approximately 11.33),<sup>41</sup> MIL-47(V) (approximately 9),<sup>42</sup> and N-doped hierarchical carbons (approximately 8.4)<sup>43</sup> at 298 K. OAC-2 displayed the highest CO<sub>2</sub>/CH<sub>4</sub> selectivity of 4.5 at 298 K and 100 kPa; these values are comparable to those of most existing adsorbents at the same condition, such as silicalite-1 (approximately 2.6),<sup>44</sup> SBA-15 (approximately 5.5),<sup>45</sup> and MCM-41 (approximately 5.5).<sup>46</sup> Modified activated carbons show higher  $CO_2/N_2$  and  $CO_2/CH_4$  selectivities due to their fine pore size and functional groups doping. It is worth noting that the  $CO_2/N_2$  and  $CO_2/CH_4$  selectivities are still increasing in high pressure because N2 and CH4 nearly reach their saturation capacities at 100 kPa, whereas the CO<sub>2</sub> loading is still far from saturation. Therefore, OACs will perform better in high pressure, which is easy to complete in industrial applications. For  $CH_4/N_2$  gas mixtures, the adsorption selectivities of OACs are almost identical. At 298 K and 100 kPa, a CH<sub>4</sub>/N<sub>2</sub> selectivity of 4.6 was obtained, which is higher than IRMOF-1 (approximately 2)<sup>47</sup> and ZIF-69 (approximately 3).48 Those comparisons indicate that OACs are promising adsorbents in the selective adsorption of flue gases and natural gas processing.

The isosteric heats of adsorption for the OACs were accomplished by using adsorption isotherms at 273, 298, and 323 K. Isosteric heat of adsorption is a value with which to directly explore the interaction strength of adsorbates and adsorbents and further elucidate the origin of the enhanced selectivity. The values of  $Q_{st}$  for  $CO_2$ ,  $CH_4$ , and  $N_2$  are provided in Table S4 in the Supporting Information. Figure 6d presents a comparison of the isosteric heats of adsorption of CO<sub>2</sub> in OAC-0, OAC-1, and OAC-2 with data from other published literature for NaX zeolite,<sup>49</sup> MgMOF-74,<sup>49</sup> Cu-TDPAT,<sup>50</sup> and NiMOF-74.<sup>51,52</sup> This implies that the cost of regeneration for OACs will be significantly lower than that of other materials. This is an important advantage in favor of OACs when applied in industry. We noted that the Qst for the three OACs are all mostly the same. The modified activated carbons are supposed to present higher values of adsorption heat; however, the value is not singly determined by decorated groups but also by the specific surface area and pore structure. The impact of the decreased specific area on OAC-1 and OAC-2 may offset the impact of oxygen doping.

Adsorption Kinetics. The adsorption kinetics data of  $CO_2$ ,  $CH_4$ , and  $N_2$  on the OACs were measured at three temperatures (273, 298, and 323 K) and low pressure (approximately 10 mmHg). The fractional uptake curves are presented in Figure 7 and Figure S5 in the Supporting

Information. At 298 K,  $CO_2$  reaches equilibrium in 50 and 60 s for OAC-1 and OAC-2, respectively, whereas  $CH_4$  and  $N_2$  reach equilibrium in 18 and 15 s for OAC-1 and OAC-2, respectively. In addition, doing so took slightly less time at a higher temperature for each gas.

The diffusion time constants  $(D_c/r_c^2, s^{-1})$  can be extracted by fitting the fractional uptake curve with the following micropore diffusion model.<sup>53,54</sup> This model assumes the mass transfer resistance for gas adsorption is controlled by the intracrystalline diffusion. When the fractional uptake is less than 0.85, the curve could be correlated by

$$\frac{m_t}{m_{\infty}} \approx \frac{6}{\sqrt{\pi}} \sqrt{\frac{D_c t}{r_c^2}} - 3 \frac{D_c t}{r_c^2}$$

where  $m_t/m_{\infty}$  is the fractional adsorption uptake,  $D_c$  (m<sup>2</sup>/s) is the intracrystalline diffusivity of gas molecules in porous media,  $r_c$  (m) is the crystal radius, and t (s) is the time. The diffusion time constants of the three gases extracted from the kinetic profiles at three temperatures are summarized in Table 3 and in Table S5 in the Supporting Information. It can be observed from Table 3 that CO<sub>2</sub> and N<sub>2</sub> are almost unable to be kinetically separated, and the kinetic selectivities of CO<sub>2</sub>/CH<sub>4</sub> and CH<sub>4</sub>/N<sub>2</sub> are less than 2. This implies that an efficient kinetic-based adsorption separation is difficult to achieve on the OACs. This is because the pore size is still too large compared to the kinetic diameters of the adsorbates, and the interaction force is not strong enough to hold one adsorbate for a longer time.

Transient Breakthrough Simulations. Separations using porous adsorbents are usually conducted in the fixed bed units; in such cases, the separation performance is dictated by a combination of adsorption selectivity and uptake capacity. For a proper evaluation of the performance of fixed bed adsorbers, it is necessary to carry out transient breakthrough simulations as described in the literature.55,56 We therefore performed transient breakthrough simulations for the separation of  $CO_2/CH_4/N_2$  gas mixtures on OACs, as shown in Figure 8. The methodology details and movies using the adsorption cycle isotherm fits are provided in the Supporting Information; the methodology used has been verified to be of good accuracy.<sup>5</sup> For all OACs, the sequence of breakthroughs is  $N_{2}$ ,  $CH_4$ , and CO<sub>2</sub>. This is dictated by the hierarchy of adsorption strengths. CO<sub>2</sub> molecules are strongly adsorbed by OACs, which elute last in the sequence. The longest retention time was obtained on OAC-1, and this material has the best separation performance due to it having the highest  $CO_2$  capacity and highest  $CO_2/N_2$ and CO<sub>2</sub>/CH<sub>4</sub> adsorption selectivities. We noted that there is a time interval in which we can get pure CH<sub>4</sub> or N<sub>2</sub> using OACs. We compare OACs with ZIF-7 using a breakthrough simulation under the same condition. The CO<sub>2</sub> and CH<sub>4</sub> adsorption data of ZIF-7 are from one recently published paper by our group, and the N<sub>2</sub> adsorption is obtained from a U.S. patent.<sup>39</sup> It is clear that the time interval of  $CO_2/CH_4/N_2$  pairs on ZIF-7 is much shorter than that of OACs. This implies OACs are more efficient and promising than ZIF-7 in industrial applications.

#### ASSOCIATED CONTENT

#### **S** Supporting Information

Figures showing the XRD patterns of OACs; representative SEM images of materials showing pore structure;  $CO_2$ ,  $CH_4$ , and  $N_2$  adsorption and fractional uptake; and a shematic of the breakthrough apparatus. Tables showing the 1-site Langmuir

parameters for the adsorption of  $CO_2$ ,  $CH_4$ , and  $N_2$ ; isoteric heats; and a summary of diffusion time constants. Additional details on IAST calculations, the isoteric heat of adsorption, simulation methodology for transient breakthrough in fixed bed adsorbers, notations, and references. Movie showing the adsorption cycle isotherm fit. The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.est.5b01652.

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#### Notes

The authors declare no competing financial interest.

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# Hydroquinone and Quinone-grafted Porous Activated Carbons for Highly Selective CO<sub>2</sub> Capture from Flue Gases and Natural Gas Upgrading

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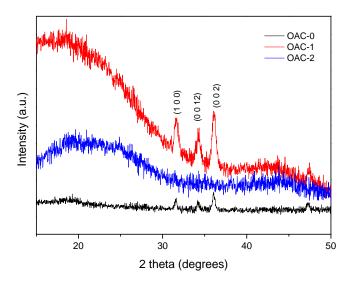
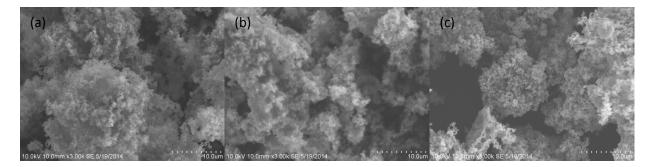


Figure S1. XRD patterns of OAC-1, OAC-1 and OAC-2.



**Figure S2.** Representative SEM images of the a) OAC-1, b) OAC-1 and c) OAC-2 showing the similarity in the pore structure of the carbon pre- and post-modification.

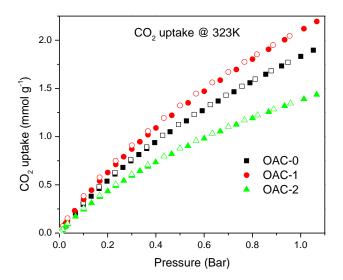


Figure S3. CO<sub>2</sub> adsorption 323K of OAC-0, OAC-1, and OAC-2.

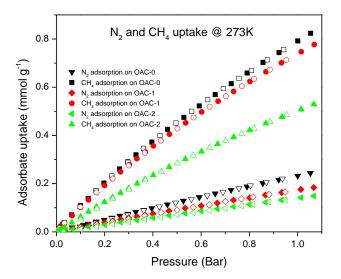


Figure S4. (a) CH<sub>4</sub> and N<sub>2</sub> adsorption 323K of OAC-0, OAC-1, and OAC-2.

## Fitting of pure component isotherms

The experimentally measured loadings for (a)  $CO_2$ , (b)  $CH_4$ , and (c)  $N_2$  were measured as a function of the absolute pressure at three different temperatures 273 K, 298 K, and 323 K.

The isotherm data for all three guest molecules in OAC-0, OAC-1, and OAC-2 were fitted with the Langmuir model

$$q = q_{sat} \frac{bp}{1 + bp} \tag{1}$$

with T-dependent parameter b

$$b = b_0 \exp\left(\frac{E}{RT}\right) \tag{2}$$

The Langmuir fit parameters are provided in Table S1, Table S2, and Table S3.

| т ·       |               | 1             | C C C ·               | different OACs.     |
|-----------|---------------|---------------|-----------------------|---------------------|
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|           |               |               |                       |                     |
|           |               |               |                       |                     |

|       | $q_{\rm sat}$        | $b_0$                  | Ε                    |
|-------|----------------------|------------------------|----------------------|
|       | mol kg <sup>-1</sup> | Pa <sup>-1</sup>       | kJ mol <sup>-1</sup> |
| OAC-0 | 6.7                  | 5.5×10 <sup>-10</sup>  | 23.8                 |
| OAC-1 | 7.6                  | 6.34×10 <sup>-10</sup> | 23.5                 |
| OAC-2 | 4.8                  | 8.16×10 <sup>-10</sup> | 23                   |

 Table S6. 1-site Langmuir parameters for CH<sub>4</sub> in different OACs.

|       | $q_{ m sat}$         | $b_0$                  | Ε                    |
|-------|----------------------|------------------------|----------------------|
|       | mol kg <sup>-1</sup> | Pa <sup>-1</sup>       | kJ mol <sup>-1</sup> |
| OAC-0 | 3.75                 | 1.67×10 <sup>-9</sup>  | 19.8                 |
| OAC-1 | 3                    | 2.11×10 <sup>-9</sup>  | 20                   |
| OAC-2 | 2.5                  | 1.802×10 <sup>-9</sup> | 20.9                 |

Table S7. 1-site Langmuir parameters for  $N_2$  in different OACs.

|       | $q_{ m sat}$         | $b_0$                  | Ε        |
|-------|----------------------|------------------------|----------|
|       | mol kg <sup>-1</sup> | Pa <sup>-1</sup>       | kJ mol⁻¹ |
| OAC-0 | 3.1                  | 2.08×10 <sup>-9</sup>  | 16       |
| OAC-1 | 2.9                  | 4.18×10 <sup>-10</sup> | 19.6     |
| OAC-2 | 2.3                  | 6.65×10 <sup>-10</sup> | 18.4     |

#### **IAST calculations**

The adsorption selectivity for the mixtures  $CH_4/N_2$  and  $CO_2/CH_4$  defined by

$$S_{ads} = \frac{q_1/q_2}{p_1/p_2}$$

were calculated according to IAST model proposed by Myers [1,2]. In above equation,  $q_1$  and  $q_2$  are the absolute component loadings of the adsorbed phase in the mixture. These component loadings are also termed the uptake capacities.

#### Isosteric heat of adsorption

The isosteric heat of adsorption,  $Q_{\rm st}$ , defined as

$$Q_{st} = RT^2 \left(\frac{\partial \ln p}{\partial T}\right)_q$$

were determined using the pure component isotherm fits using the Clausius-Clapeyron equation, where Qst (kJ/mol) is the isosteric heat of adsorption, T (K) is the temperature, p (kPa) is the pressure, R is the gas constant. The values of  $Q_{st}$  for CO<sub>2</sub>, CH<sub>4</sub>, and N<sub>2</sub> are provided in Table S4 of OACs.

Table S8. Isosteric heats of adsorption of  $CO_2$ ,  $CH_4$  and  $N_2$  in different OACs.

|       | Q <sub>st</sub> , CO <sub>2</sub> | $Q_{\rm st}$ , CH <sub>4</sub> | Q <sub>st</sub> , N <sub>2</sub> |
|-------|-----------------------------------|--------------------------------|----------------------------------|
|       | kJ mol⁻¹                          | kJ mol <sup>-1</sup>           | kJ mol <sup>-1</sup>             |
| OAC-0 | 23.8                              | 19.8                           | 16                               |
| OAC-1 | 23.5                              | 20                             | 19.6                             |
| OAC-2 | 23                                | 20.9                           | 18.4                             |

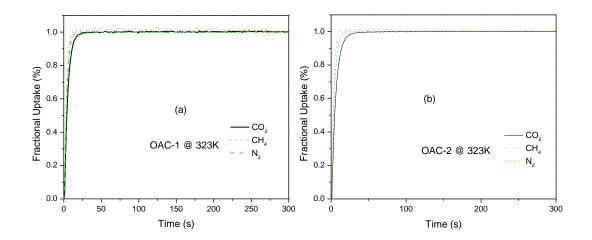


Figure S5. CO<sub>2</sub>, CH<sub>4</sub>, and N<sub>2</sub> fractional uptake on (a) OAC-1, (b) OAC-2 at 323 K.

| Sample | $\begin{array}{c} \text{CO}_2  \text{D}_c/\text{r}^2_c \\ (10^{-3}  \text{s}^{-1}) \end{array}$ | $\begin{array}{c} CH_4 D_c / r_c^2 \\ (10^{-2} s^{-1}) \end{array}$ | $\frac{N_2 D_c/r_c^2}{(10^{-3} s^{-1})}$ |
|--------|---|---|--|
| OAC-1  | 7.5   | 1.9   | 8.1                                      |
| OAC-2  | 8.2   | 1.1   | 7.4                                      |

Table S5. Summary of Diffusion Time Constants of CO<sub>2</sub>, CH<sub>4</sub>, and N<sub>2</sub> on the OACs at 323 K.

#### Simulation methodology for transient breakthrough in fixed bed adsorbers

The separation of CO<sub>2</sub>/CH<sub>4</sub>, CO<sub>2</sub>/N<sub>2</sub>, and CH<sub>4</sub>/N<sub>2</sub> mixtures is commonly carried out in fixed bed adsorbers in which the separation performance is dictated by a combination of three separate factors: (a) adsorption selectivity, (b) uptake capacity, and (c) intra-crystalline diffusivities of guest molecules within the pores. Transient breakthrough simulations are required for a proper evaluation of microporous materials; the simulation methodology used in our work is described in earlier publications [3,4].

In order to evaluate the different OACs, breakthrough calculations were performed taking the following parameter values: inside diameter of tube = 50 mm; bed length, L = 1.8 m; voidage of bed,  $\varepsilon = 0.5$ ; superficial gas velocity, u = 0.05 m/s (at inlet), interstitial velocity, v = 0.1 m/s. The mass of adsorbent packed in the tube is 2 kg; see schematic in Figure S6.

A brief summary of the simulation methodology is presented below. Assuming plug flow of an n-component gas mixture through a fixed bed maintained under isothermal conditions (see schematic in Figure 0), the partial pressures in the gas phase at any position and instant of time are obtained by solving the following set of partial differential equations for each of the species i in the gas mixture [5].

$$\frac{1}{RT}\frac{\partial p_i(t,z)}{\partial t} = -\frac{1}{RT}\frac{\partial \left(v(t,z)p_i(t,z)\right)}{\partial z} - \frac{\left(1-\varepsilon\right)}{\varepsilon}\rho\frac{\partial \bar{q}_i(t,z)}{\partial t}; \quad i = 1,2,...n$$
(1)

In equation (1), *t* is the time, *z* is the distance along the adsorber,  $\rho$  is the framework density,  $\varepsilon$  is the bed voidage, *v* is the interstitial gas velocity, and  $\overline{q}_i(t, z)$  is the *spatially averaged* molar loading within the crystallites of radius  $r_c$ , monitored at position *z*, and at time *t*.

At any time t, during the transient approach to thermodynamic equilibrium, the spatially averaged molar loading within the crystallite  $r_c$  is obtained by integration of the radial loading profile

$$\bar{q}_{i}(t) = \frac{3}{r_{c}^{3}} \int_{0}^{r_{c}} q_{i}(r,t) r^{2} dr$$
(2)

For transient unary uptake within a crystal at any position and time with the fixed bed, the radial distribution of molar loadings,  $q_i$ , within a spherical crystallite, of radius  $r_c$ , is obtained from a solution of a set of differential equations describing the uptake

$$\frac{\partial q_i(r,t)}{\partial t} = -\frac{1}{\rho} \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 N_i \right)$$
(3)

The molar flux  $N_i$  of component *i* is described by the simplified version of the Maxwell-Stefan equations in which both correlation effects and thermodynamic coupling effects are considered to be of negligible importance [4]

$$N_i = -\rho \mathcal{D}_i \frac{\partial q_i}{\partial r} \tag{4}$$

Summing equation (2) over all *n* species in the mixture allows calculation of the *total average* molar loading of the mixture within the crystallite

$$\bar{q}_{t}(t,z) = \sum_{i=1}^{n} \bar{q}_{i}(t,z)$$
(5)

The interstitial gas velocity is related to the superficial gas velocity by

$$v = \frac{u}{\varepsilon} \tag{6}$$

In industrial practice, the most common operation uses a step-wise input of mixtures to be separated into an adsorber bed that is initially free of adsorbates, i.e. we have the initial condition

$$t = 0; \quad q_i(0, z) = 0$$
 (7)

At time, t = 0, the inlet to the adsorber, z = 0, is subjected to a step input of the *n*-component gas mixture and this step input is maintained till the end of the adsorption cycle when steady-state conditions are reached.

$$t \ge 0; \quad p_i(0,t) = p_{i0}; \quad u(0,t) = u_0$$
(8)

where  $u_0$  is the superficial gas velocity at the inlet to the adsorber.

The breakthrough characteristics for any component is essentially dictated by two sets of parameters: (a) The characteristic contact time  $\frac{L}{v} = \frac{L\varepsilon}{u}$  between the crystallites and the surrounding fluid phase, and (b)  $\frac{D_i}{r_c^2}$ , that reflect the importance of intra-crystalline diffusion limitations. It is common to use the dimensionless time,  $\tau = \frac{tu}{L\varepsilon}$ , obtained by dividing the actual time, *t*, by the characteristic time,  $\frac{L\varepsilon}{u}$  when plotting simulated breakthrough curves [3].

If the value of  $\frac{D_i}{r_c^2}$  is large enough to ensure that intra-crystalline gradients are absent and the

entire crystallite particle can be considered to be in thermodynamic equilibrium with the surrounding bulk gas phase at that time t, and position z of the adsorber

$$\overline{q}_i(t,z) = q_i(t,z) \tag{9}$$

The molar loadings at the *outer surface* of the crystallites, i.e. at  $r = r_c$ , are calculated on the basis of adsorption equilibrium with the bulk gas phase partial pressures  $p_i$  at that position z and time t. The adsorption equilibrium can be calculated on the basis of the IAST. The assumption of thermodynamic equilibrium at every position z, and any time t, i.e. invoking Equation (5), generally results in sharp breakthroughs for each component. Sharp breakthroughs are desirable in practice because this would result in high productivity of pure products. Essentially, the influence of intra-crystalline diffusion is to reduce the productivity of pure gases. For all the breakthrough calculations reported in this work, we assume negligible diffusion resistances for all materials and we invoke the simplified Equation (5).

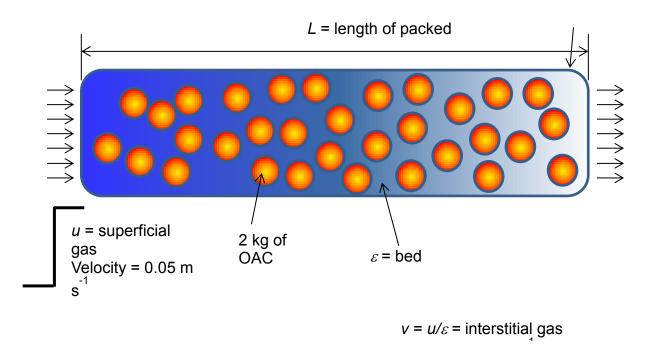


Figure S6. Schematic of the breakthrough apparatus.

## Notation

| b                   | Langmuir-Freundlich constant for species i at adsorption site A, $Pa^{-\nu_i}$           |
|---------------------|--|
| Ci                  | molar concentration of species i in gas mixture, mol m <sup>-3</sup>                     |
| C <sub>i0</sub>     | molar concentration of species i in gas mixture at inlet to adsorber, molm <sup>-3</sup> |
| Ε                   | energy parameter, J mol <sup>-1</sup>  |
| L                   | length of packed bed adsorber, m   |
| N <sub>i</sub>      | molar flux of species i, mol m <sup>-2</sup> s <sup>-1</sup>                             |
| ρ <sub>i</sub>      | partial pressure of species i in mixture, Pa   |
| $ ho_{ m t}$        | total system pressure, Pa  |
| q <sub>i</sub>      | component molar loading of species i, mol kg <sup>-1</sup>                               |
| $\overline{q}_i(t)$ | spatially averaged component molar loading of species i, mol kg <sup>-1</sup>            |
| Q <sub>st</sub>     | isosteric heat of adsorption, J mol <sup>-1</sup>  |
| r <sub>c</sub>      | radius of crystallite, m   |
| R                   | gas constant, 8.314 J mol <sup>-1</sup> K <sup>-1</sup>                                  |
| t                   | time, s  |
| Т                   | absolute temperature, K  |
| u                   | superficial gas velocity in packed bed, m s <sup>-1</sup>                                |
| v                   | interstitial gas velocity in packed bed, m s <sup>-1</sup>                               |
| Greek letters       |  |
| ε                   | voidage of packed bed, dimensionless   |
| ρ                   | framework density, kg m <sup>-3</sup>  |
| τ                   | time, dimensionless  |

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